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ANL-CANDID,

A Two-dimensional, Diffusion-theory Code
Based on CANDID2D

by

G. K. Leaf, A. S. Kennedy, and G. C. Jensen

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Applied Mathematics Division

September 1967

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ABSTRACT

ANL-CANDID is an extension and modification of the two-dimensional diffusion-theory code CANDID2D developed by Computer Applications Incorporated for use on the Control Data Corporation 3600 computer. The code, as delivered, was capable of performing a reactivity calculation as well as geometry, composition, and buckling searches. These calculations were limited to rz geometry. In addition, the code could perform an up-scattering calculation. The present code is the result of extending and modifying the C.A.I. code. The extensions include the capability of performing calculations in xy and $r\theta$ geometry, including the full periodic case. Adjoint and source calculations were added with an acceleration procedure for the source calculation. In addition, an α calculation was added. The use of two-term Chebyshev extrapolation was retained; however, the strategy for employing it was changed. In addition, the strategy employed in the search procedure was completely changed. Mesh refinement and restart capability have been provided.

I. FINITE-DIFFERENCE EQUATIONS

A. Derivation of Equations for $r\theta$ Geometry

We shall begin our discussion by deriving the finite-difference equations for $r\theta$ geometry. The derivation is based on the usual first-order approximation and is consistent with the derivation for rz geometry. In the absence of external sources, the system of multigroup diffusion equations can be written in the form

$$-\operatorname{div}(D_{g}\nabla\phi_{g}) + (\sigma_{R}^{g} + D_{g}B^{2}) \phi_{g} = \sum_{g'\neq g} \sigma^{g'g}\phi_{g'}$$

$$+ \frac{1}{k} \sum_{g'=1}^{G} \chi^{g'g}(\nu\sigma_{f})_{g'} \phi_{g'}, \quad g = 1, 2, ..., G, \qquad (1.1)$$

with external boundary conditions having the form

$$A_{g}\nabla\phi_{g}\cdot\overline{n}+B_{g}\phi_{g}=C_{g}, \qquad (1.2)$$

where B denotes the transverse buckling, and \overline{n} is the unit normal for the exterior surface of the reactor. The reactor configuration is assumed to be composed of regions such that the macroscopic cross sections are constant within each region.

The object of the continuous problem described above is to find the unique positive flux $\{\phi_g\}_{g=1}^G$ and the corresponding positive number k for which the system 1.1 has a solution subject to the conditions 1.2. In practice, we replace the continuous system by a finite-difference approximation and solve the resulting system. To this end, let the reactor domain be defined by $R_L \leq R \leq R_R$, $\theta_B \leq \theta \leq \theta_T$. In addition, we lay down a mesh $R_L = R_0 < R_1 < \ldots < R_I = R_R$, and $\theta_B = \theta_0 < \theta_1 < \ldots < \theta_J = \theta_T$. Here we include the possibility of $\theta_B = 0$, $\theta_T = 2\pi$, which is the full periodic case. The mesh is assumed to be laid down in such a manner that region boundaries occur only along mesh lines. Consider an interior cell centered at (r_i,θ_j) , as shown in Fig. 1. Integrate each member of the system 1.1 over this cell, apply Green's Theorem to the divergence term, and make the approximation

$$\int_{\theta_{j-1}}^{\theta_{j}} \int_{R_{i-1}}^{R_{i}} \phi_{g}(\mathbf{r},\theta) \mathbf{r} d\mathbf{r} d\theta \approx \phi_{g}(\mathbf{r}_{i},\theta_{j}) V_{ij} = \phi_{ijg} V_{ij}, \qquad (1.3)$$

where

$$V_{ij} = (\Delta \theta_j / 2)(R_i^2 - R_{i-1}^2)$$

is the volume of the cell and

$$\triangle \theta_{\mathbf{j}} = \Theta_{\mathbf{j}} - \Theta_{\mathbf{j}-1}.$$

Having done this, we obtain the system

$$-\int D_{g} \frac{\partial \phi_{g}}{\partial n} d\Omega + \left(\sigma_{R}^{g} + D_{g}B^{2}\right) V_{ij} \phi_{ijg} = \sum_{g' \neq g} \sigma^{g'g} V_{ij} \phi_{ijg}$$

$$+ \frac{1}{k} \sum_{g'=1}^{G} \chi^{g'g} (\nu \sigma_{f})_{g'} \phi_{ijg'}. \quad (1.4)$$

The surface integral extends over the four surfaces of the (i,j)th cell. Along the surfaces extending from points 1 to 2 and 3 to 4 in Fig. 1, we have $\partial \phi/\partial n = \partial \phi/\partial r$. Region boundaries may lie along any of these surfaces, so that ordinary interpolation will not suffice for even a first-order approximation. However, if we impose the conditions that the flux and current are continuous across the region boundaries and if we neglect the dependence of $\partial \phi/\partial r$ on θ , we can approximate two of the surface integrals as follows:³

a.
$$\int_{1}^{2} D \frac{\partial \phi}{\partial \mathbf{r}} d\Omega \cong \left(\frac{\hat{\mathbf{D}}}{\ell}\right)_{\mathbf{E}} (\phi_{\mathbf{E}} - \phi_{0}) R_{\mathbf{i}} \Delta \theta_{\mathbf{j}},$$
b.
$$\int_{3}^{4} D \frac{\partial \phi}{\partial \mathbf{r}} d\Omega \cong \left(\frac{\hat{\mathbf{D}}}{\ell}\right)_{\mathbf{W}} (\phi_{\mathbf{W}} - \phi_{0}) R_{\mathbf{i}-1} \Delta \theta_{\mathbf{j}}.$$
(1.5)

Here we have dropped the group index, and we are using the subscripts 0, E, and W for the subscripts (i,j), (i+1,j) and (i-1,j), respectively. The linkage coefficients are defined by

a.
$$\left(\frac{\ell}{\hat{\Delta}}\right)_{E} = \frac{\frac{1}{2}\Delta R_{i}}{D_{0}} + \frac{\frac{1}{2}\Delta R_{i+1}}{D_{E}},$$
b. $\left(\frac{\ell}{\hat{\Delta}}\right)_{W} = \frac{\frac{1}{2}\Delta R_{i-1}}{D_{W}} + \frac{\frac{1}{2}\Delta R_{i}}{D_{0}},$

$$(1.6)$$

where D_{0} , D_{E} , and D_{W} denote the values of the diffusion lengths in the respective cells.

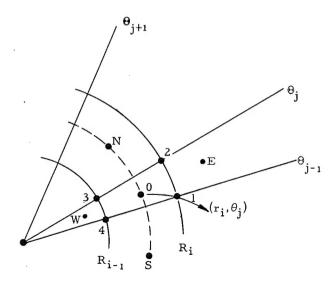


Fig. 1 Interior Cell, Centered at (r_i,θ_j)

For the remaining two surfaces, $\partial \phi/\partial n = (1/r)(\partial \phi/\partial \theta)$, and an approximation is sought for this tangential derivative along the two surfaces. Consider the surface extending from points 2 to 3, as shown in Fig. 2. To approximate the tangential derivative at the region boundary, we first approximate the derivative of ϕ in the direction of the chord joining the points 0 and N. This is done just as before, where now

$$\left(\frac{\ell}{\hat{D}}\right)_{N} = \frac{\ell_{N}}{D_{N}} + \frac{\delta_{0}}{D_{0}}.$$
 (1.7)

This approximation $(\hat{D}/\ell)_N$ $(\phi_N - \phi_0)$, is then projected in the direction of the tangential derivative and the resulting component,

$$\left(\frac{\hat{\mathbf{D}}}{\ell}\right)_{\mathbf{N}} (\phi_{\mathbf{N}} - \phi_{\mathbf{0}}) \cos \frac{\Delta \theta_{\mathbf{j}+1} - \Delta \theta_{\mathbf{j}}}{2},$$

is taken as the approximation to the tangential derivative. Thus,

$$\int_{\mathbf{z}}^{3} \mathbf{D} \frac{\partial \phi}{\partial \mathbf{n}} d\Omega \stackrel{\sim}{=} \left(\frac{\dot{\mathbf{D}}}{\ell}\right)_{\mathbf{N}} (\phi_{\mathbf{N}} - \phi_{\mathbf{0}}) \Delta \mathbf{R}_{\mathbf{i}} \cos \frac{\Delta \theta_{\mathbf{j}+1} - \Delta \theta_{\mathbf{j}}}{2}. \tag{1.8}$$

If we set $\delta_U = \delta_0 + \ell_N$, then

$$\delta_{\mathbf{U}} = (\mathbf{R_i} + \mathbf{R_{i-1}}) \sin \frac{\Delta\theta_{\mathbf{j}} + \Delta\theta_{\mathbf{j+1}}}{4}, \qquad (1.9)$$

and

$$\ell_{\mathbf{N}} = \frac{\delta_{\mathbf{U}} \omega_{\mathbf{U}}}{1 + \omega_{\mathbf{U}}}, \quad \delta_{\mathbf{0}} = \frac{\delta_{\mathbf{U}}}{1 + \omega_{\mathbf{U}}}, \quad (1.10)$$

where

$$\omega_{\mathbf{U}} = \frac{\sin \frac{1}{2} \Delta \theta_{\mathbf{j}+1}}{\sin \frac{1}{2} \Delta \theta_{\mathbf{j}}}.$$

The surface integral over the surface from 3 to 4 is handled in an analogous manner. Having finished the case of interior mesh cells, we must now consider the case when one or more of the surfaces of the cell are part of the exterior boundary when a boundary condition of the form $A \partial \phi / \partial n + B\phi = C$ is imposed. If, for example, the surface from 1 to 2 is an exterior surface, we obtain in the usual fashion³

$$\int_{1}^{2} D \frac{\partial \phi}{\partial n} d\Omega \simeq \frac{D_{0}(C - B\phi_{0})}{A + \frac{1}{2} \Delta R_{i} B} R_{i} \Delta \theta_{j}, \qquad (1.11)$$

with an analogous expression for the integral from 3 to 4. If, for example, the surface from 2 to 3 is a part of the exterior surface, then $\Delta\theta_{j+1}=0$; therefore, $\ell_N=0$, $\delta_U=\delta_0$, and $\omega_U=0$. Thus,

$$\int_{3}^{3} D \frac{\partial \phi}{\partial n} d\Omega \approx \frac{D_{0}(C - B\phi_{0}) \Delta R_{i}}{A + \delta_{0}B} \cos \frac{1}{2} \Delta \theta_{j}, \qquad (1.12)$$

where

$$\delta_0 = (R_i + R_{i-1}) \sin \frac{1}{4} \Delta \theta_i.$$

Note that if an inhomogeneous boundary condition (C \neq 0) is present, then a source term is present and the relevant calculation is a source calculation.

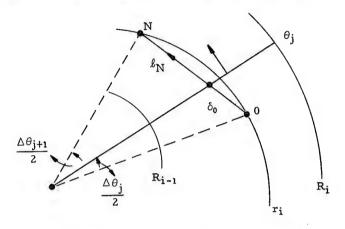


Fig. 2

Detail of Interface between Two Interior Cells

Having approximated the four surface integrals for the case of both interior and boundary points, we can make the following definitions for an interior cell:

$$a_{ij} = \left(\frac{\hat{D}}{\ell}\right)_{E} R_{i} \Delta \theta_{j},$$

$$b_{ij} = \left(\frac{\hat{D}}{\ell}\right)_{N} \Delta R_{i} \cos \frac{\Delta \theta_{j+1} - \Delta \theta_{j}}{2},$$

$$c_{ij} = a_{i-1,j},$$

$$d_{ij} = b_{i,j-1}.$$
(1.13)

At the boundary cells, we make the following definitions:

$$a_{I,j} = \frac{DBR_{I} \Delta \theta_{j}}{A + \frac{1}{2} \Delta R_{I}B}, \quad 1 \leq j \leq J;$$

$$b_{iJ} = \frac{DB\Delta R_{i}}{A + \delta_{0} B} \cos \frac{1}{2} \Delta \theta_{J}, \quad 1 \leq i \leq I;$$

$$c_{1,j} = \frac{DBR_{0} \Delta \theta_{j}}{A + \frac{1}{2} \Delta R_{1}B}, \quad 1 \leq j \leq J;$$

$$d_{i,1} = \frac{DB\Delta R_{i}}{A + \ell_{0} B} \cos \frac{1}{2} \Delta \theta_{i}, \quad 1 \leq i \leq I.$$

$$(1.14)$$

and

and

In the full periodic case, $\theta_0=0$ and $\theta_J=2\pi$; i.e., the top and bottom boundaries coincide. Thus, in this case we have

$$\mathbf{d}_{i1} \equiv \mathbf{b}_{i,I}, \qquad 1 \le i \le I. \tag{1.15}$$

Here we have dropped the group and region index on the quantities D, A, and B. Finally, if we set

$$e_{ijg} = a_{ijg} + b_{ijg} + e_{ijg} + d_{ijg} + (\sigma_R^g + D_g B^2) V_{ij}$$

the finite-difference system 1.4 takes the form

$$-(a_{ijg}\phi_{i+1,j,g} + b_{ijg}\phi_{i,j+1,g} + c_{ijg}\phi_{i-1,j,g} + d_{ijg}\phi_{i,j-1,g})$$

$$+ e_{ijg}\phi_{ijg} = \sum_{g' \neq g} \sigma^{g'g} V_{ij}\phi_{ijg} + \frac{1}{k} \sum_{g'=1}^{G} \chi^{g'g}(\nu\sigma_f)_{g'} V_{ij}\phi_{ijg}.$$
 (1.16)

Here, since we have assumed that $C \equiv 0$, the calculation is a reactivity calculation rather than a source calculation. The case of a source calculation will be discussed in Section IV.

B. Finite-difference Equations XY Geometry

For completeness, we shall exhibit the expressions for the finitedifference coefficients for XY geometry. Let the reactor domain be defined by

$$X_{L} \le x \le X_{R}, Y_{B} \le y \le Y_{T}.$$

Again we construct a mesh

$$\mathbf{X_L}$$
 = $\mathbf{X_0} < \mathbf{X_1} < \ldots < \mathbf{X_I}$ = $\mathbf{X_R}$, $\mathbf{Y_B}$ = $\mathbf{Y_0} < \mathbf{Y_1} < \ldots < \mathbf{Y_J}$ = $\mathbf{Y_T}$

in such a way that region boundaries lie along mesh lines. Letting

$$\Delta X_i = X_i - X_{i-1}$$
 and $\Delta Y_j = Y_j - Y_{j-1}$

we set

$$a_{ijg} = \frac{2D_{ijg}D_{i+1,j,g}^{\Delta Y_{j}}}{\Delta X_{i}D_{i+1,j,g} + \Delta X_{i+1}D_{i,j,g}},$$

$$b_{ijg} = \frac{2D_{ijg}D_{i,j+1,g}^{\Delta X_{i}}}{\Delta Y_{j}D_{i,j+1,g} + \Delta Y_{j+1}D_{ijg}},$$

$$c_{ijg} = a_{i-1,j,g}, \text{ and } d_{ijg} = b_{i,j-1,g}$$
(1.17)

at interior mesh cells. At the boundaries, we set

$$a_{Ijg} = \frac{D_{Ijg}B_{Ijg}\Delta Y_{j}}{A_{Ijg} + \frac{1}{2}\Delta X_{I}B_{Ijg}}, \qquad 1 \le j \le J;$$

$$b_{iJg} = \frac{D_{iJg}B_{iJg}\Delta X_{i}}{A_{iJg} + \frac{1}{2}\Delta Y_{J}B_{iJg}}, \qquad 1 \le i \le I;$$

$$c_{ijg} = \frac{D_{ijg}B_{ijg}\Delta Y_{j}}{A_{ijg} + \frac{1}{2}\Delta X_{i}B_{ijg}}, \qquad 1 \le j \le J;$$
 and
$$d_{iig} = \frac{D_{iig}B_{iig}\Delta X_{i}}{A_{iig} + \frac{1}{2}\Delta Y_{i}B_{iig}}, \qquad 1 \le i \le I;$$

where the boundary constants A_{ijg} and B_{ijg} are permitted to vary by region and group. The form of the resulting finite-difference equation is the same as that in $r\theta$ or rz geometry.

II. ADJOINT CALCULATION

The differential operator associated with the partial-differential system 1.1 and 1.2, when viewed as an operator in the space L2 over the reactor, has an adjoint operator in the same space. If the boundary conditions 1.2 are homogeneous, then the domain of the adjoint operator coincides with that of the real operator. Thus the functions that lie in the domain of the adjoint operator satisfy the same boundary conditions as the functions that lie in the domain of the real operator. Hence, the adjoint operator is, in this case, just the formal adjoint operator. Consequently the finite-difference approximation to the adjoint can be found by forming the transpose of the real finite-difference equations. Before displaying the transposed system, we shall put the finite-difference equations in the form of a matrix equation. The ordering of the space will be that of channel ordering,3 which is just a permutation of the usual group ordering. The ordering is defined in the following way: For each i, $1 \le i \le I$, the set of points $I_i = \{(i,j,g): 1 \le j \le J, 1 \le g \le G\}$ is called a channel. The points within a channel are ordered first with respect to the energy groups; then within each energy group they are ordered with respect to the points in the channel. Thus the flux vector ϕ is partitioned in the following manner:

$$\phi = (\phi_{1}, \phi_{2}, ..., \phi_{I})',$$

$$\phi_{i} = (\phi_{i1}, \phi_{i2}, ..., \phi_{iG})', 1 \leq i \leq I;$$

$$\phi_{ig} = (\phi_{ig1}, \phi_{ig2}, ..., \phi_{igJ})', 1 \leq i \leq I, 1 \leq g \leq G;$$

$$(2.1)$$

where the prime denotes the transpose, since all vectors in this discussion are column vectors. Relative to this ordering of the flux space, we shall define the following matrices:

(In the full periodic case, Jig has the element -bigJ appearing in the extreme upper right- and lower left-hand corners.)

$$K_{ig} = \operatorname{diag} \left[a_{ig1}, a_{ig2}, ..., a_{igJ} \right];$$

$$B_{i}^{g'g} = \operatorname{diag} \left[V_{i1} \sigma_{i1}^{g'g}, V_{i2} \sigma_{i2}^{g'g}, ..., V_{iJ} \sigma_{iJ}^{g'g} \right];$$

$$F_{i}^{g'g} = \operatorname{diag} \left[V_{i1} \sigma_{i1}^{g'g} (\nu \sigma_{f})_{i1}^{g'}, ..., V_{iJ} \sigma_{iJ}^{g'g} (\nu \sigma_{f})_{iJ}^{g'} \right].$$
(2.3)

Then, collecting the above matrices, we define

$$a. \ J_i = \begin{bmatrix} J_{i1} & -B_i^{21} & \cdots & -B_i^{G1} \\ B_i^{12} & J_{i2} & & & \\ \vdots & & \ddots & & \\ -B_i^{1G} & \cdots & & & J_{iG} \end{bmatrix},$$
 b. $K_i = diag \begin{bmatrix} K_{i1}, K_{i2}, \dots, K_{iG} \end{bmatrix},$ and
$$c. \ F_i = \begin{bmatrix} F_i^{11} & F_i^{21} & \cdots & F_i^{G1} \\ F^{12} & F_i^{22} & \cdots & & \\ \vdots & & & & \\ \vdots & & & & & \\ F_i^{1G} & \cdots & & F_i^{GG} \end{bmatrix}$$

Finally, we define the following matrices:

In terms of the matrices M and F, the finite-difference system 1.16 can be put in the following form:

$$M\phi = \frac{1}{k} F\phi. \tag{2.6}$$

As we have noted (see p. 14), the assumption of homogeneous boundary conditions enables us to approximate the continuous adjoint by simply solving the adjoint equations of Eq. 2.6. Thus we define ϕ^* by the following equation:

$$M^*\phi^* = \frac{1}{k} F^*\phi^*,$$
 (2.7)

where the * denotes the operation of taking the transpose. Forming the transpose of M involves taking the transpose of each of its blocks, and, in particular, would involve the transpose of the matrice J_i . In the great majority of cases, J_i is block lower triangular, and the iterative algorithms are designed to exploit this fact. However, J_i^* will then be block upper triangular, which is undesirable if we wish to use exactly the same algorithms in both cases. To circumvent this difficulty, we form the adjoint matrices and then invert or reverse the order of the energy groups. Hence we define the following matrices:

a.
$$\overline{J}_{ig} = J_{i,G-g+1};$$

b. $\overline{K}_{ig} = K_{i,G-g+1};$
c. $\overline{B}_{i}^{gg'} = B_{i}^{G-g'+1,G-g+1};$
d. $\overline{F}_{i}^{gg'} = F_{i}^{G-g'+1,G-g+1}.$

Then with \overline{J}_i , \overline{F}_i , \overline{K}_i , \overline{M} , and \overline{F} defined in terms of these matrices as before, we can solve the equation

$$\overline{\mathbf{M}}\psi = \frac{1}{\mathbf{k}} \overline{\mathbf{F}}\psi, \tag{2.9}$$

which has the same form as the real equation.

We obtain the adjoint flux ϕ^* by setting

$$\phi_{i,g,j}^* = \psi_{i,G-g+1,j}. \tag{2.10}$$

III. SEARCH STRATEGY AND α-CALCULATIONS

Consider a critical system; at time t=0, let the system be modified in some manner. Suppose that this modification or change is a constant in time; then if the matrices M and F reflect the conditions of the modified system, the flux $\Phi(x,t)$ can be described by the equation

$$V^{-1}\frac{\partial\Phi}{\partial t} = (F - M)\Phi, \qquad (3.1)$$

where V is a diagonal matrix whose elements are the speeds associated with the energy groups. After a sufficient length of time, the flux can be reasonably described by a solution of Eq. 3.1 separable in time. Letting $\Phi(\mathbf{x},t) = \phi(\mathbf{x})e^{\alpha t}$ in 3.1 leads to

$$(\mathbf{M} + \alpha \mathbf{V}^{-1}) \phi = \mathbf{F} \phi. \tag{3.2}$$

Note that since αV^{-1} is a diagonal matrix, its effect is to change the diagonal elements of M from e_{igj} to $e_{igj} + \alpha v_g^{-1} V_{ij}$. Thus the form of Eqs. 3.2 is identical to that of Eq. 2.6 with k=1. Thus if we call the result of solving 2.6 a k-calculation, we see that a means for solving 3.2 can be based on a series of k-calculations. Thus, for each given α , we solve the system

$$(M + \alpha V^{-1}) \phi(\alpha) = \frac{1}{k(\alpha)} F\phi(\alpha). \tag{3.3}$$

Then we search for that value of α for which $k(\alpha) = 1$. Thus we see that an α -calculation is just a particular type of criticality search.

To discuss the strategy employed in a search calculation, we shall have to describe briefly the method used in solving Eq. 2.6. Now recall that

$$e_{igj} = a_{igj} + b_{igj} + c_{igj} + d_{igj} + V_{ij} \left[\left(\sigma_R^g \right)_{ij} + D_{ijg} B^2 \right],$$

where for each (i,j), we have

$$\sigma_{\rm R}^{\rm g} \; = \; \sigma_{\rm c}^{\rm g} \; + \; \sigma_{\rm f}^{\rm g} \; + \; \sum_{\rm g' \neq \rm g} \; \sigma^{\rm gg'}. \label{eq:sigma_g}$$

Thus,

$$e_{igj} > b_{igj} + d_{ijg} = b_{igj} + b_{igj-1};$$

hence, J_{ig}^{-1} exists and is positive. Moreover,

$$e_{igj} > V_{ij} \sum_{g' \neq g} \sigma_{ij}^{gg'};$$

hence, J_i^{-1} exists and is nonnegative. As a consequence, the matrix $M^{-1}F$ exists and is nonnegative. The problem posed by Eq. 2.6 can be viewed in the following manner: Find the largest positive eigenvalue and the corresponding positive eigenvector of the matrix $M^{-1}F$. The method used to find these quantities is that of power iterations. The method consists of the following iterations: Starting with any nonnegative vector $\phi^{(0)}$ and any positive number $k^{(0)}$, the sequences $\psi^{(n)}$ and $k^{(n)}$ are generated as follows:

$$\psi(n) = \frac{1}{k(n)} (M^{-1}F) \psi(n-1), \quad k(n) = k(n-1) \frac{(\psi(n), \psi(n))}{(\psi(n), \psi(n-1))}, \quad (3.4)$$

where

$$(\psi,\phi) = \sum_{ijg} \psi_{igj}\phi_{igj}.$$

Then, if the matrix $M^{-1}F$ is nonnegative and irreducible, it is well known⁴ that the sequence $k^{(n)}$ will tend to the unique largest positive eigenvalue of the matrix, and the sequence $\psi^{(n)}$ will tend to the corresponding positive eigenvector. In the above situation, we see that the generation of $\psi^{(n)}$ from $\psi^{(n-1)}$ involves the solution of the equation

$$M\psi^{(n)} = \frac{1}{k^{(n-1)}} F\psi^{(n-1)}. \tag{3.5}$$

In general, the matrix M is far too large to be held in fast core; however, it is block tridiagonal with nonpositive off-diagonal blocks and nonsingular diagonal blocks, which are diagonally dominant and have nonnegative inverses. Thus, a Gauss-Seidel iterative procedure will converge to a solution of Eq. 3.5. Referring to Eq. 2.5a, we define

a.
$$D = diag[J_1, J_2, ..., J_1];$$

b. $L = \begin{bmatrix} O & & O \\ K_1 & O & & \\ & O_{K_2} & O_{\ddots} \\ & & K_{I-1}O \end{bmatrix}, \quad U = \begin{bmatrix} O & K_1 & & O \\ & O & K_2 & & \\ & & \ddots & & \\ & & & \ddots & \\ & & & & O \end{bmatrix}.$
(3.6)

Then M = D - L - U, and a Gauss-Seidel iteration relative to this splitting is defined by the following procedure:

(D-L)
$$X_n^{(\ell)} = UX_n^{(\ell-1)} + \frac{1}{k^{(n-1)}} F\psi^{(n-1)}, \quad \ell = 1, 2, ...,$$
 (3.7)

with $X_n^{(0)} = \psi^{(n-1)}$. The sequence $X_n^{(\ell)}$ will then converge to the solution $\psi^{(n)}$ of Eq. 3.5. Moreover, if there is no up-scattering, the matrices J_i can be inverted directly since they are then block lower triangular, with tridiagonal matrices forming the diagonal blocks. If up-scattering is present, the matrices J_i satisfy the necessary criteria so that a Gauss-Seidel iterative procedure can be applied to these matrices.

The present code does not go beyond the first iterate in the iterative procedure defined by Eq. 3.7. Thus, $\psi^{(n+1)}$ is defined to be equal to $\chi_n^{(1)}$; hence, the iterative procedure used in this code is as follows: Starting with any positive number $k^{(0)}$ and any nonnegative vector $\phi^{(0)}$, a sequence of numbers $k^{(n)}$ and vectors $\phi^{(n)}$ is generated by

$$(D-L) \phi^{(n)} = U\phi^{(n-1)} + \frac{1}{k^{(n-1)}} F\phi^{(n-1)}, k^{(n)} = k^{(n-1)} \frac{(\phi^{(n)}, \phi^{(n)})}{(\phi^{(n)}, \phi^{(n-1)})}.$$
(3.8)

Thus we see that the actual iterative procedure used in this code is a non-stationary, nonhomogeneous power iteration, where the product

$$\prod_{\ell=1}^{n} S_{\ell}$$

replaces

$$\frac{1}{\prod_{\ell=1}^{n} k^{(\ell)}} (M^{-1}F).$$

Here S_{ℓ} denotes the nonnegative matrix $(D-L)^{-1}[U+(1/k)(\ell)F]$.

Before we discuss the effect of the above iterative procedure on the search strategy, we shall display the above iterations in more detail. Thus, using the definitions of D, L, and U, we have the sweep through the channels.

$$J_{i}\phi_{i}^{(n)} = K_{i-1}\phi_{i-1}^{(n)} + K_{i}\phi_{i+1}^{(n-1)} + \frac{1}{k^{(n)}} F_{i}\phi_{i}^{(n-1)} \text{ for } i = 1, 2, ..., I \quad (3.9)$$

with $K_0 = K_{\bar{1}} = 0$.

If up-scattering is present, the matrices J_i will not be block lower triangular; thus, the vector $\phi_i^{(n)}$ will be found iteratively. Within a given channel, we have the following iterative scheme:

$$J_{ig}Y_{ig}^{(m)} = \sum_{g' < g} B_{i}^{g'g}Y_{ig'}^{(m)} + \sum_{g > g'} B_{i}^{g'g}Y_{ig'}^{(m-1)} + K_{i-1,g}\phi_{i-1,g}^{(n)}$$

$$+ K_{ig}\phi_{i+1,g}^{(n-1)} + \frac{1}{k(n)} \sum_{g'=1}^{\hat{G}} F_{i}^{g'g}\phi_{ig'}^{(n-1)} \text{ for } g = 1, 2, ..., G.$$
(3.10)

Here we take $Y_{ig}^{(0)} = \phi_{ig}^{(n-1)}$, and upon achieving convergence, we set $\phi_{ig}^{(n)} = Y_{ig}^{(m)}$ for g = 1, 2, ..., G, and then move on to the next channel. If no up-scattering is present, we obtain $\phi_{ig}^{(n)}$ exactly in one sweep through the groups. Except for the full periodic case, the vectors $Y_{ig}^{(m)}$ are found by a direct inversion procedure based on the factorization of a tridiagonal matrix into the product of a lower bidiagonal and an upper bidiagonal matrix. This technique is sometimes known as Choleski's method. In the full periodic case, the solution is found by means of an iterative technique, which will be discussed in Section VI.

Having thus displayed the iteration procedure at various stages of detail, we shall consider the question of how a search is carried out. In a search, we seek the value of a parameter such that the resulting system has a k_{eff} equal to a prescribed value. Thus, we seek that value X_0 such that $k(X_0) = k_0$ where k_0 is given. This is accomplished by means of linear interpolation; hence, given

$$\left(X^{(s)},k(X^{(s)})\right)$$
 and $\left(X^{(s-1)},k(X^{(s-1)})\right)$,

a straight line is passed through these points, and $X^{(s+1)}$ is taken to be the abscissa of the intersection of this straight line with the line $k=k_0$. The major problem is that an estimate for k(X) is expensive in machine time; moreover, the cost is higher than usual in this code for the following reason. Experience has shown that generally the rate of convergence of a power iteration procedure based on Eq. 3.4 is relatively fast $(r \cong 0.9)$ for the flux and even faster for the sequence $k^{(n)}$. (In this context, by a rate of convergence r, we mean that the error eventually behaves like r^n , where n is the iteration count.) Since the matrix M is so large, we cannot perform an iterative procedure based on Eq. 3.4. On the other hand, we could base a practical iterative procedure on Eq. 3.5 in conjunction with Eq. 3.6. Then the same rate of convergence would apply to Eq. 3.5; however, the rate of convergence of Eq. 3.7 would be very much slower. Since the determination of the sequence $k^{(n)}$ is based on Eq. 3.5, we would

expect the rate of convergence of $k^{(n)}$ to be comparable to a procedure based on Eq. 3.4, provided a sufficient number of iterations are performed in Eq. 3.7. However, since the flux is determined by Eq. 3.7, we would expect its rate of convergence to be slow. Thus, insofar as rates are concerned, we would expect the sequence k(n) to converge more rapidly than the fluxes. The present code uses only one iteration of Eq. 3.7 per iteration of Eq. 3.5, thereby creating the iterative procedure, Eq. 3.8. In this procedure, the rate for the sequence k(n) is now comparable to that for the flux estimates, hence converging at a rate of 0.99 or greater for the average-sized problem (Appendix A, Sample Problem 1). This situation, namely the slow convergence of the estimates k(n), causes an extreme amount of difficulty in the search procedure by compounding or accentuating the following dilemma: On the one hand, for a given $X^{(i)}$, a sufficiently accurate estimate $k^{(n)}(X^{(i)})$ must be found for $k(X^{(i)})$ to ensure that the sequence $X^{(i)}$ will converge to X₀. On the other hand, because of the curvature of k(X), an accurate determination of $k(X^{(i)})$ does not lead to a comparably accurate estimate X(i+1) of X0 when X(i) is not near X0. Taking these considerations into account, we have developed a search procedure based on the following observations:

- 1. Each $k^{(n)}(X)$ is very costly; its value in time corresponds to a complete sweep of the mesh for all groups.
- 2. The sequence $k^{(n)}(X)$ is slowly converging, its rate being of the same order as the flux.
- 3. The sequence $k^{(n)}(X)$ changes slowly as a function of X. Thus, if $k^{(n)}(X)$ is an acceptable estimate for a given X and if X' is the next estimate for X_0 , then $k^{(1)}(X')$ is close to $k^{(n)}(X)$.
- 4. After a sufficient number of iterations, the sequence $k^{(n)}(X)$ converges to k(X) in a geometric fashion.

With these observations in mind, the following strategy is employed: Suppose that we are beginning the iterations needed to approximate $k_i = k(X^{(i)})$ corresponding to a control value $X^{(i)}$. Here $X^{(i)}$ may be an initial guess or the result of interpolation. The code then generates successive estimates $k_1^{(i)}$, $k_1^{(2)}$, ..., $k_1^{(n)}$ for k_1 . Because of observations 2, 3, and 4, no change is contemplated until at least 15 iterations have been completed and a minimal level of convergence has been achieved. We then determine whether the sequence $k_1^{(n)}$ is heading towards the desired value k_0 . If the sequence is heading away from k_0 , we stop the iteration as soon as the following criterion is satisfied:

$$|k_i^{(n)} - k_i| < 0.2 |k_i^{(n)} - k_0|.$$
 (3.11)

Of course we do not know k_i ; therefore $|k_i^{(n)} - k_i|$ must be estimated. Based on the assumption that the convergence of $k_i^{(n)}$ is geometric, this difference is approximated by

$$|k_i^{(n)} - k_i| \approx \frac{r_i^{(n)}}{1 - r_i^{(n)}} |\Delta k_i^{(n)}|,$$
 (3.12)

where

$$r_i^{(n)} = \frac{\left| \triangle k_i^{(n)} \right|}{\left| \triangle k_i^{(n-1)} \right|}$$

and

$$\triangle k_{i}^{(n)} = k_{i}^{(n)} - k_{i}^{(n-1)}.$$

The test defined by 3.11 in conjunction with the approximation 3.12 is made only after at least three successive unextrapolated iterations have occurred. This restriction is necessary because extrapolation destroys the geometric character of the convergence. The case in which the sequence $k_1^{(n)}$ is heading towards k_0 is somewhat more complicated. In this case, an effort is made as soon as possible to determine whether the iterates $k_1^{(n)}$ will pass through the desired value k_0 . Again this determination is made under the assumption that the sequence $k_1^{(n)}$ is converging geometrically; thus, the same restriction with regard to flux extrapolation is in effect. If the estimate $\overline{k}_1^{(n)}$ for k_1 lies on the other side of k_0 or is close to k_0 (that is, $|k_0 - \overline{k}_1^{(n)}| < (10) \epsilon_C$), then the code continues to iterate without making a control change. On the other hand, if the estimate $\overline{k}_1^{(n)}$ lies on the same side of k_0 as $k_1^{(n)}$ and is not close to k_0 , then the test defined by 3.11 and 3.12 is applied. If the test is satisfied, a control change is made using the estimate $\overline{k}_1^{(n)}$ for k_1 . If the test is not satisfied, then the code continues to iterate (see Appendix A, Sample Problem 3).

Having discussed the decision-making process involved in seeking a value of the control parameter (X) such that $k(X) = k_0$, we turn next to the logical control of the search process. Of course, what we ultimately seek as the result of a criticality search is not the required order of X but the value of some reactor parameter(s) such that k (parameter(s)) = k_0 ; i.e., we must have parameter(s) = F(X). The functional relationship used in the program is

$$P_n^{(s)} = P_n^{(0)} (1 + X^{(s)} \delta P_n),$$

where

X^S is the value of the control parameter for the sth pass,

 $P_n^{(0)}$ are the initial values of the reactor parameter(s),

P(s) are the values of the reactor parameter(s) for the sth pass,

and

 δP_n are the parameter(s) modifier(s).

The code currently allows the following four choices of the parameter(s) P_n (i.e., four different types of criticality searches are possible):

1. Composition Search

$$P_n = VF_{m,c}$$

where VFm,c are the volume fractions for materials m in composition c.

2. Dimension Search

$$P_n = \Delta X_i \text{ and/or } \Delta Y_j$$

where $\Delta X_i(\Delta Y_j)$ are the mesh increments (i.e., interval lengths) in the X(Y) direction.

3. Buckling (B²) Search

$$P_n = B_r^2$$
,

where $B_{\mathbf{r}}^2$ is the transverse buckling for region r.

4. α Calculation

$$P_n = \alpha$$
,

where α is the asymptotic inverse reactor time period.

The logical control of the criticality search is essentially as given in the CANDID2D document, but is duplicated here for completeness.

A first guess, X_A , and either a second guess, X_B , or an estimate of dk/dX, (KDOT), and k_0 are required input. If KDOT is given, then X_B is computed from the following equation:

$$X_B = X_A + (k_0 - k(X_A)) / KDOT.$$

Once $k(X_A)$ and $k(X_B)$ are computed, linear extrapolation is done until k_0 is bracketed. A maximum number of extrapolations to bracket k_0 is provided in the input. Once k_0 has been bracketed, the values of k(X) and X that best bracket k_0 are stored in X_L , k_L , X_R , and k_R such that $k_L < k_0 < k_R$.

For the purpose of the following discussion on bracketing k, let us assume that $\underline{X} < X_A < X_B < \overline{X}$. If an extrapolation results in $X < \underline{X}$, then a midpoint is attempted for X from $X = (\underline{X} + X_A)/2$. If a later extrapolation has the same result, $X < \underline{X}$, then the bound itself is used; i.e., $X = \underline{X}$.

A third failure at the same bound will be a terminal error. The above discussion follows through in a similar manner if $X_B < X_A$ and/or the failure is at the other bound, \overline{X} .

 X_1 , k_1 , X_2 , and k_2 are the interpolation parameters and are chosen in the following manner: Once k_0 has been bracketed, set X_1 and X_2 equal to X_L and X_R , and interpolate for X from the following equation:

$$X = X_1 + \frac{k_0 - k(X_1)}{k(X_2) - k(X_1)} (X_2 - X_1);$$

then k(X) is computed, and the bounds are updated as follows:

If
$$k(X) < k_0$$
, then $X_L = X$ and $k_L = k(X)$;

or

if
$$k(X) > k_0$$
, then $X_R = X$ and $k_R = k(X)$.

The closest two values of X are used for the next interpolation:

If
$$|X - X_1| < |X - X_2|$$
, then $X_2 = X$ and $k_2 = k(X)$;

or

if
$$|X - X_2| < |X - X_1|$$
, then $X_1 = X$ and $k_1 = k(X)$.

Now interpolate for a new value of X. If X does not fall in range of the best known bounds,

$$X < X_{\perp}$$
 and X_{R} ,

or

$$X > X_L$$
 and X_R ,

then the bounds are used:

$$X_1 = X_L$$
, $k_1 = k_L$,

or

$$X_2 = X_R, k_2 = k_R;$$

this will yield a valid X.

This procedure is continued until either the convergence criterion is met or a specifiable number of interpolations is exceeded.

IV. SOURCE CALCULATION

A source calculation can arise in two ways. Either a fixed source is prescribed, or an inhomogeneous boundary condition is present. In either event, the basic assumption of a source calculation is that the reactor system in the absence of the fixed source is subcritical. In terms of the matrices introduced in Section II, a source problem can be written in the form

$$(M - F) \phi = f, \tag{4.1}$$

where f is a prescribed nonnegative source. The assumption of subcriticality ensures that the matrix M - F is not singular. In addition, it ensures that the solution ϕ is nonnegative when f is nonnegative. This last statement is by virtue of the fact that $\phi = (1 - M^{-1}F)^{-1} M^{-1}f$, and by the hypothesis that the spectral radius of $M^{-1}F$ is less than one. Thus $(1 - M^{-1}F)^{-1} \ge 0$.

The method used to solve Eq. 4.1 takes advantage of the existing method for solving the homogeneous problem. Hence, in terms of the matrices introduced in Section III, the iterative procedure employed in the absence of up-scattering can be written in the form

$$(D-L) \phi(n) = (U+F) \phi^{(n-1)} + f.$$
 (4.2)

As far as the mechanics are concerned, this procedure differs from the k_{eff} case only to the extent that the eigenvalue estimates $k^{(n)}$ are not used in the fission source. Recalling the structure of D - L from Section III for no up-scattering, we see that $(D-L)^{-1}$ exists and is nonnegative. Furthermore, U + F is a nonnegative matrix; thus, the splitting M - F = (D-L) - (U+F) is a regular splitting.³ As a consequence, the spectral radius of $(D-L)^{-1}(U+F)$ is less than one; thus, 4.2 is a convergent iterative procedure.

When up-scattering is present, we shall assume that a fixed number of up-scattering iterations are performed. From Eq. 3.10, we recall that within each channel i, $1 \le i \le I$, the up-scattering iterations have the following form:

$$J_{ig}Y_{ig}^{(m)} = \sum_{g' < g} B_{i}^{g'g}Y_{ig'}^{(m)} + \sum_{g' > g} B_{i}^{g'g}Y_{ig'}^{(m-1)} + K_{i-1,g}\phi_{i-1,g}^{(n)}$$

$$+ K_{ig}\phi_{i+1,g}^{(n-1)} + \sum_{g'=1}^{G} F_{i}^{g'g}\phi_{ig'}^{(n-1)} + f_{ig}, \quad g = 1, 2, ..., G.$$

$$(4.3)$$

Here,

$$Y_{ig}^{(0)} \equiv \phi_{ig}^{(n-1)},$$

and at the end of ℓ iterations we set

$$\phi_{ig}^{(n)} \equiv Y_{ig}^{(\ell)}$$
.

Referring to the definition of J_i in Eq. 2.4a, we form the block decomposition $J_i = D_i - L_i - U_i$, where

$$D_{i} = diag[J_{i1}, J_{i2}, ..., J_{iG}],$$

and L_i and U_i are the remaining block lower and upper triangular matrices, respectively. In terms of this decomposition, the up-scattering iterations take the form

$$(D_{i} - L_{i}) Y_{i}^{(m)} = U_{i} Y_{i}^{(m-1)} + K_{i-1} \phi_{i-1}^{(n)} + K_{i} \phi_{i+1}^{(n-1)} + F_{i} \phi_{i}^{(n-1)} + f_{i}.$$
 (4.4)

If we set $R_i = (D_i - L_i)^{-1} U_i$, and recall that $Y_i^{(0)} \equiv \phi_i^{(n-1)}$ and $\phi_i^{(n)} \equiv Y_i^{(\ell)}$, then $\phi_i^{(n)}$ has the following form:

$$\phi_{i}^{(n)} = R_{i}^{\ell} \phi_{i}^{(n-1)} + (I - R_{i}^{\ell}) J_{i}^{-1} \left\{ K_{i-1} \phi_{i-1}^{(n)} + K_{i} \phi_{i+1}^{(n-1)} + F_{i} \phi_{i}^{(n-1)} + f_{i} \right\}. \tag{4.5}$$

We can write this iterative procedure in a manner analogous to Eq. 4.2 if we define the following matrices. Let \hat{M} and \hat{N} be defined by

$$\hat{M} = \begin{bmatrix}
I \\
-(I - R_2^{\ell}) J_2^{-1} K_1 & I \\
-(I - R_3^{\ell}) J_3^{-1} K_2 & I \\
-(I - R_1^{\ell}) J_1^{-1} K_{I-1} & I
\end{bmatrix} (4.6)$$

and

$$\hat{N} = \begin{bmatrix}
R_{1}^{\ell} + (I - R_{1}^{\ell}) J_{1}^{-1}F_{1} & (I - R_{1}^{\ell}) J_{1}^{-1}K_{1} \\
R_{2}^{\ell} + (I - R_{2}^{\ell}) J_{2}^{-1}F_{2} & (I - R_{2}^{\ell}) J_{2}^{-1}K_{2} \\
R_{1}^{\ell} + (I - R_{1}^{\ell}) J_{1}^{-1} + F_{1}
\end{bmatrix}$$
(4.7)

With these matrices defined, the iterative procedure 4.5 can be written in the form

$${\stackrel{\wedge}{M}}\phi^{(n)} = {\stackrel{\wedge}{N}}\phi^{(n-1)} + g, \tag{4.8}$$

where

$$g_i = (I - R_i^{\ell}) J_i^{-1} f_i \text{ for } i = 1, 2, ..., I.$$

At this point we can make several observations:

- 1. For each channel i, the matrix $\left(I R_i^{\ell}\right)J_i^{-1}$ is nonnegative and nonsingular.
- 2. The iterative procedure 4.8 is convergent.
- 3. The solution thus obtained satisfies Eq. 4.1.

Starting with the first observation, recall that $R_i = (D_i - L_i)^{-1} \ U_i$ is the block Gauss-Seidel iteration matrix associated with the matrix J_i . The matrix J_i has on its block diagonal nonsingular matrices which have positive inverses. Moreover, the block off-diagonal matrices are nonpositive; thus, R_i is nonnegative and has a spectral radius less than one. For the same reasons, J_i^{-1} and $(D_i - L_i)^{-1}$ are nonnegative. As a consequence, we see that $J_i^{-1} \ge R_i J_i^{-1}$ since

$$(I - R_i) J_i^{-1} = (I - R_i)(I - R_i)^{-1}(D_i - L_i)^{-1} = (D_i - L_i)^{-1} \ge 0.$$
 (4.9)

From this it follows that $\left(I-R_i^{\ell}\right)J_i^{-1}\geq 0$ for any $\ell\geq 1$. Since the spectral radius of R_i is less than one, the same is true of R_i^{ℓ} ; thus, $\left(I-R_i^{\ell}\right)^{-1}$ exists and is nonnegative. Hence, $\left(I-R_i^{\ell}\right)J_i^{-1}$ is not singular.

Addressing ourselves to the second observation, we see from the first observation that \hat{M}^{-1} exists and is nonnegative; moreover \hat{N} is also nonnegative. Thus the matrices \hat{M} and \hat{N} used in 4.8 form a regular splitting³ of the matrix \hat{M} - \hat{N} ; consequently, the process is convergent.

Concerning the third observation, suppose that ψ satisfies

$$(\stackrel{\wedge}{M} - \stackrel{\wedge}{N}) \psi = g. \tag{4.10}$$

Now $\stackrel{\wedge}{M}$ - $\stackrel{\wedge}{N}$ = SJ⁻¹(M - N), where

$$J = \operatorname{diag}[J_1, J_2, ..., J_I]$$
and
$$S = \operatorname{diag}\left[\left(I - R_1^{\ell}\right), ..., \left(I - R_I^{\ell}\right)\right].$$
(4.11)

From the first observation, the matrix SJ^{-1} is not singular; moreover, $g = SJ^{-1}f$. Hence, ψ satisfies Eq. 4.1; therefore $\psi = \phi$, where ϕ satisfies the equation

$$(M - F) \phi = f. \tag{4.12}$$

Based on the above observations, we can draw the following conclusion: When up-scattering is present, the solution does not depend on the number of up-scattering iterations performed. The use of up-scattering iterations would then rest on the possibility that they will improve the rate of convergence or at least the asymptotic rate. The possibility of improvement rests heavily on the behavior of \hat{N} defined in 4.7 as a function of the number of up-scattering iterations, ℓ . However, although the asymptotic rate of the process defined by Eq. 4.8 does increase with increasing ℓ , it is still not sufficient to offset the consequent increase in computation time.

V. ACCELERATION

In this section we shall describe the acceleration scheme that was implemented for source calculations. The scheme used was two-term Chebyshev extrapolation applied to the iterative scheme 4.2 (or 4.8 for up-scattering). This scheme was chosen because a two-term Chebyshev procedure is used to accelerate the $k_{\mbox{eff}}$ calculation. Thus a minimum amount of effort was involved in adapting this procedure to a source calculation.

Recall from Section IV that the iterative procedure used to solve a source problem can be written in the form

$$\phi^{(n+1)} = R\phi^{(n)} + b. \tag{5.1}$$

Here $R = (D-L)^{-1}(U+F)$ and $b = (D-L)^{-1}$ f, if no up-scattering is present. If up-scattering is present and we assume that a fixed number of up-scattering iterations are performed, then $R = \hat{M}^{-1}\hat{N}$ and $b = \hat{M}^{-1}f$, where \hat{M} and \hat{N} are defined by Eqs. 4.6 and 4.7. Neither the matrix R nor its associated Jacobi matrix is symmetric; thus, to justify the use of Chebyshev extrapolation, we shall have to assume that the matrix R has real eigenvalues. The two-term Chebyshev extrapolation as the described briefly as follows: Let ϕ be the solution of Eq. 5.1, and let $e^{(n)} = \phi - \phi^{(n)}$. It then follows that

$$e(p) = Re^{(p-1)} = Rpe^{(0)}.$$
 (5.2)

If we assume for simplicity that the eigenvalues λ_q of R are simple with $1>\lambda_1>\lambda_2>\ldots\geq 0$, then $e^{\left(p\right)}$ can be expressed in the form

$$e(p) = e_1 \lambda_1^p u_1 + c_2 \lambda_2^p u_2 + \dots$$
 (5.3)

If, on the other hand, we had used

$$\iint_{j=1}^{\mathbf{p}} \frac{\mathbf{R} - \alpha_{j}}{1 - \alpha_{j}} = \frac{\mathbf{R} - \alpha_{\mathbf{p}}}{1 - \alpha_{\mathbf{p}}} \cdot \frac{\mathbf{R} - \alpha_{\mathbf{p}-1}}{1 - \alpha_{\mathbf{p}-1}} \cdot \frac{\mathbf{R} - \alpha_{1}}{1 - \alpha_{1}}$$

in place of R^p on $e^{(0)}$, we would have obtained

$$e^{(p)} = c_1 \prod_{j=1}^{p} \frac{\lambda_1 - \alpha_j}{1 - \alpha_j} u_1 + c_2 \prod_{j=1}^{p} \frac{\lambda_2 - \alpha_j}{1 - \alpha_j} u_2 + \dots$$
 (5.4)

Here we can see the possibility of making $e^{(p)}$ as defined by Eq. 5.4 smaller than $e^{(p)}$ as defined by Eq. 5.3. For if we set

$$Q_{p}(\lambda;\lambda_{1}) = \prod_{j=1}^{p} \frac{\lambda - \alpha_{j}}{1 - \alpha_{j}}, \qquad (5.5)$$

and then choose $Q_p(\lambda;\lambda_1)$ to be the minimal polynomial in the supremum norm of degree p over the interval $[0,\lambda_1]$, then $e^{(p)}$ as defined by Eq. 5.4 will be smaller than $e^{(p)}$ as defined by Eq. 5.3. The solution to the above minimization problem is given by

$$Q_{p}(\lambda; \lambda_{1}) = \frac{C_{p}\left(\frac{2\lambda}{\lambda_{1}} - 1\right)}{C_{p}\left(\frac{2}{\lambda_{1}} - 1\right)},$$
(5.6)

where $C_p(x)$ is the Chebyshev polynomial of degree p normalized to be 1 at x = 1. The roots of the polynomial in Eq. 5.6 are then given by

$$\alpha_{j}^{(p)} = \frac{\lambda_{1}}{2} \left[1 + \cos \left\{ (2(p-j) + 1) \frac{\pi}{2p} \right\} \right] \quad \text{for } j = 1, 2, ..., p.$$
 (5.7)

Since the matrix polynomial $Q_p(R;\lambda_I)$ is expressed in its factorial form, we accomplish the extrapolation by applying one factor at a time. Suppose that starting with a flux $\phi(n)$ and an estimate $\overline{\lambda}_1$ for λ_1 , we decide to perform an extrapolation cycle of order p. The roots $\{\alpha(p)\}_{j=1}^p$ are then computed from

Eq. 5.7 by using $\overline{\lambda}_1$ in place of λ_1 . For use in the actual application, we define the numbers $\beta_j^{(p)} = \left(1 - \alpha_j^{(p)}\right)^{-1}$. Then we generate in succession

$$\phi^{(n+1)} = R\phi^{(n)} + b;$$

$$\hat{\phi}^{(n+1)} = \frac{1}{1 - \alpha_1^{(p)}} \left(\phi^{(n+1)} - \alpha_1^{(p)} \phi^{(n)} \right) = \beta_1^{(p)} \phi^{(n+1)} + \left(1 - \beta_1^{(p)} \right) \phi^{(n)};$$

$$\phi^{(n+2)} = R\hat{\phi}^{(n+1)} + b;$$

$$\hat{\phi}^{(n+2)} = \frac{1}{1 - \alpha_2^{(p)}} \left(\phi^{(n+2)} - \alpha_2^{(p)} \hat{\phi}^{(n+1)} \right) = \beta_2^{(p)} \phi^{(n+2)} + \left(1 - \beta_2^{(p)} \right) \hat{\phi}^{(n+1)};$$

$$\vdots$$

$$\hat{\phi}^{(n+p)} = \frac{1}{1 - \alpha_p^{(p)}} \left(\phi^{(n+p)} - \alpha_p^{(p)} \hat{\phi}^{(n+p-1)} \right) = \beta_p^{(p)} \phi^{(n+p)} + \left(1 - \beta_p^{(p)} \right) \hat{\phi}^{(n+p-1)}.$$

With the generation of $\hat{\phi}(n+p)$ we have completed an extrapolation cycle of order p. The error vector, $\hat{e}(n+p) = \phi - \hat{\phi}(n+p)$, then satisfies

$$\hat{e}^{(n+p)} = Q_p(R; \overline{\lambda}_1) e^{(n)}. \tag{5.9}$$

As is well known,^{3,4} the effectiveness of this extrapolation is sensitive to the accuracy of the estimate $\overline{\lambda}_1$ for λ_1 . In this code, the initial estimate for λ_1 is found from the fact that, for sufficiently large n,

$$\frac{||\phi^{(n+1)} - \phi^{(n)}||_{2}}{||\phi^{(n)} - \phi^{(n-1)}||_{2}} \cong \lambda_{1}, \tag{5.10}$$

where

$$||\phi||_2 = \left(\sum_{ijg} \phi_{ijg}^2\right)^{r/2}$$
.

Subsequent estimates are made using an up-dating procedure introduced by Varga.³ This up-dating procedure is used in the following way: Suppose an extrapolation cycle of order p was begun with the nth iteration, so that $\hat{\phi}(n+p)$ is the result of this cycle. An unextrapolated pass is then made generating $\phi(n+p+1)$, and the following quantity is formed:

$$E_{n,p} = \frac{||\phi^{(n+p+1)} - \hat{\phi}^{(n+p)}||_{2}}{||\phi^{(n+1)} - \phi^{(n)}||_{2}}.$$
 (5.11)

It can be shown³ that

$$\mathbf{E}_{\mathbf{n},\mathbf{p}} \cong \left[Q_{\mathbf{p}}(\lambda_1; \overline{\lambda}_1) \right]. \tag{5.12}$$

Now

$$|Q_{p}(\lambda; \overline{\lambda}_{1})| \leq Q_{p}(\overline{\lambda}_{1}; \overline{\lambda}_{1}) = \frac{1}{C_{p}(\frac{2}{\overline{\lambda}_{1}} - 1)} \quad \text{for } 0 \leq \lambda \leq \lambda_{1}, \tag{5.13}$$

and

$$Q_{p}(\lambda; \overline{\lambda}_{1}) \ge Q_{p}(\lambda_{1}; \overline{\lambda}_{1}) \quad \text{for } \overline{\lambda}_{1} \le \lambda \le 1.$$
 (5.14)

Hence, using Eq. 5.12, if

$$E_{n,p} \leq \frac{1}{C_p \left(\frac{2}{\overline{\lambda}_1} - 1\right)},\tag{5.15}$$

then we can conclude that $\lambda_1 \leq \overline{\lambda}_1$ and also that the extrapolation was as effective as we can expect since $1/[C_p(2/\overline{\lambda}_1-1)]$ is the maximum of the modulus of $Q_p(\lambda;\lambda_1)$ in the interval $[0,\overline{\lambda}_1]$. On the other hand, if

$$E_{n,p} > \frac{1}{C_p\left(\frac{2}{\lambda_1} - 1\right)}, \tag{5.16}$$

then, using Eq. 5.14, we can conclude that $\overline{\lambda}_1 \leq \lambda_1$. In this case, the estimate for $\overline{\lambda}_1$ will need to be moved up. Since $\overline{\lambda}_1 < \lambda_1$, we know that $Q_p(\lambda_1; \overline{\lambda}_1)$ is positive; thus, using Eq. 5.12, we define the new estimate $\overline{\lambda}_1'$ as that value of λ for which

$$E_{n,p} = Q_p(\overline{\lambda}_1^!; \overline{\lambda}_1). \tag{5.17}$$

This yields a new estimate $\overline{\lambda}_1'$ for which $\overline{\lambda}_1 < \overline{\lambda}_1' < 1$.

A source calculation then proceeds in the following manner: Starting with an initial guess for the flux, we perform 15 unextrapolated iterations, at which point we make an initial estimate $\overline{\lambda}_1$ for λ_1 using Eq. 5.10. If this estimate, which is also a measure of the rate of convergence, is greater than 0.75 and less than 0.995, then we perform an extrapolation cycle of order 6. If the estimate is greater than 0.995 but not greater than 0.999, the order is limited to 5. On the other hand, if the estimate exceeds 0.999, then this value is used in place of the estimate. On the other hand, if the estimate is less than 0.75, then the unextrapolated iterations are continued until the estimates exceed 0.75 or convergence is attained. At the end of an extrapolation cycle, we make a test according to inequality 5.15 as to whether the cycle was effective. If it was not, we make a new estimate according to Eq. 5.17. Three unextrapolated iterations are then performed before entering the next extrapolated cycle. Generally speaking then, the pattern is extrapolated cycles of length p, interspersed with three unextrapolated iterations. Of course, there are exceptions to this general pattern. For example, if the rate of convergence (as estimated by Eq. 5.10) drops below 0.75 or rises above one during unextrapolated iterations, then these unextrapolated iterations are continued.

The problem is continued until the following criteria are satisfied:

$$||\phi^{(n)}||_{2}^{-1}||\phi^{(n)} - \phi^{(n-1)}||_{2} \le \epsilon_{\phi}$$
 (5.18)

and

$$\overline{k}(n) - \underline{k}(n) \le \epsilon_k,$$
 (5.19)

where

$$\overline{k}(n) = \max_{ijg} \frac{\phi_{ijg}^{(n)}}{\phi_{ijg}^{(n-1)}}, \quad \underline{k}(n) = \min_{ijg} \frac{\phi_{ijg}^{(n)}}{\phi_{ijg}^{(n-1)}}, \quad (5.20)$$

and the extrema are taken over the set of points (i,j,g) for which $\phi_{ijg}^{(n-1)} \geq \underline{\epsilon_{\phi}}$ (see Appendix A, Sample Problem 4). Also note the following fact: If any component of the flux should become negative during, or at the conclusion of, an extrapolation cycle, this component is set equal to zero. The intuitive justification of this is as follows: The external source is nonnegative; thus it might be expected to be strongly biased in the direction of the dominant eigenvector of R, which is positive. Thus the desired solution is positive and thereby biased in the direction of the dominant eigenvector. Thus, if a component becomes negative in the course of extrapolation, it indicates a buildup of a lower harmonic. This means the approximation is heading away from the solution and will have to pass through zero on its way toward the solution. Thus, setting the negative component equal to zero can do no harm and indeed prevents the difficulty that arises when any component of the flux goes negative.

Many of the procedures used to accelerate a source calculation have been incorporated into the $k_{\mbox{eff}}$ acceleration procedures. To describe these procedures, as well as others that were developed, we shall describe the acceleration procedure in some detail. For simplicity in description, we shall assume that no up-scattering is present and that the problem is not a full periodic problem.

Recall from Section III that a k_{eff} calculation involves finding the largest positive eigenvalue and corresponding eigenvector of a nonnegative matrix. We saw from Eq. 3.8 that this code solves the problem by relying entirely on power iterations (disregarding up-scattering and periodic iterations). Thus, from Eq. 3.8, the basic iterative procedure is given by

$$\phi^{(n)} = (D - L)^{-1} \left[U + \frac{1}{k^{(n-1)}} F \right] \phi^{(n-1)}, \quad k^{(n)} = k^{(n-1)} \frac{\left(\phi^{(n)}, \phi^{(n)}\right)}{\left(\phi^{(n)}, \phi^{(n-1)}\right)}. \quad (5.21)$$

If we set $S(k) = (D - L)^{-1}[U + (1/k)F]$ with $S_{n-1} = S(k^{(n-1)})$ we can write Eq. 5.21 in the form

$$\phi^{(n)} = S_n \phi^{(n-1)} = \prod_{m=1}^{n} S_m \phi^{(0)}, \quad k^{(n)} = k^{(n-1)} \frac{(\phi^{(n)}, \phi^{(n)})}{(\phi^{(n)}, \phi^{(n-1)})}.$$
 (5.22)

It is not hard to see that if k is the dominant eigenvalue and ϕ the corresponding eigenvector which satisfy

$$M\phi = \frac{1}{k} F\phi, \qquad (5.23)$$

then

$$S(k) \phi = \phi. \tag{5.24}$$

Moreover, 1 is the dominant eigenvalue of the nonnegative matrix S(k). Thus if we know that the sequence $k^{(n)}$ as generated in Eq. 5.22 converges to k, it follows that the sequence $\phi^{(n)}$ will converge to ϕ . However, the coupling between $k^{(n)}$ and $\phi^{(n)}$ is so complex that it is difficult to show that the procedure defined by Eq. 5.22 will converge.

Let us assume that the process defined by Eq. 5.22 is convergent. The application of Chebyshev acceleration to power iteration is predicated on two basic assumptions:

- 1. The eigenvalues of the matrix are real.
- 2. The eigenvalue estimates are converging at a faster rate than the eigenvector iterates. Moreover, the extrapolation is not applied until the eigenvalue iterates have almost converged.

Suppose both assumptions are satisfied; so that $k(n) \cong k$. Then the process becomes an iterative process with a fixed matrix, S = S(k). Thus,

$$\phi(\ell+1) = S\phi(\ell) \quad \text{for } \ell \geq n, \tag{5.25}$$

where S has its dominant eigenvalue equal to one. Let $\overline{\sigma}$ be an estimate for the second largest eigenvalue σ of S, whose eigenvalues lie in the interval [0,1]. Then a Chebyshev extrapolation cycle of order p applied to the process Eq. 5.25, beginning with $\ell = n$, would yield

$$\hat{\phi}^{(n+p)} = Q_p(S; \overline{\sigma}) \phi^{(n)}, \qquad (5.26)$$

where

$$Q_{p}(x; \overline{\sigma}) = \frac{C_{p}(\frac{2x}{\overline{\sigma}} - 1)}{C_{p}(\frac{2}{\sigma} - 1)}.$$
 (5.27)

The function $C_p(x)$ is, as before, the Chebyshev polynomial of degree p normalized to be 1 at x=1. If we expand $\phi^{(n)}$ in terms of the eigenvectors of S, we see that the $\hat{\phi}^{(n+p)}$ has the form

$$\hat{\phi}^{(n+p)} = a_1 Q_p(1; \overline{\sigma}) u_1 + a_2 Q_p(\sigma_1; \overline{\sigma}) u_2 + ...,$$
 (5.28)

where $\phi^{(n)} = c_1 u_1 + c_2 u_2 + \ldots$ and $1 > \sigma_1 \ge \sigma_2 \ge \ldots \ge 0$ are the eigenvalues of S. Now $Q_p(1; \overline{\sigma}) \equiv 1$ and $Q_p(x; \overline{\sigma})$ is minimal over $[0, \overline{\sigma}]$. Therefore we have performed an effective extrapolation depending on where our estimate $\overline{\sigma}$ is relative to σ_1 .

We have described the standard technique for applying Chebyshev extrapolation to eigenvalue problems. Accounts may be found in Refs. 3 and 6. Note that here the dominant eigenvector estimates play the same role that the error vectors did in the source calculation. Based on the above analysis, we could use the up-dating method for finding new estimates for σ_1 .

When the two-term Chebyshev extrapolation is applied to a k_{eff} problem, the results are disappointing (see Appendix A, Sample Problems 1 and 2). This is in contrast to the source-calculation case (see Appendix A, Sample Problems 4 and 5). There the results are quite good; for example, the extrapolation cycles make their theoretical error-reduction criterion $(E_{n,p} \leq C_p(2/\overline{\lambda}-1)^{-1})$ more often than not. On the other hand, for a k_{eff} calculation, an extrapolation cycle will almost never make its theoretical criterion, frequently does little better than the same number of power iterations would have done, and occasionally even causes apparent divergence. Of course, we can always recover from this latter contingency by performing power iterations.

We shall explain why extrapolations in the $k_{\mbox{eff}}$ case behave as they do. We shall also describe what remedial steps were taken in the use of extrapolation to minimize the effects described above.

The problem in the k_{eff} case is that assumption 2 is not satisfied. That assumption 1 is essentially satisfied is shown by the following empirical evidence: If k=1, then the matrix S(k) is essentially the same as the matrix R which appeared in the source calculation. In fact, $S(k)=R+(1/k-1)(D-L)^{-1}$ F. The source extrapolations were very effective, indicating that R had essentially real eigenvalues. Thus we would expect that at least for k near one, S(k) would also have essentially real eigenvalues. This argument is somewhat weak since S(k) can be viewed as a perturbation of R by a nonsymmetric matrix.

On the other hand, the evidence that assumption 2 is not satisfied is overwhelming. In almost every case, the rate of convergence of the eigenvalue iterates $k^{(n)}$ is of the same order as the eigenvalue iterates. Thus, it is not feasible to wait until the sequence $k^{(n)}$ has converged before extrapolating. Let us consider the effect of performing an extrapolation cycle before the eigenvalue iterates have converged. Two effects are present. First, the estimate $k^{(n)}$ for k is not accurate, and second, the iterates $k^{(n)}$ are changing during the extrapolation cycle. Consider the first effect. Assume that we have an estimate $\mu = k^{(n)}$ and an estimate $\overline{\sigma}$ for the dominance ratio of $S(\mu) = S(k^{(n)})$; i.e., if $\lambda_1(\mu) > \lambda_2(\mu) \ge \dots$ are the eigenvalues of $S(\mu)$, then we have an estimate for $\lambda_2(\mu)/\lambda_1(\mu)$. Note that since $\mu \ne k$, $\lambda_1(\mu) \ne 1 = \lambda_1(k)$. Let $\phi^{(n)}$ be considered in terms of the eigenvectors of $S/(\mu)$; i.e.,

$$\phi(n) = c_1 u_1 + c_2 u_2 + ..., \quad u_i = u_i(\mu).$$
 (5.29)

Then, as a result of an extrapolation cycle, we obtain

$$\hat{\phi}(n+p) = c_1 Q_p(\lambda_1(\mu); \overline{\sigma}) u_1 + c_2 Q_p(\lambda_2(\mu); \overline{\sigma}) u_2 + \dots$$
 (5.30)

Figure 3 is a graph of $Q_p(x; \overline{\sigma})$.

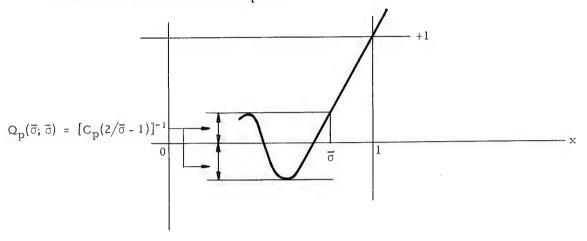


Fig. 3. Graph of $Q_D(x; \overline{\sigma})$

Now $\lambda_1(\mu) \neq 1$, and since μ is in general not even close to k, we see that $\lambda_1(\mu)$ need not be less than one; or if it is less than one, it need not be greater than $\overline{\sigma}$. Thus we see that the "reduction factor"

$$\left| \frac{Q_{\mathbf{p}}(\lambda_{2}(\mu); \, \overline{\sigma})}{Q_{\mathbf{p}}(\lambda_{1}(\mu); \, \overline{\sigma})} \right| \tag{5.31}$$

need not be less than the theoretical reduction factor $[C_p(2/\overline{\sigma}-1)]^{-1}$ regardless of how well the dominance ratio $\overline{\sigma}$ is estimated. Indeed we can see that for $\lambda_1(\mu)<\overline{\sigma}$, the reduction factor can be greater than one. Thus, we can begin to see why the extrapolations are so ineffective, and how the higher harmonics can be built up during a cycle when $\lambda_1(\mu)<\overline{\sigma}$.

The problem of a changing estimate $k^{(n)}$ during the extrapolation cycle adds an additional degree of complexity by preventing us from applying a Chebyshev polynomial.

Having examined the causes for the ineffectiveness of the extrapolation procedure, we find that two courses of action are open. First, change the basic iterative structure. Second, modify the technique of applying the extrapolation procedure. The first course of action requires scrapping the computational portion of the code and starting from the beginning. This is being done. As far as this code is concerned, the procedures for applying the extrapolation have been modified to at least prevent disaster.

The following procedure is available, on option, to help prevent the occurrence of $\lambda_1(\mu) < \overline{\sigma}$. The code has a built-in sequence of numbers 0.8, 0.85, 0.9, 0.925, 0.95, 0.96, ..., 0.99, 0.991, ..., 0.995. The user

specifies one of these numbers, say $\overline{\sigma}_0$. Then in the first extrapolation cycle, the code will not use any estimate for $\overline{\sigma}$ that exceeds $\overline{\sigma}_0$. If $\overline{\sigma}_1$ is the next larger number in the above sequence, then on the second extrapolation cycle the code will not use any estimate $\overline{\sigma}$ which exceeds $\overline{\sigma}_1$. This process is continued until 0.995 is reached. The above procedure then has the effect of restricting the growth of $\overline{\sigma}$ as the iterations proceed.

In this connection, a second technique is available for avoiding the situation $\lambda_1(\mu) > \overline{\sigma}$. This involves running a coarse-mesh problem and then refining the mesh and using the results from the coarse problem.

The up-dating method for finding estimates for the dominance ratio σ is used only when the sequence k(n) has essentially converged. As we have noted before, this seldom occurs before the flux has converged. The reason the up-dating method is not used before this point is that it forces the estimates for σ to be large very early in the iteration process. Most of the estimates for σ are made by means of

$$\overline{\sigma}_{n} = \frac{\|\phi^{(n)} - \phi^{(n-1)}\|_{2}}{\|\phi^{(n-1)} - \phi^{(n-2)}\|_{2}} = \sigma.$$
 (5.32)

This monitoring is performed during the three- to four-power iterations that follow each Chebyshev extrapolation cycle.

VI. SOLUTION OF THE FULL PERIODIC PROBLEM

The full periodic case arises in $r\theta$ geometry when the reactor domain is annular. Thus the flux must satisfy $\phi_g(r,0)\equiv\phi_g(r,2\pi)$. This condition leads to the following form for the matrices J_{ig} :

$$J_{ig} = \begin{bmatrix} e_{ig1} & -b_{ig1} & -b_{igJ} \\ -b_{ig1} & e_{ig2} & -b_{iJ2} \\ & & & & \\ & & -b_{ig,J-1} \\ -b_{igJ} & -b_{ig,J-1} & e_{ig,J} \end{bmatrix}. \tag{6.1}$$

We recall from Eq. 3.10 that for each channel i and group g we seek the solution of the equation

$$J_{ig}Y_{ig} = f_{ig}, (6.2)$$

where f_{ig} denotes the entire right-hand side of Eq. 3.10. In the nonfull periodic case, $b_{ig}J$ = 0, and the Choleski algorithm is used. Because of memory limitations and programming difficulties, Eq. 6.2 is solved iteratively.

To effectively describe the iterative scheme used to solve Eq. 6.2 in the full periodic case, we shall briefly describe the Choleski algorithm as used in this code. In the remainder of this section, we shall drop the subscripts i and g. The basic idea of the method is to factor the matrix into the product of a lower triangular matrix and an upper triangular matrix (Gauss elimination). Thus, starting with a system Ax = g, where A is a tridiagonal matrix, we obtain the pair of equations

$$L\omega = g;$$

$$Ux = \omega.$$
(6.3)

Here L is lower triangular, and U is upper triangular and has the form

$$U = \begin{bmatrix} 1 & -\gamma_1 & & & & & \\ & 1 & -\gamma_2 & & & & \\ & & 1 & -\gamma_3 & & \\ & & & & -\gamma_{p-1} \\ & & & & 1 \end{bmatrix}.$$
 (6.4)

A forward sweep yields the sequences $\{\gamma_j\}$ and $\{\omega_j\}$; the backward sweep yields the solution $\{x_i\}$. Thus if

$$A = \begin{bmatrix} e_1 & -b_1 & & & \\ -b_1 & e_2 & -b_2 & & \\ & & & & \\ & & -b_{p-1} & e_p \end{bmatrix}, \tag{6.5}$$

then

$$\gamma_{1} = \frac{b_{1}}{e_{1}}, \quad \omega_{1} = g_{1}/e_{1};$$

$$\gamma_{j} = \frac{b_{j}}{e_{j} - b_{j-1}\gamma_{j-1}}, \quad \text{for } j = 2, 3, ..., p - 1;$$

$$\omega_{j} = \frac{g_{j} + b_{j-1}\omega_{j-1}}{e_{j} - b_{j-1}\gamma_{j-1}}, \quad \text{for } j = 2, 3, ..., p.$$
(6.6)

The backward sweep that generates the solution is then given by

$$\mathbf{x}_{p} = \omega_{p};$$
 $\mathbf{x}_{j} = \omega_{j} + \gamma_{j} \mathbf{x}_{j+1}, \quad \text{for } j = p-1, p-2, ..., 1.$
(6.7)

Note that the sequence $\{\omega_j\}$ depends on the vector g, whereas the sequence $\{\gamma_j\}$ depends only on the matrix elements; thus the sequence $\{\gamma_j\}$ is calculated only once and then stored.

We shall now consider the iterative method used to solve the periodic problem,

$$J_{y} = f, (6.8)$$

where J (without the ig subscripts) is defined by 6.1. Let the matrix A be the result of setting $b_J = 0$ in the matrix J. Then we rewrite 6.8 in the following form:

$$Ay = \begin{bmatrix} f_1 + b_J y_J \\ f_2 \\ \vdots \\ f_{J-1} \\ f_J + b_J y_1 \end{bmatrix}.$$
 (6.9)

Let $y^{(0)}$ be an initial guess. Then we will generate $y^{(1)}$ from

$$Ay^{(1)} \begin{bmatrix} f_1 + b_J y_J^{(0)} \\ f_2 \\ \vdots \\ \vdots \\ f_J + b_J y_1^{(0)} \end{bmatrix}. \tag{6.10}$$

Since A is tridiagonal, we can use Choleski's algorithm to generate $y^{(1)}$. Thus we form

$$\gamma_1 = b_1/e_1$$
, $\gamma_j = \frac{b_j}{e_j - b_{j-1}\gamma_{j-1}}$, for $j = 2, 3, ..., J - 1$,

$$\begin{aligned} \gamma_1 &= b_1 \big/ e_1, & \gamma_j &= \frac{b_j}{e_j - b_{j-1} \gamma_{j-1}}, & \text{for } j &= 2, 3, ..., J-1, \\ \\ \omega_1^{(1)} &= \frac{f_1 + b_J y_J^{(0)}}{e_1}; & \omega_j^{(1)} &= \frac{f_j + b_{j-1} \omega_{j-1}}{e_j - b_{j-1} \gamma_{j-1}}, & j &= 2, 3, ..., J-1, \\ \\ \text{with} & \\ \omega_J^{(1)} &= \frac{f_J + b_J y_1^{(0)} + b_{J-1} \omega_{J-1}}{e_J - b_{J-1} \gamma_{J-1}}. \end{aligned}$$

With the sequence $\left\{\omega_{j}^{(1)}\right\}$ thus generated on the forward sweep, we generate the sequence $\left\{\hat{y}_{j}^{(1)}\right\}$ on the backward sweep:

$$\hat{y}_{J}^{(1)} = \omega_{J}^{(1)};$$

$$\hat{y}_{J-1}^{(1)} = \omega_{J-1}^{(1)} + \gamma_{J-1}\hat{y}_{J}^{(1)};$$

$$\vdots$$

$$\hat{y}_{2}^{(1)} = \omega_{2}^{(1)} + \gamma_{2}\hat{y}_{3}^{(1)};$$

$$\hat{y}_{1}^{(1)} = \omega_{1}^{(1)} + \gamma_{1}\hat{y}_{2}^{(1)}.$$
(6.12)

At this point, we shall modify the definition of the sequence $\left\{ \hat{y}_{j}^{(1)} \right\}$. We define the sequence $\left\{ y_{j}^{(1)} \right\}$ by

$$y_{j}^{(1)} = \hat{y}_{j}^{(1)} \text{ for } j = 2, 3, ..., J,$$

and

$$y_{1}^{(1)} = \frac{f_{1}}{e_{1}} + \frac{b_{J}}{e_{1}} y_{J}^{(1)} + \gamma_{1} y_{2}^{(1)} = \omega_{1}^{(1)} + \gamma_{1} y_{2}^{(1)} + \frac{b_{J}}{e_{1}} \left(y_{J}^{(1)} - y_{J}^{(0)} \right).$$

Thus the modification consists of requiring $y_1^{(1)}$ to satisfy the first equation of the system 6.9 with the latest estimate available for y_J .

With $\{y^{(1)}\}$ thus generated, the process is repeated, thereby generating $y^{(2)}$, ..., $y^{(\ell)}$. The iteration is complete when the last equation in the system 6.9 is satisfied to a sufficient degree of accuracy. Hence we terminate the iteration when

$$\left| -b_{J} y_{1}^{(\ell)} - b_{J-1} y_{J-1}^{(\ell)} + e_{J} y_{J}^{(\ell)} - f_{J} \right| < \epsilon_{\phi_{I,I}}.$$
 (6.13)

The code actually stores $\left\{\gamma_j\right\}$ rather than $\left\{e_j\right\}$; however, the code does form γ_J , defined by

$$\gamma_{\rm J} = \frac{1}{e_{\rm J} - b_{\rm J-1} \gamma_{\rm J-1}}.$$
 (6.14)

From this, e_J can be found and used in Eq. 6.13. Note that this error criterion is on one component of the residual and thus has nothing to do with the rate of convergence of the iterative process.

The periodic iterations differ fundamentally from up-scattering iterations in the following sense. We saw in Section IV that for a source calculation (and the same thing is true for a keff calculation) it does not matter how many up-scattering iterations are performed per sweep through the channels. The iterative process will still converge to the same solution although the rate of convergence will be affected (see Appendix A, Sample Problems 8 and 9). The situation with regard to periodic iterations is entirely different. Within each channel and for each group, we must find the solution, or at least an approximation to the solution, of the system 6.2. However, this solution is the solution to a periodic system. The iterates $y^{(1)}$, $y^{(2)}$, ..., $y^{(\ell)}$ we generate to approximate the periodic solution are themselves solutions of a nonperiodic system. Moreover, the higher levels of iterations (up-scattering or outer) will not modify the lack of periodicity. Thus the periodicity of the solution to the source problem or keff problem is determined by the systems 6.2. From these we must obtain essentially periodic solutions (see Appendix A, Sample Problems 6 and 7).

VII. INPUT DATA PREPARATION

In general, ANL-CANDID conforms to the Argonne Standard Reactor Code input as specified in ANL-7194. A type 100000 identification card is followed by a set of CANDID code-dependent cards (Type 1).

A. Type 1 Cards

These cards contain data peculiar to ANL-CANDID: user options, problem constants, convergence criteria, criticality search specifications, and buckling and volume source data. All data values listed below take the default options noted, unless specifically input otherwise. The card type is punched beginning in column 1. CANDID reads the cards into one of three arrays (entitled INTERFCE, IOPTIONS, and B) depending on column 2.

1. Type 11 cards for user Options. A blank means the option will not be exercised; any nonblank character causes execution of the option. Standard problems have nonblank columns 8 and 10.

INTERFCE Subscript	Columns	Contents
	1-2	Must contain 11.
	3-6	Undefined.
1	7	Process all input cards, but do not execute the problem (input check only).
2	8	Print all fluxes (normally selected).
3	9	On-line output (for programmer use only).
4	10	Print interface fluxes and currents (normally selected).
5	11	Execute slower critical search procedure (i.e., run each intermediate k calculation to full, specified convergence). This option is normally blank, implying the program will decide when a control change is to be made.
6	12	Do not use Chebyshev acceleration during a k calculation. This option is normally blank, implying Chebyshev acceleration is to be used to allow acceleration to speed convergence.
7	13	Print modified cross sections with every control change during a composition search. Normally selected if a composition search.
8	14	Must be blank.

INTERFCE Subscript	Columns	Contents
9	15	Fission fraction (χ) vector is the same for all materials in cross-section set and does not need homogenization into a (space-consuming) matrix for every composition, normally selected, if possible, to conserve storage.
10	16	Perturbation problem is to be executed.
11	17	Two-tape perturbation (real flux on logical unit 49, adjoint flux and original cross sections on logical unit 48). Otherwise, the perturbation program will expect one flux tape consisting of real and adjoint flux and original cross sections.

2. Type 12 cards for Integer data. Integers must be right adjusted and fully contained within their six-column fields.

IOPTIONS Subscript	Columns	Contents
121 card	1-3	Must contain 121.
	4-6	Undefined.
1	7-12	Number of energy groups in cross-section set.
2	13-18	Number of energy groups in problem.
3	19-24	Number of regions in problem.
4	25-30	Maximum number of outer iterations desired. Normally set to 10000 since program will stop on time limit.
5	31-36	Maximum number of up-scattering iterations (for up-scattering problems only). Enter 1 if not up-scattering problem.
6	37-42	Maximum number of fixed fission source iterations (generally not used).
7	43-48	Maximum number of extrapolations on $k_{\mbox{eff}}$ during a search problem (4 assumed).

[†]If up-scattering iterations are performed, the convergence rate of the problem is increased. However, experience has indicated that the increase is not sufficient to offset the cost in time for performing the iterations. Thus, it is recommended that up-scattering problems be tried without using up-scattering iterations first, and these iterations only resorted to if unusual convergence difficulties arise.

IOPTIONS Subscript	Columns	Contents
8	49-54	Maximum number of interpolations on $k_{\mbox{eff}}$ during a search problem (20 assumed).
9	55-60	Print frequency during a search (i.e., number of control changes in criticality search before indicative printing is performed). Normally enter a 1 to imply printing after each control change.
10	61-66	Number of general interfaces (feature not available).
11	67-72	= 0 if new problem.= 49 if a restart problem. In addition, input final k from previous run on 500000 card.
122 card	1-3	Must contain 122.
	4-6	Undefined.
12	7-12	X direction mesh-refinement factor (relative to mesh of input flux guess or restart flux) (1 assumed).
13	13-18	Y direction mesh-refinement factor (1 assumed).
14	19-24	Perturbation-point output key: 61 for printer; 62 for card punch; 0 if point perturbation output not desired; another value (1-49) if output is desired on user logical unit (0 assumed).
15	25-30	Maximum number of periodic iterations when doing an $r\theta$ full-circle problem. Normally set to 2.

3. Type 13 cards for Decimal data. All numbers must have decimal points.

B Subscript	Columns	Contents
131 card	1-3	Must contain 131.
	3-6	Undefined.
1	7-18	Desired final value of $k_{\mbox{eff}}$ when doing a critical search (1.0 assumed).
2	19-30	Upper bound on search-control parameter (X) (0.0 assumed).

B Subscript	Columns	Contents
3	31-42	Second guess (X_2) of search-control parameter (0.0 assumed). (First guess on Type 500000 card.)
4	43-54	Lower bound on search-control parameter (0.0 assumed).
5	55-66	Estimate of dk/dx (used to calculate X_2 if nonzero). Normally used, if k(X) is approximately known, to eliminate the second control pass required to compute dk/dx. Estimate may also be available from previous restart runs.
132 card	1-3	Must contain 132.
	3-6	Undefined.
6	7-18	Flux-convergence criterion on outer bounds (k, k) of k (0.001 assumed).
7	19-30	Not used.
8	31-42	Smallest flux used in computing monitor information (10^{-6} assumed).
9	43-54	Agreement required in neutron balance equation (0.001 assumed).
10	55-66	Estimate of initial $k_{\rm eff}$ in critical search problem (1.0 assumed).
133 card	1-3	Must contain 133.
	3-6	Undefined.
11	7-18	Extrapolation parameter for fixed fission source iteration (1.0 assumed). Normally not used.
12	19-30	Up-scatter iteration convergence criterion (0.001 assumed). Normally not used. (See footnote on p. 45.)
13	31-42	Flux convergence criterion on sum of flux differences [†] (0.001 assumed).

For a k-calculation, three convergence criteria are required on the outer iterations:

We recommend that these all have the same order of magnitude, ranging from 10^{-4} to 10^{-7} , with a "normal" value of 10^{-5} . That is, a convergence of 10^{-4} may be desired for some purposes, but does not generally yield sufficient accuracy. A convergence of 10^{-7} yields sufficient accuracy, but at a large cost in computer time. Hence, this convergence is not generally used unless the results are to be input to the perturbation code where small perturbations are to be made.

^{1.} Outer bounds of k: $\underline{\mathbf{k}} \leq \mathbf{k}^{(n)} \leq \overline{\mathbf{k}}$ (132 card)

2. Sum of flux differences: $\left[\sum_{ijg} \left(\phi_{ijg}^{(n)} - \phi_{ijg}^{(n-1)}\right)^2\right]^{1/2}$ (133 card)

^{3.} k_{eff} differences: $k^{(n)} - k^{(n-1)}$ (Type 500000 card)

B Subscript	Columns	Contents
14	43-54	Not used.
15	55-66	Asymptotic inverse reactor time period (α).
	1-3	Must contain 134.
	3-6	Undefined.
16	7-18	Initial estimate of outer iteration convergence rate used in computing Chebyshev extrapolation coefficients (0.995 assumed). Normally set to 0.995. (See page 37.)
17	19-30	Periodic iterations convergence criterion (r θ full circle). Normally set to 0.0001.
18	31-36	Interval shift factor for Chebyshev acceleration (0.0 is assumed, implying no shift). Normally set to 0.0.

4. Type 14 and 15 cards for Buckling and Volume Source (decimal) data. Data can be input as one constant, as varying with energy only or with region only, or as varying both with energy and region number.

Columns	Contents
1 - 2	Must contain 14 for Buckling data and 15 for Volume Source data.
3-6	Region number, given only when data are to be distributed into regions.
$ \begin{array}{c} 7-18 \\ 19-30 \\ 31-42 \\ 43-54 \\ 55-66 \end{array} $	Buckling or Volume Source data, as named in columns 1-2, given on as many cards as necessary to cover all energy groups for this region, before beginning next region.
67-69	Energy-group number of first group on this card, given only when data are to be distributed by energy groups.

The remaining cards in the input deck are standard, except as noted below.

B. Type 2XXXXX Cards

The Type 210000 cards for geometry specification must be consistent with mesh of the input flux guess. If a problem is being restarted using an old flux tape, the number of mesh points and the region structure given on the Type 2 cards must correspond to the tape structure. Mesh "multiplying" in

X or Y directions is accomplished by setting the values of the X and Y mesh-refinement factors (IOPTIONS (12) and (13) on the 122 card). The factor input is the number of mesh points to be made from each input mesh point. The problem is then run using the new, "full" mesh. Subsequent restart problems to be run using the new mesh will require a new set of Type 2 cards consistent with this mesh (and the removal of mesh-multiplying factors IOPTIONS (12) and (13)). That is, mesh refinement is used only once for each desired mesh size without redefining the mesh in the Type 2 cards.

The Type 220000 cards for the ratio-method geometry specification is not implemented.

C. Type 3XXXXX Cards

Column 60 is ignored; i.e., general boundary conditions are not implemented.

On the remaining Type 3XXXXX cards the only restriction on card ordering is that the sides be input one at a time, until all regions and groups on a side are covered, with any ordering of data within a side allowed.

D. Type 4XXXXX Cards

Both Type 400000 and 410000 cards must always be present. That is, direct isotope-to-composition homogenization is not implemented.

E. Type 5XXXXX Cards

On the Type 500000 card, the value for δ (columns 13-24) is used as the convergence criterion on differences between k values; for generation n, the criterion is satisfied if $|\mathbf{k}(n) - \mathbf{k}(n-1)| < \delta$ (see footnote on p. 47). The initial guess supplied in columns 46-57 is used as an initial value for k in a flux calculation, or for μ in a source calculation (blank or 0 signals the default option of 1.0). The k corresponding to a restart flux must be used for a restart problem. If the problem is a criticality search, the initial guess, i.e., the control parameter (X_1) , is placed in this field, and the initial k is put on the 132 card. The remaining Type 5XXXXX cards are implemented as specified. All cards with region numbers must have them in increasing order. Energy-dependent buckling modifiers can be input using the concentration change format, with columns 2-6 containing a region number, if desired, and columns 19-24, 37-42, and 55-60 containing group numbers corresponding to modifiers in columns 7-18, 25-36, and 43-54.

F. Type 6XXXXX Cards

On the Type 600000 card CANDID requires the initial guess to be a flux guess and not a source; therefore column 7 must be nonblank. Spectrum averaging is not performed. External source must be read from Type 15 cards.

G. Type 7XXXXX Cards

These cards are as specified in ANL-7194. In addition to the normal use of cross-section set modification, these cards are used to input cross sections for use in perturbation calculations.

[†]ANL-CANDID expects microscopic cross-section data on a magnetic tape prepared by the XLIBIT Code.

An XLIBIT cross-section set must be available on tape for each problem even if all cross-section data are to be input via the 700000 cards.

VIII. EXECUTION CARD DECK ARRANGEMENTS

A. Normal Execution

The control cards and deck structure necessary to execute a normal ANL-CANDID run are $\binom{7}{9}$ punches are assumed in column 1):

JOB, User Identification
MOUNT XLIBIT Cross-section Tape on LUN 44
MOUNT CANDID AOV Tape on LUN 47 (AOV means absolute overlay)
SCRATCH TAPES on LUN 42, 43, 45, 55 (SYSTEM scratch unit)
EQUIP,44=(XLIBIT Cross-section tape name, Version) RO, SV
EQUIP,3=45 (Equivalent reference saves one scratch tape)
EQUIP,47=(CANDID AOV tape name, Version), RO, SV
LOAD MAIN, 47, Time estimate, Print estimate
Data
END OF PROBLEM (column 3)

B. Program Compilation

The volume of source program cards (some 20000) in ANL-CANDID prohibits the handling of source decks directly. Thus, the cards are put on tape along with SCOPE cards denoting the end of overlays. The following SCOPE control cards (in addition to special mounting instructions) will compile the source code from LUN 10 (unlabeled) and prepare the binary load and go (LGO) tape on LUN 11. Note that CANDID is compiled as a one-bank code.

```
MOUNT TAPE (ANL-CANDID source tape reference number) ON
  LUN 10
EQUIP,10=SV,**
EQUIP, 11=(CANDID LGO tape name, Version), SV
FILE, 11
MAIN,47
FILE END
                        Main Control Section
FTN,L,X=11,I=10,R,*
COMPASS, L, X=11, I=10
FILE, 11
OVERLAY,47,1
                        Overlay 1: Processes card Types 1 and 2.
FILE END
FTN,L,X=11,I=10,R,*.
FILE, 11
OVERLAY,47,2
                        Overlay 2: Processes card Types 3-5.
FILE END
FTN, L, X=11, I=10, R, *.
```

FILE, 11 Overlay 3: Creates pseudo cross-section OVERLAY,47,3 FILE END library tape from XLIBIT tape. FTN,L,X=11,I=10,R,*. FILE, 11 OVERLAY,47,4 Overlay 4: Processes card Type 7 if FILE END they exist. FTN,L,X=11,I=10,R,*. FILE, 11 OVERLAY,47,5 Overlay 5: Performs search and k FILE END calculations. FTN,L,X=11,I=10,R,*. Overlay 6: Output edit. FILE, 11 OVERLAY,47,6 FILE END FTN,L,X=11,I=10,R,*. FILE, 11 OVERLAY,47,7 FILE END FTN, L, X=11, I=10, R,

C. Program Modification Using LGOEDIT

Additions, corrections, or deletions to the load-and-go tape are made using the program LGOEDIT (see Appendix B). The card deck structure listed below will accomplish the following (assuming the existence of an old load-and-go tape):

- 1. Compile the indicated source programs and place them, together with the associated LGOEDIT control cards (REP,DEL,INS), on the merge tape (LUN 45).
- 2. Select the first program name (NAME1) from the control card on the merge tape, and search the old LGO tape (LUN13) until NAME1 is found. Preceding subroutines are simply copied from LUN13 onto the new LGO tape (LUN11). The instruction on the control card is executed as described in Appendix B, and LGOEDIT proceeds to the next subroutine on the merge tape. Note that LGOEDIT does not space the old LGO tape, indicating that the ordering on the old LGO tape must be preserved on the merge tape.

MOUNT TAPE (Old LGO tape) on LUN13
EQUIP,13 = (Old LGO tape name, Version), RO,SV
EQUIP,11 = (New LGO tape name, Version), SV
FILE, 45
REP NAME 1
FILE END
FTN,L,X=45,*.

Source for Subroutine NAME1

FILE,45 DEL NAME2 FILE END FILE,45 INS NAME3 FILE END FTN,L,X=45,*.

Source for Subroutine NAME3
SCOPE

Binary Deck for LGOEDIT RUN 11 | 13 | 45 (Tape numbers as input to LGOEDIT columns 1-6)

D. Modification Followed by Execution

The normal execution structure (A) may be merged with the program modification procedure (C) to produce a run in which modification is followed by execution. The structure is as follows:

Normal control cards (A) Modification structure (C)

LOAD, 11 (Prepares AOV Tape on LUN 47) RUN Problem Data

E. Restart

A restart feature is provided for the following reasons:

- 1. Machine failure usually in the form of tape parity error.
- 2. Time estimate exceeded.

CANDID checks internally for an error condition, and when one is encountered, the current flux vector is written, as indicated in Fig. 4, onto a flux save tape (LUN 48). Normal termination will also result in the flux being written on LUN 48. Thus, to save the flux, the user must insert the following control card:

EQUIP,48=(Flux save tape name, Version), SV.

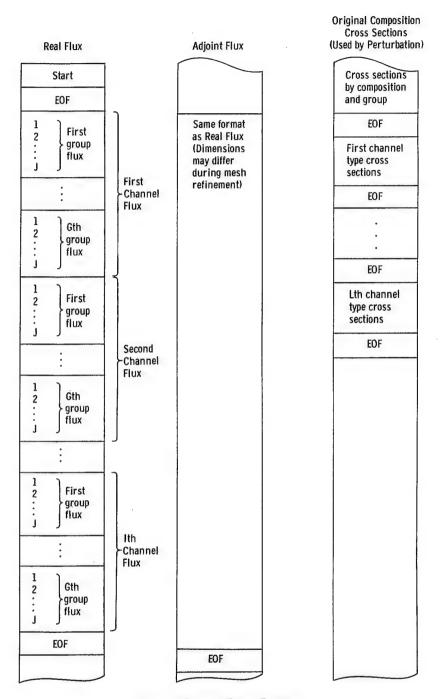


Fig. 4. Restart Tape Format

To restart the problem from this point, the user must remount the tape on the restart unit (LUN49) by inserting the following SCOPE cards:

Mount tape (Old Flux tape Identification) on LUN 49 EQUIP,49=(Old Flux Tape Name, Version), RO,SV

The original input-data deck is used with the exceptions (1) the number 49 must be punched on the 12 card (see Section VII), and (2) the

 $\binom{(n)}{\text{keff}}$, where n is the last iteration of the previous run, must be punched on the 500000 card (or on the 132 card if a search problem). Normal execution will then proceed with the flux guess taken from LUN 49 instead of the standard input unit. The new flux may again be saved on LUN 48. (See Section IX.)

F. Back-to-back Runs

The stacking of runs back-to-back is provided, but is not usually done except when perturbation problems are run. The reason for this is that two-dimensional problems of any size require so much machine time that they are run in stages. (See Section IX.)

A complete exit to the SCOPE System is made following each End of Problem card. Thus, the program must be reloaded for each problem. This is done to allow for any available mounting or dismounting type instructions to be recognized between problems. Thus, in stacking problems, the deck structure is:

Normal execution control cards (A)
Program modification structure, if any (C)

LOAD, RUN

OR

LOADMAIN,47

Data Problem 1

END OF PROBLEM

Tape-handling instructions (i.e., unloading old flux tape, mounting new flux tapes for next problem, etc.)

LOADMAIN,47 Data Problem 2 END OF PROBLEM

Tape-handling instructions LOADMAIN,47 Data Problem n END OF PROBLEM End of File

G. Sample Deck Structure

A sample deck structure including program modification, restart, and back-to-back problems might appear as follows:

JOB, User Identification

MOUNT OLD LGO Tape ON LUN11

MOUNT XLIBIT Cross-section Tape on LUN 44

MOUNT Old Flux Tape for Problem 11 on LUN 49

SCRATCH Tapes on LUN's 42,43,45,55

EQUIP, 11= (NEW CANDID LGO, L), SV

EQUIP,13=(OLD CANDID LGO,),SV

EQUIP,3=45

EQUIP,44=(XLIBIT Cross-section Tape Name, Version),SV

EQUIP,47=(CANDID AOV,2),SV

EQUIP,48=(NEWFLUX1,1),SV

EQUIP,49=(OLDFLUX1,1),SV

FILE,45

REP NAME1

FILE END

FTN,L,X=45,*.

Source deck for Subroutine NAME1

SCOPE

Binary deck for LGOEDIT

RUN

131145

(LGOEDIT input card)

End of File

Release, 13

(Special ANL Control Card which unloads 13)

Load, 11

Release, 11

RUN

CANDID Data Deck for Problem 1

End of Problem

End of File

Release, 48

Release, 49

Mount old flux tape for Problem 2 on LUN 49

EQUIP,48=(NEWFLUX2,1),SV

EQUIP,49=(OLDFLUX2,1),SV

LOADMAIN,47

CANDID Data Deck for Problem 2

End of Problem

End of File

IX. NORMAL EXECUTION PROCEDURES

A two-dimensional diffusion problem is usually run in stages (a sequence of computer runs) for two reasons:

- 1. The probability of error in problem input data (even if it gets by the input processor, which checks for consistency only) is significant. Thus, we would like to be assured that the problem is proceeding in the "right" direction before spending a significant amount of computer time.
- 2. Even when we are assured that the problem is progressing smoothly, it is worthwhile to continue in stages of less than 60 min of computing time to measure convergence rates, progress, etc.

The stages mentioned above are of two types:

- l. Mesh size: i.e., starting with a coarse mesh, running the coarse-mesh problem, and refining the mesh using the coarse-mesh flux as an input flux guess for the refined problem.
- 2. <u>Computer time</u>: i.e., putting a time limit on execution of something less than 60 min, coming off the computer to measure progress, and restarting and continuing the problem from that point.

Experience has shown that a 20-group, 40 x 40 mesh may take 3 to 4 hr of computing time (depending on accuracy requirements and the rate of convergence of the problem). Hence, it is wise not to be hasty in expending this amount of time only to find (possibly because of an error or incorrect assumptions) that the results are not what was expected. The problems are generally run as follows:

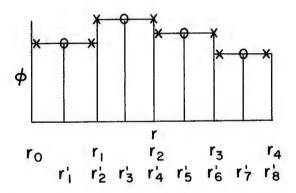
- 1. The mesh is reduced or "coarsened" to a problem that will run in 20 or 30 min (e.g., 10 x 10 x number of groups); the results are expected for "ballpark" accuracy. Experience indicates that this result is certainly within 10% of the final answer and may be as good as 1%. Errors in configuration should show up at this time. At the end of the job, the flux tape is saved as indicated under Execution Card Deck Arrangements, Section VIII.
- 2. After step 1 is completed and the refined problem is to be attempted, the mesh-refinement feature is used. This amounts to placing a multiplication factor (integer) for the X and/or Y direction on the 122 card (see Section VII). For example, a 2 placed in the field for the X direction will double the mesh along this axis. (The coarse problem must be planned so that an integral multiple of the coarse mesh will give the final desired mesh.) CANDID now automatically spreads the coarse mesh to adjacent points in the fine mesh, as indicated in Figs. 5 and 6. The resulting flux guess is then used to start the problem. Refinement can take place as often as desired along either axis.

_		X	V	W					
/ -	L.	^	X	X	X	X	X	X	X
Z ₂	_	X	X	X	X	X	X	X	X
7	Z	X	X	X	X	X	Х	X	X
Zı		X	X	X	X	X	X	X	X
zo		X	X	X	X	×	X	X	X

Coarse mesh: O Flux point Refined mesh: X Flux point

The orginal 3 x 4 mesh (points 0) 6 x 8 (points X)

Fig. 5. Mesh Refinement* (Doubling in r and z directions)



Note: An interpolation scheme would give a better approximation, but one has not been incorporated into ANL—CANDID.

Fig. 6. Typical Flux Shape for Starting the Refined Problem*

^{*}Unprimed mesh lines denote coarse mesh.
Primed mesh lines denote refined mesh.

X. COMPUTING REQUIREMENTS AND EXECUTION TIME

- A. Machine: CDC 3600 with one-bank 32K fast memory.
- B. Operating System: SCOPE
- C. Tape Units: Normal Execution: Four scratch tapes, cross-section tape, AOV tape.

Program Modification: Four scratch tapes, cross-section tape, old LGO tape, new LGO tape, AOV tape.

Restart: Four scratch tapes, cross-section tape, AOV tape, old flux tape, new flux tape.

D. Execution Time: Execution time depends on the convergence criterion and the problem convergence rate. Typically, a 40 x 40 mesh, 20-group problem will run 3-4 hr. A 10 x 10 mesh, 20-group problem will run 20-30 min.

XI. RESTRICTIONS AND LIMITATIONS

- 1. Roughly, the largest problem that can be handled is 20 groups by 40-45 points in a channel (i.e., Y or Z or θ direction). Any number of points can be handled in the R direction if one has enough money.
- 2. The number of compositions cannot exceed the number of materials in the cross-section library. "Dummy" isotopes should be inserted in the cross-section set if this situation arises.
- 3. Composition cross sections cannot be directly homogenized from isotopes. The material level cross sections must be inserted even though they may be decimals (i.e., volume fractions are 1.0).
 - 4. A source guess cannot be input. Only a flux guess is allowed.
- 5. The distributed source for a source calculation must be input by region and group on the Type 15 cards. CANDID will not accept an external source input via the 610000 cards.
- 6. Generally, the output is erroneous for adjoint calculations, but the flux printout is correct.

XII. SUMMARY OF REQUIREMENTS

Execution of ANL-CANDID requires familiarity with the following documents, in addition to ANL-7305:

- 1. Code independent input and cross-section library tape preparation are described in the following documents:
 - M. Butler and H. Greenspan, General Input Specifications for ANL Reactor Programs, ANL-7194 (April 1966).
 - S. D. Sparck, XLIBIT: An ANL Cross-section Library Code, ANL-7112 (Feb 1966).
- 2. To make modifications and/or additions to the CANDID load-and-go tape, the use of LGOEDIT is helpful:

LGOEDIT2, reproduced as Appendix B by permission of the authors.

- 3. For perturbation calculations, the following publication is required:
 - G. K. Leaf and A. S. Kennedy, PERC, A Two-dimensional Perturbation Code Based on Diffusion Theory, ANL-7304 (May 1967).
- 4. Finally, if available, the Control Data Corporation documentation of CANDID2D may be helpful:

CANDID2D, A Two-dimensional Neutron Diffusion Program, Control Data Corporation (1966).

The following items are necessary to run ANL-CANDID:

- 1. Source (BCD FORTRAN Card images) tape.
- 2. Load-and-go tape prepared by compiling source tape.
- 3. Absolute overlay tape prepared from the load-and-go tape.
- 4. XLIBIT generated cross-section library tape.
- 5. Proper SCOPE control cards.
- 6. Input deck which conforms to ANL-7305 and ANL-7194.
- 7. Lots of luck!!

APPENDIX A

Sample Problems

The sample problems in this appendix reinforce the conclusions drawn in the body of the report. Hence only relevant parts of each output listing are included. For completeness, sample problem 1 is listed in its entirety.

For reference, the tests for convergence are listed below:

1. Sum of flux differences =
$$\left[\sum_{i,j,g} \left(\phi_{ijg}^{(n)} - \phi_{ijg}^{(n-1)}\right)^{2}\right]^{1/2},$$

where

(i,j) are the point indices,

g is the group index,

and

n is the outer iteration index.

2. Bounds of $k_{eff} = \overline{k} - \underline{k}$,

where

$$\overline{k} = \max_{ijg} \left\{ \frac{\phi_{ijg}^{(n)}}{\phi_{ijg}^{(n-1)}} \right\}$$

and

$$\underline{\mathbf{k}} = \min_{\mathbf{ijg}} \left\{ \frac{\phi_{\mathbf{ijg}}^{(n)}}{\phi_{\mathbf{ijg}}^{(n-1)}} \right\}.$$

3.
$$k_{eff}$$
 difference = $\begin{vmatrix} k_{eff} - k_{eff} \end{vmatrix}$.

4. Criticality search =
$$k_{eff}$$
 - $k_{desired}$.

5. Up-scattering =
$$\frac{\sum_{jg} \left| \phi_{jg}^{(m)} - \phi_{jg}^{(m-1)} \right|}{\sum_{jg} \left| \phi_{jg}^{(1)} - \phi_{jg}^{(0)} \right|}$$

6. Periodic =
$$\left| -b_{J}y_{1}^{(\ell)} - b_{J-1}y_{J-1}^{(\ell)} + e_{J}y_{J}^{(\ell)} - f_{j} \right|$$
.

1. Sample Problem 1

a. Description

(1) Problem Type

Real keff calculation without Chebyshev acceleration.

(2) Configuration

(a) Geometry

rz

(b) Region Definition

The reactor consists of three regions as follows:

Region No. 1: Core region composed of uranium $(U^{235} \text{ and } U^{238})$, stainless steel, molybdenum, zirconium, iron, nickel, chromium, and sodium.

Region No. 2: Control- and safety-rod region composed of uranium (U²³⁵ and U²³⁸), molybdenum, niobium, iron, nickel, chromium, and sodium.

Region No. 3: Radial-blanket region composed of uranium (U²³⁵ and U²³⁸), plutonium (Pu²³⁹), iron, nickel, chromium, and sodium.

(c) Mesh Definition

r direction: 27 points

z direction: 20 points

(d) Boundary Conditions

Left: $\phi' = 0$.

Right: $\phi = 0$.

Bottom: $\phi^{\dagger} = 0$.

Top: $\phi = 0$.

(e) Number of Energy Groups: 2

(3) Convergence Criteria

 $k_{\mbox{eff}}$ difference = 10^{-5} .

 k_{eff} bounds = 10^{-5} .

Sum of flux difference = 10^{-5} .

PAGE NO.											
27 MESH ZERO,				N N	10000	404	64	1.0250+000	1.5000+000	1.0000 0003	1,0000+000
PROB, NO, 2:00000-080 20CANDID K=CALC (NO ACC), RZ, TWO GROUPS, 20 X 2 CO CODE DEPENDENT DATA OR INPUT IF DIFFERENT THAN A	BUCKLING DATA HAVE NOT BEEN INPUT.	VOLUME SOURCE DATA HAVE NOT BEEN INPUT.	ALL FLUX VALUES WILL BE PRINTED, CHEBYCHEV ACCELERATION WILL NOT BE USED.	ENERGY GROUPS IN CROSS SECTION DATA	SNO	OF EXTRAPOLATIONS ON X*EFFECTIONS ON X*EFFECTI	FREQUENCY ON FISSION SOURCE K	DESTRED KEEFFECTIVE	ARCH PARAMETER	FUR SOURCES:	SS AT K-EFFECTIVE

b. Output Listing

CV

PAGE NO.

PROB, NO. 2:00000+080 RDCANDID K=CALC (NO ACC), RZ, TWO GROUPS, 20 X 27 MESH

GEOMETRY	LEPT	RIGHT	BOTTOM	TOP	
	BOUND	BOUND	BOUND	BOUND	
RZ WITHOUT INVERSION	0.000000000	7,8100+001	0.0000.000	6.0000+001	
		REGION BOUNDAHIES	HES		
REGION NO.	LEFT BDRY.	RIGHT BDRY.	BOTTOM BDRY.	TOP BDRY.	
wit (0.0000+000		000000000000000000000000000000000000000	6.0000+001	
ed ke	2,5902+001 2,5902+001	7,8100+001 7,8100+001	0.000040000	6.0000#001	
PROB, NO. 2 000000+000	ADGANDID K-CALC	(NO ACC) , RZ, MESH DEFIN	TWO GROUPS, 20 X 27 MESH ITION	O E	in
	OZI	GREMENT METHOD - XC	X(H) DIRECTION		
X(R) INCREMENTS	XCR) INC	GREMENTS MCR.	LINDREMENTS	ENGREENE (X)X	X(R)
0,0000+000 12 IMAX 8 27	2,3168+001	3 2,5902+001	1 12 7,8100+001	1000	
	X(R) PLUS DE	ELTA X(R) FROM LEFT	TO RIGHT OF REACTOR		
MESH NO. 1.9307+000 3,8614+806	5.7920+000 7.	7227*000 9.6533+600	1,1584+001 1,3515+001	1 1.5445+001 117876+001	100+10001
MESH NO. 12 1237+001 2,3168+808	13 2,4079+001 2,	14 1991+001 2.5902+001	16 1602+01 3,4602+0	17 18 18 18 17 17 17 17 17 17 17 17 17 17 17 17 17	19 20 20
ABOUT NO. IN COLUMN IN CASE OF STREET	23	24	26	27	

			INGREMENT METHOD . Y(2) DIRECTION	# Y(Z)	DIRECTION				
Y(Z)	Y(Z) INCREMENTS	Y(Z)	INCREMENTS	(Z) Å	Y(Z) INCREMENTS	'n	(Z) A	Y(Z) INGREMENTS	Y(Z)
0.0000.0	20	£00000#9							

1000-000	2000+000
001 3,	19 100
217800+	5:7800*
2,4000+001	5,4000+001
2,1000+001	5,1000+001
1,8000+001	4.8000+001
(2) MESH NO. 1 8 8 9 10000+000 1,2000+000 1,2000+001 1,5000+001 1,8000+001 2,1000+001 2,4000+001 2,7800+001 3,0000+001	12) MESH NO. 11 18 1.8 1.9 1.0 1.4 1.2000+001 4.5000+001 4.8000+001 5.1000+001 5.4000+001 5.7800+001 6.0000+001
1,2000+001	14,2000+001
9,0000+000	3,9000+001
6.0000+800	3,6000+000
3,0000+000	3,3000+001
Y(Z) MESH NO. ORDINATE 3.0	OKDINATE 3,3000

THE STATE OF THE S		1		-					(RE	REGION	NUMBER		BY ME	I	POINT	7										MESH
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8 6857698 004

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989904*U01 9.9996964*P990141*U01 9.9997040*P90904119*U01 9.9997114*P909056*U01 9.9997184*P9090904*P01				
01 9.9997040 01 9.9997114 01 9.9997186	-001 Z.Y1/8/08-00		+00000000	.06-2386-00
01 9.9997114 01 9.9997186	-001 2.8451082-00	00+0000000.	+00000000	.8787296-0h
9,9997186	-001 2.7741745-00	.00000000.	+00.0000.	.7021218-0n
1 0.9907257	-001 2.7050210-00	0000000000	+0000000	.53.2505-0n
	-001 2.6376005-00	00+0000000.	+0000000	.3629952-0r
1 9,9997325	-001 2.5718713-00	.0000000+00	+0000000	.20.2087-0n
1 9,9997392	+001 2.5077898-00	00+00000000	+0000000	.0417689-0n
9,9997457	-001 2.4453138-00	00+0000000	+0000000	.8875390-0n
9,9997521	+001 2.3844021-00	.00000000.	+00000000	.7374040-0n
9,9997583	-001 2,3250165-00	.00000000.	+0000000	.5912475-0n
9,9997643	-001 2.2671172-00	.00000000.	+00000000	.4489487-0n
6.9997708	-001 2.2106676-00	00+00000000	+0000000	.31.3940-0n
9,9997760	-001 2,1556299-00	.00000000.	.0000°.00•	.1754963-00
9,9997815	-001 2,1019693-00	000000000.	+00-0000	. U441275-00
9,9997870	-001 2.0496491-00	00+00000000	+0000000	.9162045-0n
9,9997923	-001 1.9986397-00	0000000000	+00°0000.	.7916372-0n
9.9997975	-001 1,9489047-00	.00000000.	+0000000	.67-3106-0P
9,9998025	-001 1.9004117-00	.00000000.	+0000000	.5521418-0P
9,9998075	-001 1,8531303-00	00+00000000	+0000000	.437-478-00
.9998123	5-001 1.8070298-005	1,00000000001	1,0000,000+5	4.3249383-005
,9998169	-001 1.7620788-00	.000000000.	+00000000	.215/422-00
.9998215	-001 1.7182516-00	.00000000.	+0000000	.1093779-DP
,999826	-001 1.6755183-00	00+00000000	+0000000	.0057537-00
.999830	-001 1.6338520-00	.000000000.	+0000000	.9848013-0n
,9998345	-001 1.593238-00	.00000	+000000	4565-00

AN ERROR WAS ENCOUNTERED IN SUBROUTINE KEFFCALC AT STATEMENT NUMBER 5.4743382*003 5.0747050*003 4.9759841*003 4.8457066*003 4.8457066*003 4.8457066*003 4.2778406*003 4.2778406*003 3.47868*003 3.47868*003 3.46978*003 2.7987402*003 2.7987402*003

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20 IMAXOUT=

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20
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                                                27 JMAXOUT=
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   8.4912379-004
5.5544858-004
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1.3057572-003
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                                               20 IMAXOUT=
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NI I	コロチウウオオウルロ	9.224802±00	9,998427*0	1.5149896#00	000000000	000000000	3.6172743-00
*	8932355#00	9,9994939=00	9,9998466 6	1,4773270=00	00+00000000	00*0000000	3,5263176*00
4	8930285=00	9,9995066-00	9,9998504=0	1.4406020*00	00000000	00+000000	3,4376862,00
in	8928288-00	9,9995190 00	9,9998541*0	1,4047937=00	000000000	000000000	3,3513235*00
•	8926341-00	9,9995310=00	9,9998578*0	1,3698775=00	00000000	000000000000000000000000000000000000000	3.2671829=00
7	8924442m00	9,9995428=00	9,9998613=0	1,335831900	000000000	000000000	3,1851756=00
00	8922590-00	9,9995542=00	9,9998647.0	1,3026340=00	0000000000	000000000000000000000000000000000000000	3.1052710-00
0	8920785 00	9.9995654#80	9.9998681+0	1.2702635000	004000000	004000000	3.0273965*00
40	8919024*00	9 9 9 9 5 7 6 2 ± 0 n	9.0008714m	4.2384004#00	000000000000000000000000000000000000000		9.9545007+00
	8017307=00	0.909EAAR	0 0000 74Ke	1.2070204			8 8775505-00
4 0	00-1001740	000000000	0.0000000000000000000000000000000000000	00 P / 0 / 0 / 0 / 0 / 0 / 0 / 0 / 0 / 0			DECALCION D
4 5		00-07/40444	0 - / / 0 A A A A	00 200 6//10	00.00000000	0000000000	M 44604/4/400
7	00.00.000	W129900/6"U	V 99990017	11486434811	00000000	00000000	Z 132234/*UU
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2	8910822#00	8.9996266m00	0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0	1.0922791.00	000000000	0000000000	2.6000533±00
16	8909341 = 00	9,9996359+00	9,9998894*	1.0651473#00	00000000	00000000	8,5350091=00
17	8907865*00	9,9996450 = 00	9,9998922#0	1,0386894 00	00000000	00+0000000	2.4716224-00
∞	8906425*00	9,9996539*00	0.9998949	1,0128901,00	000000000	000000000	£,4098335*00
10	8905022*00	9,9996625*00	9.99989750	9.8773216=00	000000000	00+00000	2.3496090000
20	8903653*00	0.9996709=00	00006666666	9,6319914#00	000000000	0000000000	£.2909080 = 00
21	8902318-00	9,9996791=00	9.9999025	9.3927854#00	000000000	00*000000	2.2336826e00
25	8901016900	9,9996871=00	9.9999949	9.1595339*00	000000000	00000000000	9.1779066-00
23	8899746=00	9.9996949*0	9.9999073#	8.932064300	00000000	00000000	9.1235348 m00
2.5	8898509+00	9,9997025#0	9.0666666	8.7102580±00	00000000000		2.0705338=00
25	8897301 00	9.9997099e0	9,9999118+	8.4939650*00	0000000	000000000000000000000000000000000000000	9.0188643w00
26	889612400	9,9997172=0	9,9999140*	8,2830567 = 00	00000000	000000000	£.9684987=00
27	889497600	9,9997242"0	9,9999162*	8.0773800*00	00000000	0000000000	1.9194049±00
200	8893857=00	9.9997311+0	9.9999182	7.8768217000	00000000	00+0000000	4.8715393e00
58	9,8892765#001	9,9997378=001	9,9999203*	7,681	4.0000000+000	0.0000000	Hei
30	8891701=06	9,9997443=0	9,9999223*	7,4905669m0	0000000000	0000000	1.7793835 b0
31	8890663-06	9,9997507=01	9,9999242*	7,3045961m0(000000000	0000000000	1.7350365=0
21	8888650 = 01	9,9997569+01	9,9999261+	7,1232544=0	000000000	0000000000	\$,6918013e0
33	8888653e0(9,9997629"0	9,9999279	6,9464209=0	00000000	00000000	1,6496531=0
34	8887701-00	9.9997688*D	9,9999297*	6.7739731=0	00000000	00000000	£,6085614-0
35	,8886762•0	8.9997746×0	9,9999314	6.6058211 m Q (00000000	00000000	1,5685015=0
36	8885847.01	9,9997802=0	9,9999351*	6,4418510 m0	00000000	000000000	1,529435400
37	8884954e0	9,9997857=0	9.999348*	6.281945200	0000000	.0000000	# 4913559=0
38	8884083*0(0 4016466666	9,9999364	6.1260197=0	000000000	00000000	4,4542253*0
9	8883234 0	9,999796200	9,999380 e	5.9739717m0	00000000	.00000000	1 4180274 = 0
0.4	8882407.00	9,9998013=0	9,9999395	5,8256966#0	000000000	.0000000.	1,3827375#0
4 .	8881599#0	9,9998062=0	9,9999410	5,6811069#0	000000000	0.00000000	1,3483208m0
A .	8880812*0	9.998110=0	9 9999425	5,5400905*0	000000000	0.0000000	1,3147685=0
200	000000448U	9.9998157=0	9,999439*	5,4025879m0	00000000	0.0000000	4 . 2820572 0
4 1	8879296*0	9 9 99 8203 m0	9,999453	5,2684991#0	00000000	0+00000000	1.2501696.0
4	8878506*0	9,9998248*0	9,9999467	5,1377491°0	000000000	0.00000000	1,219066300
4	8877854*0	9,9998291=0	9,9999480	5,0102349m0	.00000000	.00000000	1.1887503.0

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2DCANDID K*CALC (NO ACC) , RZ, TWU GROUPS, 20 X 27 MESH

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99 SON. PAGE MEST 27 × K+CALC (NO ACC) , RZ, TWU GROUPS, 20 12 ARE 2 S FLUXES IN CHANNEL NO. FLUXES IN CHANNEL ZMESH ZMESH 45000000 125 5 20 4 6 2DCANDID 4500 10 0 **®** 0 # # # # 20 PS 3,05996=003 2,72895=003 2,33075=903 8,37791m004 1,37338m003 2,33075 903 2,72895 903 3,05996 903 1,85536m003 2,30615e003 2,70015e003 3,02767e003 3,02767#003 2,70015#003 2,30615#003 1,87515=003 8,28949-004 3,28063m003 3,45282m003 3,53999m003 3,53999m00 3,45282m003 3,28063m003 3,31563m003 3,48965*003 3,57775=003 3,57775=003 3,48965=003 3,31563+003 8,37794=104 2,81575-004 2,81575-004 1,87515=003 1,85536-003 1,35888#003 8,28949#004 2,78603=004 1,35888 003 GHOUR 2,000000+000 1,95324e002 2,42781e002 2,64261e002 3,18740e002 3,45371e002 3,72675e002 3,18740 m 0 0 2 2,84261 m 0 0 2 2,42781 m 0 0 2 1,43058-002 2,93302=005 3,44120 002 3,72872 002 3,92442 002 3,63498-002 3,45371.0002 1,95324=002 2,933020003 8,72683-003 1,43058-002 3,06895-002 4,02349=002 3,92442=002 3,72872=002 3,44120 - 002 3,06895 002 2,62113 002 1,54449=002 9,42170=003 3,16656#003 4,02349=002 2,10877-002 3,16656=003 2,10877-002 2,62113#002 9,42170m005 1,54449=002 GROUP GROUP PROB, NO. THESH 7 HS3H7 4592880 10

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RDCANDID KOCALC (NO ACC) , RZ, TWO GROUPS, 20 X 27

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4 PAGE NO X III NI 2 KECALC (NO ACC) , RZ, TWU GROUPS, 20 X ARE 20 ARE 21 FLUXES IN CHANNEL NO. FLUXES IN CHANNEL NO. 100 00 0 10 TOM 450000 SDCANDID 1,07877e003 9,62072e004 8,21687e004 6,61070e004 4,841749004 2,95857*004 9,926719005 6.91120 m009 2.05634m004 3.37093m004 8.13814#004 8.56528#004 8.76151#004 2,95857=005 6,61070 004 8,21687 004 1,26131+003 7,51062-004 8,78151=004 7,510625004 6,69816=004 5,72077=004 4,60251m004 3,37093 004 2,05634,004 6,91120=005 9,62072=004 1,23025-903 4,60251=004 5,72077=004 8.13814=004 CU 4,84174m004 1,16890-003 1,16890#803 1,07877-003 GROUP GROUP 2,00000+000 1,98555=003 2,42478=003 2,60677=003 3,53234,003 3,01690,003 2,42718,003 2,824575003 2,97282-003 3,04786-003 1,77769m003 2,42718m003 3,01690m003 1,16997-005 2,97282=003 2.82457=003 2,32478w003 1,98555-003 1,16997=003 2,59872-004 1,77769-003 1,08443=003 2,39872e004 3,047866003 2,00677=003 1,59743=003 3, 53234 m 103 .29173m003 4,51699-003 4.63102e003 7,13710=004 . 53102=003 3,96080=003 3,64469w004 1,08443,003 3,64469=004 GROUP PROB, NO. HESH 7 HS3H7 450000 450

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PROB	, NO. 2.	00000+000 ADCANDID	K-CALC (NO ACC) , RZ, TWU GROUPS, 20 X 27 MESH	PAGE NO.	94
			FLUXES IN CHANNEL NO. 23 ARE		
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2.00000.080 2DCANDID K.CALC (NO ACC) , RZ, TWO GROUPS, 20 X 27 MESH

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			FLUXES IN CHANNEL NO. 27 ARE		
S	ROUP	and	TAE'SH Z		
	5354900	12514-00			
01	9445500	32310 a 90	2		
9	3,18768e005	1,03654=009	r 04		
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2:00000+080 ADCANDID K#CALC (NO ACC) , RZ, TWU GRUUPS, 20 X 27 MESH PHOB, NO.

I	E	MESH.	HESH	TOUT	HESH		ZMESH
000	94839+00	39877-00	32500=00	72788-00	.60846m00	46805900	
00	17480 = 00	16003-00	13808*00	10918-00	.07365m00	03188408	C
0000	92583 # 00	0162-00	36564=00	81827-00	.76002=00	69154800	מו
0000	62943m00	59639=00	54726*00	48258-00	.40305#00	30955#00	•
0000	26830-00	22722-00	16616*00	08576-00	98691=00	87070m00	in
000	82668#00	77859-00	70709-00	61296-00	.49722=00	36115#00	•
0000	29085-00	23692-00	15675 00	05120-00	92142=00	76885#00	7
3.005	4,64935=005	4,59092=005	4,50405=005	4,38969-005	C	083	3 0
000	89838#00	33187-00	74045=00	62008-00	.47208#00	29808#08	0-
000-7	01691-00	95386=00	86012000	73671-00	58498=00	40658#00	10
0007	01691-00	95386=00	86012-00	73671-00	.58498#00	4065880B	다
0000	89838-60	83187-00	74045=00	.62008-00	.47208#00	.29808808	4
3000	64935-00	59092-00	50405-00	38969-00	24906.00	08374500	13
00-9	.29085-00	23692-00	15675-00	.05120-00	.92142m00	76885400	*
7.00	82668#90	77859-00	70709-00	,61296-00	49722=00	36115400	1.5
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02.00	62948-00	59639-00	5472600	48258*00	.40305*00	38955#00	17
00.	.92583×40	90162=00	86564 00	.81827 - 00	.76002=00	,69154#08	9
22 a 0 0	17480-00	16003-00	13808-00	10918-00	07365=00	.03188800	49
3400	94838=00	00-2186	82500=00	,72788-00	,60846¤00	46805408	20
(6)	RMESH	HESH A	KESH 1	NE S	2 5 2	HESH	ZMESH
25 e 0	13090*00	93810.00	.73227 . 00	51613-00	.60934=00	54536m0W	rt
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0.09	.52709=00	43306=00	33266=00	22724-00	.84952m00	.5374960U	ю
13.0	.08502=00	95663-00	81956#00	67561-00	.07174-00	.02913800	*
4200	59161=00	43202-00	.26165*00	.08273-00	.33213=00	.27918#08	ī.
28 = 0	.03439=00	84753-00	.64805=00	43857-00	.55973=00	49773808	9
1.8.0	40244900	19293-00	96925 00	,73435-00	.74892m00	.67939#0B	7
57.0	68672 00	45970=00	.21733-00	96281-00	.89504m00	81971#08	30
0300	.88022 # 00	64129-00	38619-00	11832-00	.99450 m00	91522400	0
53.0	3,97818 \$005	3,73321-005	3,47168=005	3,19704=005	2,04485m005	1,963576005	10
300	97818-00	,73321-00	47168-00	19704-00	.04485e00	96357#00	ef ef
0380	,88022-00	64129-00	.38619=00	11832-00	.99450m00	.91522800	CV
5700	,68672rg0	45970=00	,21733-00	96281-00	.89504m00	.81971#00	193 Ti
180	40244=00	19293=00	.96925=00	,73435.00	.74892m00	67939600	y el
280	.03439-00	84753-00	64805+00	43857-00	.55973=00	49773800	15
4. Cr	159161-00	43202-00	*26165 m00	.08273-00	.33213=00	,27918#0B	9
3.0	.08502*00	95663=00	81956-00	.67561.00	.07174w00	,02913#0B	17
0.09	.52709=0D	.43306-00	,33266*00	,22724-00	.84952m00	,53749800	3 0
3000	31560=00	,74195-00	12954.00	.48642=00	.78838m00	9803#08	64
200	000	C C T C	240040	E444 AA	40 4 400 4	1 1 2 2 2 4 1	•

2,00000+090 RDCANDID K*CALC (NO ACC) , RZ, TWU GROUPS, 20 X 27 MESH PROB, NO.

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ESE	I	NESE A	MESH 4	MEST	KEST	HESH	HESH	ZMESH
	48261.00	69140500	.10070-00	.20527 mg	73726*0	1237000	9480	
	1133-00	7 8 0	7499=0	384 0	0951-0	29418=0	1350580	N
	23143000	.24981=00	36864-00	.51436*0	31059-0	.52358*0	0057180	כיו
	.87346°00	12639000	33010-00	.79835=D	15478-0	.08023.0	37315#0	4
	.22724 00	40007-00	.11107-0D	.96419"0	92128=0	58565#0	7067880	50
	43691-00	63927=90	.06677-00	.98317*0	59123-0	.02741*0	9983880	•
	.61120-00	83810-00	19616"00	.83020*0	14813-0	39462#0	2407830	7
	14582m00	00-89166	.29610-00	4844200	.57827-0	67825#0	4280090	00
	83745-00	.09621s00	.36413-00	,92973"0	87104-0	.87130*0	.5554340	0
	88384 . 00	14913400	39857=00	.15516=0	.01926-0	.96903m0	6199490	0
	+8384-00	14913*00	39857-00	.15516-0	.01926-0	.96903#0	61994#0	11
	,83745 = 00	,09621.00	36413-00	.92973 0	87104-0	.87130=0	5554360	CVI wil
	,74582 000	.99168 .00	29610=00	.48442=D	.57827=0	67825 m	4280080	13
	.61120m00	83810 . 00	19616-00	.83020m0	14813-0	.39462=0	24078#0	*
	43691 a 0 0	63927=90	.06677-00	.98317=0	59123-0	.02741mG	99838#0	in.
	.22724000	. 40007m00	.11107-00	.96419=0	.92128=0	.58565#B	7867880	16
	.87346m00	12639-00	33010-00	.79835=0	.1547800	.0802300	3731560	17
	,23143e00	24981-00	36864-00	.51436=0	.31059-0	52358#0	.0857160	18
	.41133e00	.03257.00	27499-00	.14384=0	40951-0	.29418s0	1350580	1.9
	8261.00	69140-00	.10070-00	.20527	4,73726-008	23	619	20
I	N. H.S	ESH	Z Z	(2)	RMESH 26	RMESH 27	ZMESH	
	. 55842=00	.87969m00	.68754=00	46503=00	84417-00	.65701=01	-	
	04181.00	64204 90	,69226-00	.03098-00	.48711 .00	,68317w00	N	
	,62567000	33106.00	.77409-00	.6900600	.99493-00	.75920 . 0 D	P	
	.04639.00	,91843m00	,78762=00	0753*00	,22813-00	.76729#00	4	
	.12444e00	.35020=00	,70788m00	.86819"00	,52652=00	,68261w00	E.	
	.31655-00	.60597-00	.51222-00	35822 - 00	,78733.00	48263000	9	
	47624000	64984=90	.18083-00	.7655500	.00412-00	14765-00	^	
	.59958.00	.04964m90	,69725.00	.08017-00	17157-00	.66130m00	œ	
	+ 68353 m 00	10040 m 90	.04876-00	29432-00	28555-00	.01092.00	0	
	,72603-00	12827=00	,22671=00	.40273=00	,34324 00	.18791 * 00		
	.72603-00	12827 # 90	.22671-00	.40273-00	.34324-00	18791 = 00		
	, 68353 w00	10049r00	.04876-00	.29432=00	28555-00	01092-00		
	. 59958e00	04561-00	6972500	.08017-00	,17157-00	.66130 .00		
	47624000	. 64984m DO	18083-00	.76555-00	.00412-00	.14765m00		
	. 31655000	60597=80	.51222-00	35822 00	,78733-00	48263=00		
	.12444000	35020=00	.70788-00	.86819 m 0 0	,52652=00	.68261=00		
	.04639-00	91343mg0	,78762-00	.30753=00	.22813-00	.76729=00		
	6,62567.008	4,33106mp08	2,77409=008	1,69006#008	40	2,75920#009	18	
ď	044 84 00	10001	** > > > > > > > > > > > > > > > > > >	44.000	44			
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2.00000+000 2DCANDID K+CALC (NO ACC) , RZ, TWU GROUPS, 20 X 27 MESH

	PAGE NO. 54		PAGE NO. 55							
GROUP 1 2	(NO ACC), RZ, TWU GROUPS, 20 X 27 MESH P.		ACC) , RZ, TWO GROUPS, 20 X 27 MESH BY POINT AND GROUP							
REGION 3 6,58017+001 1,42225+001	K CALC		NDID K-CALC (NO BOUNDARY FLUXES	ZMESH 1	Ø 10 4	n o r	80 C	+ C F	4 N 0	F 80
REGION 2 1,38243+001 1,40211+000	2.00000+000 2DCANDID DISTRIBUTION BY REGION POWER	5303 = 001 5609 = 001 1885 = 002	000+000 2DCA	1239±00	400	041	0 44	440	2,73029e004 2,43494=004 2,07964=004	004
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2:00000+080 ADCANDID K-CALC (NO ACC) , RZ, TWO GROUPS, 20 X 27

50 PAGE NO. KaCALC (NO ACC) , RZ, TWU GROUPS, 20 X 27 MESH CHANNEL NO. CURBENT ON LEFT INTERFACE OF CHANNEL NO. ELUX ON LEFT INTERFACE OF ZEESE ZMESH NW 4 500 00 0 RDCANDID 45 8440 in or GROUP 2,16496=105 6,44158=005 2,68311+004 2,54934=004 2,35273=004 2.03884.006 2.14056.006 2.19460.006 2.19460.006 2.14056.006 1,15022=006 8,42434=007 5,13904=007 2,098235004 1,05596=004 6,44158-005 2,16496-405 1,05596-004 1,44176-004 1,44176-004 1,87699-006 1,67395=006 1,42969=006 1,15022-006 1,42969#006 1,67895#006 1,72719-007 8,42434-007 1,87699-006 72719-007 5,13904-007 2,000000+080 1,90346e003 2,13434e003 2,41267e003 2,43405e003 2,43550e003 2,43405e003 2,31267e003 2,31267e003 1,40346e003 1,40346e003 1,40793e003 1,40793e003 1,96400=004 5,84363=004 9,57938=004 "2,16410"004 "2,16410"004 1,96400-004 1,507930003 1,62571-003 -1,27518-004 ₩3.24014m004 *2,53383*004 -1.74107-004 e7,77889e005 *2,61442*005 =2,61442=005 =7,77889=005 -1,27518e004 =2,84117=004 -3,07856-004 .3.24014.004 -3,52194s004 *3,52194*004 *3,07856*004 *2,84117*004 -2,16410 0004 GROUP ZMESH 4004007 ZMESH 100 0 0 0 C 5 4

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2.00000+080 RDCANDID K+CALC (NO ACC) , RZ, TWU GROUPS, 20 X 27 MESH PROB, NO.

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64 PAGE NO. TOUT 23 × ADCANDID KACALC (NO ACC) , RZ, TWO GROUPS, 20 TOP BOUNDARY CURRENTS BY POINT AND GROUP RMESH U 10 10 10 10 10 800 GROUP 2 *1,74766*005 *1,73949*005 *1,72845*005 11,670461005 11,638751005 11,597761005 11,558781005 11,521801005 *1,10384*006 *7,21456*007 -1,47029-005 -1,46709-905 -1,71620-905 *1,73684+905 *8,62384=906 "6,62615"006 93,44883#906 *2,39767#\$06 *1,64185-006 84,45724m007 *2,39827s907 "7,37266m908 *1,72617#005 04,84787±906 2,000000+080 #5,11673#004 #4,98263#004 #4,81719#004 *2,28790*005 *1,30576*005 *9,89767 =4:13181=006 =2:51500=006 =1:33779=006 *3,01881#004 *1,26913#004 *6,18806*005 *4,102550007 *5,28660*004 -3,25780m004 -3,32765w004 .3,48320 0005 -5,32088-004 *4,62181*004 *4,39811s004 04,14789e004 *3,87309a004 -3,17187s004 0 N GROUP PROB, RESE 4500

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3 m3,33177m006 m5,98749m007 1 4 m3,07486m006 m5,52980m007 1 5 m2,74224m006 m4,92805m007 1 6 m2,34209m006 m4,20895m007 1 7 m1,88428m006 m3,38622m007 1 8 m1,38006m006 m3,48018m007 1 9 m8,41862m07 m1,51292m007 1	4	-3,50664800	6,30175-00	2 4
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GROUP	LEFT	BOTTOM	10.	RIGHT		
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a.	PKOB, NO, 2,	2,00000+080 RDCANDID	D K+CAL	C (NO ACC), RZ, TWU GROUPS, 20 X 27 MESH SIDE OF EACH REGION BY GROUP	PAGE NO.	89
GROUP	REGION 1	REGI	REGION	GROUP		
CI	000+00000.0	3,90025-003		5		
a.	PROB, NO, 2,	2,00000+000 ZDCANDID	ADID KECALC (NO	, RZ, TWU GF	PAGE NO.	69
GROUP	*2,24058*001 *8,53003*003	LEAKAGE REGION 2 *4,25138-002 *2,31644-003	ON BOTTOM SIDE REGION 3 *1.11964-001	GROUP 1		
er.	PROB, NO. 2.	2.00000+090 RDCANDID	K*CAL	C (NO ACC) , RZ, TWO GROUPS, 20 X 27 MESH SIDE OF EACH REGION BY GROUP	PAGE NO.	7.0
GROUP A	5,90376-001 -3,90025-003	REGION 2 6.25926=001 -5,20816=003	REGION 3 2,15386-002 3,87069-003	GROUP 1 2		
G.	PROB, NO, 2;	2.00000+000 ADCANDID	NDID K#CALC (NO	ACC), RZ, TWO GROUPS, 20 X 27 MESH OF EACH REGION BY GROUP	PAGE NO.	7.5
GROUP	2.24058+001	REGION 2	REGION 3	GROUP 1		

PAGE NO.

2:00000+000 2DCANDID K=CALC (NO ACC) , RZ, TWU GROUPS, 20 X 27 MESH

PROB. NO.

PAGE NO.

2,000000+000 2DCANDID K=CALC (NO ACC) , RZ, TWO GROUPS, 20 X 27 MESH

PROB. NO.

		:			
	73	7.			
	PAGE NO.	PAGE NO.			
	27 MESH	27 MESH			
	20 x	* 20 X			
	GROUPS	GROUPS			
	RZ, TWU GROUPS,	RZ, TWU GROUP			
GROUP 2	AND SOUP	A ND	, -1 or	000	1000 1000 1000
387-001 385-003	CALC (NO AC BY REGION N 3 GH 19-001 1	D &	36-003	4,0000+000 8,0345+001 8,0345+001 5,8359+00 0,0000+000	444099
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2 2 5 0	ABSORP ABSORP A 2 R 5-001 5-002 1.	EDCANDID	1-001	SCAT 0+0000+000 0+0000+000 1+6242+000 0+0000+000	+ + + + + + + + + + + + + + + + + + +
REGION 2 3,55499-002 1,30791-003	2,00000+000 2DC 1 REGION 2 00 1,34655+001	080+000	2,54972-001 2,00211-003	04400	NO
4 40 0	2,000 1,00	0	200	1X1 1 2 2 8 8 3	425004
REGION 90376-0	NO. EGIUN U5884+0	. NO.	7,04662-002	I 40404	01 4 01 4 01 4 01 4 01 4 01 4 01 4 01 4
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- 1		SCATTER		SOURCE BY REGION	IN AND GROUP				
-	REGION 1	REG10N 2		REGION 3	GROUP				
1	1.50493=001	2,38891+002		2,17422*001	co.				
ō	PROB, NO, 2.01	.00000+000 ZDC	RDCANDID K	CALC	(NO ACC) , RZ, TWU	GROUPS, 20 X 2	27 MESH	PAGE NO.	76
	REGION 1	REGION 2		2 2				:	:
	1,66381=002	4,61164-003		2,67448*002	- 03				
ō	PKOB, NO, 2,0	2,00000+000 &DC	RDCANDID B	KTCALC (NO	ACC) , RZ, TWU BY REGION	GROUPS, 20 X 2	27 MESH	PAGE NO.	77
REGION	LEAKAGE	ABSORRTION	NO	BUCKLING	FISSION	SCATTER	EXTERNAL	TOTAL	REGION
	5,86476*001 3,42420*002 *5,95308*001	01 1,19306+000 02 1,57242+001 01 7,66147-001	000	0,00000+000	2,11346+000 2,56974=001 2,04061=001	1,30493"001 2,38891"002 2,17422"001	0.0000000000000000000000000000000000000	4.64416=001 8,93789=002 2,50649=001	4 01 15
	2,54093*002	02 2,11645+000	000	0 * 0 0 0 0 0 0 + 0 0 0	2,57449+000	3,71804=001	0.00.000.0000.0	8,04439"001	ALL
0	PROB; NO, 2,0	2.00000+000 2DC	2DCANDID	KaCALC (NO	, RZ, TWU	GRUUPS, 20 X 3	27 MESH	PAGE NO.	78
	LEAKAGE	ABSORMTION	NO	BUCKLING	FISSION	SCATTER	EXTERNAL	TOTAL	GROUP
	2,15586*002 3,87069*003	02 1 _e 77645+000 03 3 _e 39997-001	000	0,000000+0000	2,55443+000 2,00581*002	3.71804 = 001	0.0000000000000000000000000000000000000	7,56444=001 4,79946=002	નળ
	S KANGTOND	000 4 4 4 4 5 4 0 0 0							

2. Sample Problem 2

- a. Description

 - (2) <u>Configuration</u>

 Same as Sample Problem 1
 - (3) Convergence Criteria

 Same as Sample Problem 1

PR08. NO.

23

PAGE NO.

3.00000+000 2DCANNID K-CALC (W/ACC), RZ, TWO GROUPS, 20 X 27 MESH

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* *	00+/609//	/ 0 - 0 0	1/05+05	Z40occo.	.o. } }	.) □ •	.032/303700
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8	1.3007313+000	4.3227500-001	1.1166810+000	6.7356246-002	1.00000000+380	1.00000000+000	6.8440604-001
4	.2770434+00	.5538333-00	.0741421+00	.4212842-00	0040000000.	.00000000.	.1875873-00
יטו	.2593540+00	.3341824.n	,0527936+00	.6239459-00	25+6000000.	0040000000.	.1937532-00
40	.2460554+00	.9264789-00	.0403481+00	.0566803-00	0040000000	000000000.	.4770019-00
7	.2359993+00	4214974-n	.0331508+00	.6130244=00	00+0000000.	00+00000000	.9100104-00
60	.2283286+00	.8503622-n	0317257+00	.2465991-00	00+00000004	.00+0000000.	.4668948-00
0	.2223844+00	.2203109-0	.0303336+00	.9337948-58	0040000000	00000000.	.0830246-00
10	.2176492+00	.5360629-0	.0289328+00	.6610098-03	0040000000	.00000000.	.7532654-00
11	.2137044+n0	.8061696-n	.0275521+00	.4196198-00	00.000000.	00+00000000	.4693513-00
12	.2101739+00	.9155540-0	.0256558+03	.2037956-00	.0000000±00	.00000000.	.3410042-00
13.	.2067851+00	.9699044-n	.0232527+03	+0082774-00	004000000.	.00000000.	.2626031-00
4	.2033687+00	.0181376-n	.0205935+00	.8305807-38	0040000000	.00000000.	.1877973-00
45	.1998310+00	.0536106-0	.0179032+03	.6696202-08	00+0000000.	00000000.	.1254212-00
91	1961292#00	.0817362-00	.0153233+00	.5486144=00	.1207133-00	.0157843+00	.0714973-00
17	.1922483+00	.1108706-0	.0128980+00	.6070502-00	.1207133-00	.1541610+00	.0181095-00
41 60	.1881516+00	.1453672-00	.0110641+00	.9015912-00	.1207133-20	.5105853+00	.6527356-00
19.	.1837515±00	.1934397 mn	.0098562+00	.5999353=00	.1207133-00	.3477837+00	.0462262-00
20	.1788621+00	.2761269-00	.0080144+00	.1376828-00	.1207133-00	.5147047+00	.04n1710-00
21	.1731275+00	.4684469-n	.0050043+00	.4633023-00	.1207133-00	.6648508+00	.8159641-00
22	.1666187+0.0	.7340278-0	.0006331+00	.5482408=00	.120/133-00	00+0000000.	.7197333-00
23	.1604084+00	.7434306-0	.0006125+00	.3284806-00	.120/133-10	000000000.	.6269419-00
24	.1544774+00	.7525706-0	.0005489+00	.1226426-00	.1207133-00	.00000000.	.5291833-00
52	.1488063+00	.7614630-n	.0004689+00	.9924957-00	.120/133-00	.0157843+00	.432261.1-00
56	.1433797+00	.7701327-0	.0003749+00	.3274390-00	.120/133-00	1541610+00	.3361611-00
27	.1381862+00	.7792441-0	.0002138+00	.3930639-00	.120/133-09	.5105853+00	.2289375-00
28	1332186+00	7908584-0	.9991021+0C	.9395698-00	.1207133-00	3477837+90	.0824177-00
58	.1284732+00	.8074002-0	.9933376-00	.4301811-00	.120/133-00	.5147047+00	.8593736-00
30	.1239474+00	.8385790-0	.9823993-00	.8053446-00	.1207133-08	.6648508+00	.4382030-00
31	.1196198+00	.9032313-0	.9695521-00	*2923073=00	.1207133-00	00+00000000	.6320769-00
32	1154387+00	.9066249-0	.9696029-00	.5092589-00	.1207133-00	000000000	.2978076-00
33	.1114055+00	.9098670-0	.9719377-00	.4197259-68	.120/133-00	.00000000.	.2070737-00
34	.1075157+00	.9129685-0	.9738810-00	.3330860=00	-120Z133±00	00+0000000.	.0912480-00
35	.103762A+00	.9159428-n	.9755294-00	.2880641-00	.6419436+0C	.0167014+00	.9586676-00
36	.1001418+00	.9188456-n	.9769051-00	.5277702-60	.6419436-00	1644195+00	.8059492-00
37	.0966479+00	.9220279-0	.9778413-00	.2564947-60	+6419436-00	.55598/3+00	.5813413-00
35	.0932759+00	.9260648+0	.9780507-00	.1256609-00	.6419436-00	.54370B0+00	.1985938-00

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4,617							
	,0022946+00	9943272+00	9960833-00	.0617798-00	00+0000000.	.000000000.	.7561878-00
100	.0018723+00	9943950-00	9962746-00	.9879536-00	.00000000.	.000000000	.8796067-00
4	,0014652+00	9944653-00	9964386-00	.9168714-00	.0000000	.00000000.	.9732976-00
S	,0010729+00	9945387=00	9965840-00	.8485222-00	.00000000	.00000000.	.0453440-00
40	.0006947+00	9946147-00	9967139-00	.8126731-00	.6434338-00	.0167040+00	.0991643-00
7	,0003300+00	9946943-00	.9968263-00	.0027757-00	.6434338-00	,1644491+00	.1320526-00
00	,9997829-00	9947849-00	9968943-00	.5825787-00	.6434338-00	,5561211+00	.1093128-00
0	.9963879-00	9949039-00	9968816-00	.0769378-00	.6434338-00	.5443150+00	.9777182-00
10	.9931029-00	.9950926-00	9969458-00	.7630108-00	.6434338-00	,6535044+00	.8531368-00
11	.9898930-00	,9954871+00	.9970703+00	.9019440-00	.6434338-00	,9198935+00	.5831660-00
12	,9865933-00	,9955335+00	.9972314=00	.5407437-00	.6434338-00	.00000000.	.6979383-00
13	.9834006-00	,9959041+00	.9972048*00	.4896361-00	.6434338±00	00+00000000	.30n6896-00
14	,9803124-00	,9962259-00	,9971759-00	,4403271-00	.6434338-00	00+0000000*	.50n4871-00
t. ←	.9773263-00	.9965051+00	.9971568-00	.3924930-00	.6434338-00	.00000000.	.5169181-00
16	,9744401-00	.9965890+00	.9972467-00	.3684635-00	.6678942-00	.0167471+00	.5770029-00
17	,9716512-00	.9966181+00	.9974419-00	.5151471=00	+6678942-00	1649350+00	.2373968-00
18	,9689567-00	.9966544-00	,9975918-00	.9585058+00	.6678942-00	,5583192+00	.3747876-00
19	,9663532-00	9967051+00	.9976797+00	.1022030+00	*6678942-00	.5543206+00	.7467011-00
20	,9638375-00	.9967869-00	.9976648"00	.7120668#00	.6678942-00	.7210328+00	.7790409-00
21	,9614081-00	.9969649-00	.9976653-00	.2780028-00	.6678942-00	,0128072+00	.0037626-00
22	,9589896-00	,9966111-00	,9980156=00	.1219412=00	.6678942-00	.00000000.	.4045429-00
23	.9566555-00	.9969186-00	,9979135-00	.0819777-00	.6678942-00	.00000000.	.9490033-00
24	,9544082-00	9971798-00	.9979162-00	.0415257-00	,6678942-00	.00000000.	.3639254-00
25	,9522362-00	,9974026-00	.9979475-00	.0064396-00	.6678942-00	00+0000000*	.4490942-00
56	.9501357-00	.9975181-00	,9979883=00	.8952543+00	.6631278-00	,0167387+00	.7020701-00
27	.9481050-00	.9975387=00	,9981147-00	.0960429-00	,6631278-00	1648403+00	.7596320-00
28	.9461421-00	,9975645-00	,9982302-00	.4171780-00	.6631278-00	.5578904+00	.6572014-00
58	. 9442444-00	.9976008-00	,9983332-00	.2448330-00	.6631278-00	.5523647+00	.9936737-00
30	9424100-00	.9976594-00	.9982950-00	.8519520-00	.6631278-00	,7077478+00	.3561805-00
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20	,9388768-00	.9975290-00	.9985607-00	.1356782-00	.6631278-00	.00000000.	.0316767-00
33	,9371766-00	.9977521-00	.9984770=00	.8442643-00	,6631278-00	*00000000	.2493276=00
4 1	9355405-00	.9979419-00	.9984797-00	.5471563-00	,6631278-00	000000000	.3781740+00
35	,9339594-00	.9981040=00	.9985027-00	*5918670-00	,6631278-00	*0000000	.9871520-00
30	,9324305-00	.9981857-00	9985331-00	.1690123-00	.6617411-00	.0167363+00	.4743789-00
15	, 9509524-00	.9982009-00	,9986252=00	*9403604=00	.661/411-00	1648128+00	,2231346±0U
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A 10	9281424-00	998246/-00	004/86/866	.6200047-00	• 661/411 = UU	1040047166	.1261/43*00
40	,92680/2-00	9982899*00	9987551-00	.5133072-00	. 661/411-00	, /038945+00	.6524117-00
4	,9255178-00	.9983840 .00	.9987553=00	,1822386±00	.6617411+00	,9885983+00	.7134247=00

PAGE NO.

3,00000+000 2DCANDID K-CALC (W/ACC), RZ, TWO GROUPS, 20 X 27 MESH

PROB. NO.

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          3,2989899=005
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                               2,4008606+004
                                          1,2009241-005
                                                   1,1580883-005
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                                                                                              1,1724156-005
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.0827735=005
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                                                                                                                                                            9,9998337*001
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          9,9997447-001
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                                                    9,9997718+001
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          ,9996389*001
                     9,9996479"001
                              .9996674=001
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                                                    9,9996629#001
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                                                                        .9997157-001
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3,4814672*003
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                    9,8935571-001
                              9,8932928-001
                                          9,8930299-001
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7 644 L41 FLUXE 8.6834609400 9 641 L81 FLUXE 8.7444447400 10 641 L81 FLUXE 8.46681467400 10 641 L81 FLUXE 1.68681467100 2 641 L81 FLUXE 1.6864834710 3 641 L81 FLUXE 4.6864834710 5 641 L81 FLUXE 4.487489800 5 641 L81 FLUXE 7.4874898710	9	-4	wi.	wi	3694657=00					
6 6 8 1	7 6	v i	7	00	8334609+00					
9 9 9 1 1 FLUX	9	-1	-i	Rυ	744137900					
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2 GR1 JR1 FLUXE 1,6844389#00 2 GR1 JR1 FLUXE 1,0677709#00 3 GR1 JR1 FLUXE 6,9679182#00 4 GR1 JR1 FLUXE 4,48743087#00 5 GR1 JR1 FLUXE 4,48743087#00 5 GR1 JR1 FLUXE 4,48743087#00 5 GR1 JR1 FLUXE 4,48743087#00	0	3	-		4682142-00					
2 6 8 4 1 8 1 FLUX	9	7	4	wi	6244349+00					
4 Get Jet FLUXE 6,9079122=00 4 Get Jet FLUXE 4,4974303>00 5 Get Jet FLUXE 2,7432087=00 6 Get Jet FLUXE 2,7432087=00	2	7	+1	*	067770900					
6 Get Jet FLUXE 4,4974303500 5 Get Jet FLUXE 2,733087=00 6 Get Jet FLUXE 1,4432172=00	3	7	-	9	9679172-00					
5 Get Jet FLUXE 2,7432087-00 6 Get Jet FLUXE 1,4432172-00	4	-	-		4974303+00					
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3. Sample Problem 3

a. Description

(1) Problem Type

Dimension search varying the z direction uniformly to obtain $k_{\mbox{eff}}$ = 1.025.

(2) Configuration

Same as Sample Problem 1.

(3) Convergence Criteria

 k_{eff} difference = 10^{-5} .

 k_{eff} bounds = 10⁻⁵.

Sum of flux differences = 10^{-5} .

 $|k_{eff} - 1.025| \le 5 \times 10^{-5}$.

CODE DEPENDENT DATA OR INPUT IF DIFFERENT THAN ZERO.

BUCKLING DATA HAVE NO! BEEN INPUT.

VOLUME SOURCE DATA HAVE NO! BEEN INPUT.

ALL FLUX VALUES WILL BE PRINTED.

The state of the s	3
ENERGY GROUPS TO USE IN PROBLEM	N
NUMBER OF REGIONS	8
	10000
MAXIMUM NUMBER OF INNER ITERATIONS / CHANNEL CALC	4
XIMUM NUMBER OF EXTRAPOLATIONS ON K-EFFECTIVE	4
XIMUM NUMBER OF INTERPOLATIONS ON K-EFFECTIVE	O
PRINT FREGUENCY FOR CRITICALITY SEARCH,	₩
•	1000001

1.0250+00	,				
DESIRED K-EFFECTIVE	SECOND GUESS ON SEARCH PARAMETER	CONVERGENCE CRITERION ON OUTER BOUNDS OF K	SMALLEST FLUX VALUE USED IN COMPUTING K-EFFECTIVE	FIRST GUESS AT K-EFFECTIVE	CONVERGENCE CRITERION FOR SUM OF FLUX DIFFERENCES

PAGE NO.	
DIMENSION SEARCH, RZ, TWO GROUPS, 20 X 27 MESH	PROBLEM TYPE
2DCAND1D	
1.00000+000	
PR08, NO.	

CRITICALITY SEARCH VARYING THE Y-DIMENSION UNIFORMLY

FRACTIONAL MODIFIER.......... 5.0000-002 CONVERGENCE CRITERION (EPSILON)..... 5.0000-005

FLUX CONVERGENCE CRITERION (DEL!A)... 1.0000-005

13	
PAGE NO.	
TWO GROUPS, 20 X 27 MESH	0
N SEARCH, RZ, TWO G	II LNIN
DIMENSION S	NXT = 0
.00000+000 2DCANDID	1,444000+000
+	LUE (X) =
PROB, NO.	CONTROL VALUE (X)

PAGE NO.	
DIMENSION SEARCH, RZ, TWO GROUPS, 20 X 27 MESH	REACTOR GEOMETRY
2DCANDID	
1.00000+000	
PROB. NO.	

N

TOP	6,0000+001
BOTTOM	0.00+0000.0
RIGHT BOUNE	7,8100*001
LEFT BOUND	0.00+0000+0
GEOMETRY TYPE	RZ WITHOUT INVERSION

REGION BOUNDARIES

TOP BDRY.	6.0000+001 6.0000+001 6.0000+001
BOTTOM BDRY.	0.000+0000.0
RIGHT BDRY,	2,3168+001 2,5902+001 7,8100+001
LEFT BORY.	0,0000+000 2,3168+001 2,5902+001
REGION NO.	нам

Ŋ

PAGE NO.

1.00000+000 2DCANDID DIMENSION SEARCH, RZ, TWO GROUPS, 20 X 27 MESH

MESH LEFINITION

INCREMENT METHCD - X(R) DIRECTION

X(R) INCREMENT INCREMENTS X(R) X(R) INCREMENTS X(R) INCREMENTS

X(R)

7.8100+001 2,5902+001 2,3168+001 0000+0000+0

IMAX = 2

X(R) PLUS DELTA X(R) FROM LEFT 10 RIGHT OF REACTOR

MESH NO. 1.9307+000 3.8613+000 5,7920+000 7,7227+00 9,6533+000 1.1584+901 1.3515+001 1,5445+001 1.7376+3€1 1.9307+001 11 12 12 13 14 15 14079+001 2,4079+001 2,5902+001 3,0252+001 3,4602+001 3,8951+001 4,3361+001 4,7651+001 21 22 23 24 25 5,2001+001 5,6351+001 6,0701+001 6,5050+001 6,9400+001 7,3750+001 7.8100+001 MESH NO. ABSCISSA

1.00000+000 2DCANDID DIMENSION SEARCH, RZ, TWO GROUPS, 20 X 27 MESH PROB. NO.

MESH LEFINITION

4

PAGE NO.

INCREMENT METHCD - Y(Z) DIRECTION

Y(Z) Y(Z) INCREMENTS Y(Z) INCREMENTS Y(Z) INCREMEN'S 6,0000+001 INCREMENTS 20 Y(Z) 0.00000.0 20

Y(Z) PLUS DELTA Y(Z) FRCM BOTTOM TO TOP OF REACTOR

Y(Z) MESH NO. 11 12 12 13 14 14 5000+001 4.5000+001 5.1000+001 5.1000+001 5.1000+001 5.7000+501 6.0000+001

23 PAGE NO. THE LAST Y(Z) COORDINATES USED IN THIS PROBLEM ARE GIVEN BELOW. NOTE - Y(Z) BOTTOM = 0.0000+000 1.00000+000 2DCANDID DIMENSION SEARCH, RZ, TWO GROUPS, 20 X 27 MESH

MESH NO. 11 12 13 19 20 20 OKDINATE 3,5383+001 3,8599+001 4,8132+001 5,1466+001 5,4682+001 5,7899+001 6,1115+3,1 6,4332+001 MESH NO. 1 ORDINATE 3,2166+000 6.4332+000 9.6498+000 1.2866+001 1.6083+001 1.9300+001 2.2516+001 2.5733+301 2.8949+3.1 3.2166+001

CANDIDED IS PREPARING X SECTIONS AND COEFFICIENTS ON LUN 4

																																			•						
ERRLAM 3.0777590+000	0129531+00	6.9989977-001	3379057-00	2733140-00	5217598-00	9396533-00	47:8034-00	.0654141-0n	7251916-00	5024138-00	4048266-00	318.652-00	236, 354-00	11/2105/11	0487754-00	824333210	12.3878-00	0367111-00	.7842958-DP	,8639953-DP	7225297-00	.5947279-0n	4733338-00	.3555512-00	.2268238-00	.065.567-BR	.8405506-DR	4555185-01	00000000000000000000000000000000000000	2314630=00	8814219-00	5458013-00	1859350-00	.7448846-00	.1168519-00	.0704109-00	.422,516-00	3419140-06	2664864-0	19534/5-01	.12/1522-01
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CHANGE IN PHI 1.6001128-001	5238998=00	7.0901127-002	8765122-00	1152899-00	5551726+00	1034577-00	7201811-00	3851776-00	0867123-00	8169426-00	5693228-00	3392236=00	1257319+00	929529-00	001/2/0///	1406400-	8754356400	4115122-00	1825005-00	,0910585-00	,9002621-UD	7254182-00	,6043499-00	.7821125-U ₀	3947643-00	.8099436-00	9824661-00	2749399+00	001000000000000000000000000000000000000	480860	244810210	7161038-00	2487260-00	.1774291-00	,9725234-00	.1864231-00	,8115272*UC	.7210671-00	.6296316-00	5668051-00	,6834519*UC
K+UPPER 3.1440882+000	2502260+00	1.1342587+000	0899794+00	0673236+00	.0537236+00	.0469086+00	.0442636+00	.0417356+00	.0393630+00	.0371593+00	.0346438+00	.0317438+00	.0286574+00	.0255798+00	0726440+00	017074410	.0152860+00	0132723+00	.0101569+00	.0062390+00	.0057638+00	.0053407+00	.0049552+00	.0045972+00	.0042210+00	+0037695+00	,0032896+00	.0024926+00	.0012412+00	0.437.1100.	00+0540400	.0010044+00	.0009460+00	.0008740+00	,0007704+00	.0005958+00	.0003147+00	.0002996+00	.0002855+00	,0002723+00	.0002595+00
K-LOWER 6.6329144-002	3727288-00	4.3435897-001	5918884-00	3999222+00	0154760-00	5294330-00	9718329-00	3519415-00	6684379-00	8691789-00	9416112-00	9993725-00	0505383*00	085692/-00	11/400070	188311010	2408210=0	3290515-00	5231395-00	7759903-00	7853851-00	7939345-00	8022186-00	8104171-00	8195274-00	8311897-00	8488407=00	6813737-00	04444740	0488416#0	9517737=00	9545863-00	9576009-00	.9612909-00	,9665354-00	,9752537-00	,9889265-00	,9895768-00	.9901904-00	,9907693-00	.9913238-00
ATION HISTORY K-EFFECTIVE 9.7052596-001	5549969-00	9.4885777-001	4756525-00	4982784-00	5442377-00	6050041-00	6744533-00	7481008-00	8226053-00	8953977-00	9642263-00	0027509+00	0084387+00	0134558+00	01/8108+00	0245003+00	02693940	0283674+00	.0285543+00	.0276902+00	.0269160+00	.0262197+00	.0255914+00	,0250237+00	.0245134+00	,0240626+00	.0236785+00	0233741+00	000000000000000000000000000000000000000	0042284260	0226128+0	0224541+0	.0223072+00	.0221736+00	.0220566+00	.0219654+0[.0219167+00	.0218718+0	•0218310#D(0217942+0	•021/611+00
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2-003		-004		1	1		ŧ	E	ł	-004	1		- 8	,	-005		1	ŧ
1.054,832	9.6466494	8.377,079	6.2729736	2.9791650	2.8128290	2,6557595	2.5073742	2,365,046	2.2125212	2.0262715	1.7627073	1,3278832	6,6716078	6,338,809	6,3051767	5,6819583	5,3751632	5.0852017
1.5066170+11.	2.3323867+35	4.4385000+38	9.2713734+38.	1.00000000+000	1.0000000+100	1.00000000+300	1.0157058+00	1.1532902+080	1.5068170+000	2.3323867+33	4.4383000+30	9.2715734+00	1.0000000+30.	1.00000000+007	1.00000000+50	1.00000000+30.	1.0167885+00	1.1654018+00
		1	1	8	1		1	8	1.	3		9.0766404-001			1		1	1
2.0661518-004		4.9275447-064	0	0	0	0	0	0		0		2,2419544-004	1,7811546*005	1.7230024 - 005	1.6689240-005	1,6174166-005	1.5942094-005	1,7714563-005
1,0002458+000	1,0002288+000	1,0002042+000	1,0001625+000	1,0000948+000	1.0000909+000	1,0000872+000	1,0000837+000	1,000803+000	1.0000766+000	1,0000720+000	1.0000655+000	1.0000547+000	1,0000374+000	1,0000362+000	1,0000351+000	1,0000339+000	1,0000328+000	1,0000317+000
.9919	66.	.9936	.995352	6266.	960866	.998216	.998329	,9984375+0	.998553	.998693	.99889	9,9992194-001	1666.	.999728	-0057999.	.99977	.9997904-	.99980
00+9	.0217064+00	.0216866+00	.0216755+00	.02	.0216820+00	.0216861+00	,0216909+00	.0216962+00	.02170	.0217086+00	.0217163	61+00	.021.73	28+00	.0217658+00	17785+0	.0217910+00	.0218033+
45	4	47	48	49	50	51	52	53	4	S S	96	57	58	29	9	61	62	63

A CONTROL CHANGE WILL BE MADE USING KEFF= 1.0221076+000

	X Y YP YPP GAMMA RTO TEST 6.2000000+001 1.0217910+000 1,2368342-005 -2.4162000-007 1,0221076+000 5.1189232+001 6.3306552·054	
DELKNM1 1.2489152-005	RT0 5.1189232+001	
FNTH FAUNT TAUNTH TAUNY DELKNTH DELKNTH 5.0802017-005 5.3751632-005 -3,1967269-003 -3.2089744-003 1.224/532-005 1.2489152-005	1.0221076+000	
TAUNF1 +3.2089744*003	YPP -2.4162000=007	CNVGCRIT
TAUNTH -3,1967269-003	YP 1,2368342-005	ENT NUMBER THRU
FNM1 5.3751632-005	Y 1.0217910+000	TRACE BY STATEM
FNTH 5.0802017-005	X 6.2000000+001	LOGICAL BRANCH TRACE BY STATEMENT NUMBER THRU CNVGCRIT

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BRANCH NO	∓ 1	N	m	4	ī	œ	7	6 0	σ.	10	

PROB, NO. 1.00000+000 2DCANDID DIMENSION SEARCH, RZ, TWO GROUPS, 20 x 27 MESH

I LUIN

NXT =

CONTROL VALUE (X) = 1.500000+000

PAGE NO. 25

56 PAGE NO.

THE LAST Y(Z) COORDINATES USED IN THIS PROBLEM ARE GIVEN BELOW.

NOTE - Y(Z) BOTTOM = 0.0000+000

CANDID2D IS PREPARING X SECTIONS AND COEFFICIENTS ON LUN

7,6757802-005 5,1188574-005 3,9181410-005 3,3822871-005 2,4637615-005 1,3355870-005 1,8263160-004 ,1867128-004 8,8625733-005 6,6692010-005 5.8268284-005 4,9182936-004 3,0985600-004 2,7158402-004 2,39.5750-004 2.0946717-004 1.5858124-004 1,3729758-004 .0251999-DD4 4.5097549-005 3,1188887-005 57 PAGE NO. 2.0670786+331 1.00000000. 1.00000000+300 1.0000000.t 1.00000000+33 1.00000000+00001 1.00000000.1 1.0000000. 1.00000000+50 1.0167788+30 1.1652926+33 5.7717265+00 1.00000000+00 1.000000000 1.0000000. 2DCANDID DIMENSION SEARCH, RZ, TWO GROUPS, 20 X 27 MESH SIGMA 1.00000000+0000 1,00000000+000 1,00000000+000 1.00000000.1 1.0000000+000 1,00000000000 1.0000000000 1.00000000+000 1.00000000+0000 1.00000000+000 000+000000001 9.6858806-001 9.6858886-001 1,00000000.1 1.00000000+0000 9.6858866-001 100-90886896 9.6858806-001 9.6858806-001 9.6858806-001 1.000000000 CHANGE IN PHI 4,6467302-005 2.9592621-005 2,6040217-005 3,7874628-005 3,4989745-005 3,3771488-005 3,2640844-005 3,0562673-005 2,7757329-005 .6884718-005 .8476237-005 3,6948845-005 4,6195375-004 2,3890014-005 3,9759127+005 3,6333739-005 3,1575769-005 2.5647139-005 5,8887073-005 4,2315767 "005 1,0001507+000 1,000636+000 1,0000455+000 K+UPPER 1,0004925+000 1,0003186+000 1,0002512+000 1,0002221+000 1,0001956+000 1,0001718+000 1,0001159+000 1,0001020+000 1.0000710+000 1.0000571+000 1.0000510+000 1,0000379+000 1,0000308+000 1,0003655+000 1,0002829+000 1,0001321+000 1,0000000+000 1,0000798+000 1,0000430+000 1,000000004000 .0000133+000 .0000128*000 .0000124+000 .0000127+000 1,0000134+000 1,0000134+000 .0000134+000 .000134+000 .000131*000 .000120+000 .000#118*000 .,000117+000 .0000119+000 ..0000133+000 1,0000205+000 1,0000050+0000 .000+780000. 1,000121+000 1,0000132+000 1,0000113+000 1,0000130+000 1.00000+000 ,0225193+000 1.0225998+000 1.0221485+000 1.0221880+000 1,0222261+000 1,0222986+000 1.0223332+000 .0224906+000 1,0225738+000 1,0226250+000 1.0226494+000 .0226730+000 1.0226960+000 1.0227184+000 1.0227407+000 1.0227650+000 K-EFFECTIVE 1.0222630+000 1,0223667+000 1.0223992+000 .0224306+000 .0224611+000 1.0225470+000 ITERATION HISTORY 4624624 20 20 20 20 I TERATION 4,441

7.5347198-006 5.1484967-006 4.556167-006 6.7489163-006 6.7489163-006 6.0774910-006 1.0252930-005 7.2693510-006 7.2693510-006 7.3468823-006 7.3468823-006 7.3468823-006 7.3468823-006 7.3468823-006 7.3468823-006 7.3468823-006 7.3468823-006 7.3468823-006 7.3468823-006 7.3468823-006 7.3468823-006	.4967102-00 .2433775-00 .9017759-00
7	# + + + + + + + + + + + + + + + + + + +
689 689 689 689 689 689 689 689 689 689	6747436-0 6747436-0 6747436-0
2. 31413241005 2. 134133541005 2. 134133561005 3. 05599691005 4. 84729401005 1. 05181601005 1. 7498721005 1. 57494171005 1. 57494171005 1. 5503531005 1. 74958891005 1. 74958891005 1. 74958891005 2. 22717881005 3. 5366789105 7. 68460105	2903741-00 2453492-00 1993477-00
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1.00002004 1.000002011 1.000002074+000 1.0000110011001 1.000011684+000 1.000011684+000 1.000011684+000 1.000011684+000 1.000011684+000 1.000011684+000 1.000011684+000 1.00001284+000 1.00001284+000 1.00001284+000 1.00001284+000	.0000110+00 .0000110+00 .0000110+00
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0 4 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0	4444 4010410

A CONTROL CHANGE WILL BE MADE USING KEFF= 1.0233260+000

3.0251902+001 3.6405660.004 DELKNM1 1.2236414-005 GAMMA 1.0233260+000 DELKNTH 1.1838507-005 4.4000000+001 1.0231440+000 1.2037461-005 +3.9794757-007 3.9017759-006 -1,8441939-003 -1,8560324-003 FNTH 2.8935319-006

LOGICAL BRANCH TRACE BY STATEMENT NUMBER THRU CNVGCRIT

STATEMENT NO	m	35	40	4	55	80	105	0	+1	115	S
BRANCH NO	₩.f	C)	ю	4	ī.	•	7	œ	0		Ħ

H -Z Z CONTROL VALUE (X) = 1.576935+000 PAGE NO. 2DCANDID DIMENSION SEARCH, RZ, TWO GROUPS, 20 X 27 MESH 1.00000+000 PROB, NO.

50

THE LAST Y(Z) COORDINATES USED IN THIS PROBLEM ARE GIVEN BELOW. NOTE - Y(Z) BOTTOM = 0.0000+000

MESH NO. 1 6 7 8 9 10946+00 1 9 10946+00 1 1.2946+01 1.6183+00 1 1.9419+00 1 2.2656+00 1 2.5892+00 1 2.9129+0 1 3.2365+00 1

MESH NO. 11 12 12 13 14 14.8548+001 5.1785+001 5.5021+001 5.602+001 5.602+001 5.602+001 5.602+001 6.1494+501 6.4731+001

54 CANDIDED IS PREPARING X SECTIONS AND COEFFICIENTS ON LUN

· 🛶 ·	1.0233855+000	K-LOWER 1.0000320+000	K-UPPER 1.0006601+000	CHANGE IN PHI 6.4984712-005	SIGMA 1.0000000+000	1.00000000.1	ERRLAM 6.2813351-004
4. 4. 0.		-		000000000000000000000000000000000000000			0000
V 10	1.0234979+000	1.000396+000	1.0004867+000	5.5887757-005	1.000000+050		3.8641508-014
4	0235510400	000414400	0004780+00	3186870-00	00+0000000	C + 0 0 0 0 0 0 .	3659299-88
. rv	*0236023+00	.0000424+000	.0003348+00	0916839-00	00+0500000	000000000000000000000000000000000000000	.9239207-0n
9	.0236517+00	.0000423+00	.0002951+00	.8897568-00	000000000	35+60000000	.5276220-0n
7	.0236995+00	.0000415+00	.0002588+00	.7042814-00	.00000000.	.0000000.	.172.785-00
00	.0237455+00	.0000408+00	.0002260+00	.5306940-00	.00000000.	C+8000000°	.8527388-DP
0	.0237899+00	.0000395+00	.0001970+00	.3663578-00	00+0000000.	5+6000000°	.5754232-0P
10	.0238327+00	.0000383+00	.0001715+00	,2096272-00	000000000000000000000000000000000000000	0000000°	.3324924-DP
T T	.0238741+00	.0000370+00	.0001494+00	00283886-00	00+00000000	C+0000000°	.124, 717-0n
12	.0239139+00	,0000357+00	.0001303+00	.9147924-00	000000000	·00000000·	.466 500-0n
13	.0239524+00	,0000345+00	.0001140+00	.7751532-00	.00000000.	·00000000·	.9545483-0P
∜ 1	.0239894+00	,0000332+00	.0001001+00	.6401241-UO	.00.000000.	· 0 0 0 0 0 0 0 0 •	.69. 422-DD
 	.0240251+00	.0000320+00	.0000883+00	.5096319-00	00+0000000	35+0000000	.6288060-0P
10	0240596+00	000000000000000000000000000000000000000	00+28/0000	4402910+00	.64151/3-00	.016/20/+35	000000000000000000000000000000000000000
1	,0240928+00	.000029/+00	.0000696+00	. /990349-U0	.6415173-00	. 1.046111+00	9899198-00
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22	.0242440+00	.0000247+00	.0000380+00	9236675-00	.6415173-00	00000000	3317476-00
23	.0242729+00	.0000252*00	.0000351+00	,8341350-00	.6415173-00	000000000000000000000000000000000000000	.8722812-0n
24	.0243009+00	.0000254+00	,0000325+00	.7450885-00	.6415173-00	.00000000.	.0921378-00
25	,0243280+00	.0000250+00	.0000303+00	.6580169-00	.6415173-00	55+00000000.	.3055701-0n
26	.0243543+00	,0000240+00	.0000556+00	.6161594-00	.6828094-00	.016/734+00	.565 598-00
27	.0243797+00	.0000224+00	.0000294+00	.9015460-00	.6828094-00	.1652315+30	.9487432-DR
200	.0244042+00	.0000212+00	.0000291+30	,7586108-00	.6828094-00	.5596626+88	.8984303-0P
01 (.0244280+00	*000020>*00	.000028/+00	.9726778=00	.6828094-00	.5e04604+35	.2322161-DF
30	.0244510+00	.0000207*00	.0000281+00	.3012900-00	.6828094-00	.76303/0+3	. 3642295-DR
4 0	00.400.44.00		00.002000.	00-7307077	00-1600609		267267
3 10	0245169400	0004774000	004202000.	1043264-00	00-1400389		1079498-00
45	.0245376+00	.0000185*00	.0000252+00	.0280621-00	6828094-00	20+0009000	.7179499-UN
35	.0245576+00	.0000184+00	.0000233+00	.9621351-UD	.6828094-00	0000000	.8848742-Dn
36	.0245770+00	.0000180+00	.0000221+00	9315224-00	.6749259-00	.0167595+00	.0733430-00
37	.0245957+00	.0000170*00	.0000219+00	.1423917-00	.6749259-00	.1650748+00	.941.446-0P
38	.0246138+00	.0000160+00	.00000217+00	,7747963-UO	.6749259-00	.5589523+00	.7461148-DP
39	.0246314+00	.0000153+00	.0000214+00	.4062409-UO	.6749259-00	.5572115+30	.0697203-00
4 0	.0246484+00	.0000153+00	.0000209+00	.5737450-00	.6749259-00	.7407450+00	.56. 831-00
4	.0246648+00	.0000154+00	.0000198+00	.2898390-00	.6749259-00	.0412052+00	.3958426-DN
42	.0246811+00	.0000129*00	.0000222+00	.6062597+00	.6749259-00	0000000000	.28/3815-00

PAGE NO.

1.00000+000 2DCANDID DIMENSION SEARCH, RZ, TWO GROUPS, 20 X 27 MESH

PROB, NO.

	4	1.0247121+00	1.000013/+00	.0000185+00	4929347-00	.6749259-00	0000000000	.8596121-0n
	45	1.0247269+00	1.0000135+00	.0000171+00	4441428-00	.6749259-00	000000000	.5958074-0n
	46	1.0247411+00	1,0000133+00	.0000163+00	4215536-00	.6731810-00	.0167564+38	.0023220-00
	47	1.0247549+00	1.0000125+00	.0000161+00	5766799+00	6731810-00	.1650401+03	.6155689-00
	8	1.0247682+00	1.0001118+00	.0000160+00	0419446-00	.6731810-00	.5587951+03	.2074244-DR
	49	1.0247812+00	1.0000113+00	.0000157+00	2419914-00	.6731810-00	.5564935+0C	.4461049-0n
	50	1.0247937+00	1.0000113+00	.0000154+00	.0406947-00	.6731810-00	.7358408+00	.0727027-0n
	51	1.0248057+00	1,0000113*00	.0000146+00	4137482-00	.6731810-00	.0340839+85	.2472017-0n
	52	1.0248178+00	1.000005+00	.0000163+00	1825563-00	.6731810-00	00+0000000.	.7757792-0n
	53	1.0248294+00	1.0000101+00	.0000149+00	.1415252-UD	.6731810-00	00000000000000	.7827489-DD
	40	1.0248406+00	1,0000101+00	.0000136+00	.0995768-00	.6731810-00	.000000000	.5596022-00
	55	1.0248514+00	1.000100+00	.0000126+00	.0636371-00	.6731810-00	000000000000000000000000000000000000000	.6400376-00
	56	1.0248619+00	1.000009*00	.0000120+00	0469958-00	6731490-00	.0167564+98	.22.1857-0n
	57	1.0248721+00	1.0000092*00	.0000119+00	1612410-00	.6731490-00	.1650395+00	.6716443-00
	50	1.0248819+00	1.0000087+00	.0000118+00	.5038974"UD	6731490-00	5587922+30	1058735-00
	29	1.0248914+00	1.0000083+00	.0000116+00	3877098-00	.6731490-00	.5564804+00	.2801472-Dn
	60	1,0249006+00	1,0000083+00	.0000113+00	1853695-00	.6731490-00	.7357508+00	.0039228-00
	61	1,0249095+00	1.00000083*00	.0000107+00	.7776632-40	.6731490-00	.0339536+3€	.3961184-DD
	62	1.0249184+00	1.0000070+00	.0000120+00	,7103664-00	.6731490-00	.00000000.	.0469826-00
	63	1,0249269+00	1,0000074*00	.0000109+00	.4074899-00	.6731490-00	000000000.	.5472040-0n
	64	1.0249352+00	1,0000074+00	.0000100+00	.0985807-00	.6731490-00	00+0000000.	.6216730-0n
	62	1.0249431+00	1,0000073+00	.0000003+00	.8342898-UD	.6731490-00	90+00000000	.9434374-DD
	99	1.0249509+00	1,0000072*00	+0000088+00	,7116541=00	.6736577-00	.0167573+06	.6343838-00
	67	1.0249583+00	1,00000068*00	.0000088+00	,5531445-UD	.6736577-00	.1650496+38	.9667496-DD
	68	1,0249656+00	1.0000064+00	.0000087+00	.1077206=00	.6736577-00	.5588381+30	.2864842-0n
	69	1.0249726+00	1,0000001+00	.00000085+00	,7587963 nu	.6736577-00	.5566897+38	.4147157-00
	70	1.0249794+00	1,0000001+00	.00000083+00	.8201890-00	.6736577-00	.7371798+30	.2113090-00
	71	1.0249859+00	1,0000001+00	.0000000+00	,3106516-00	.6736577-00	.0360245+00	.7644488-DD
	72	1.0249924+00	1,0000001+00	.0000088+00	.4143492-00	.6736577-00	00000000	.7035788-00
	73	1.0249987+00	1,00000054+00	.00000081+00	.1926933-00	,6736577-00	000000000	.6162888-00
	4	1.0250048+00	0 1.0000055+000	1,0000074+000	5,9650454-006	9.6736577-001	1.00000000+3000	1,9333092-006
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PREPARING X SECTIONS ONLY UN LUN

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4. Sample Problem 4

a. Description

(1) Problem Type

Source calculation with Chebyshev acceleration and $\mu = 1/0.9$.

(2) Configuration

(a) Geometry

rθ (quarter circle)

(b) Region Definition

The reactor consists of 18 regions as follows:

Region No. 1: Internal thermal column (water).

Region No. 2: Stainless steel shroud.

Regions No. 3-11: Fuel regions composed of enriched uranium, stainless steel, and water.

Region No. 12: Control-rod region composed of zirconium and water.

Regions No. 13-15: Reflector composed of beryllium, aluminum, and water.

Region No. 16: Outer reflector composed of aluminum and water.

Region No. 17: Iron vessel wall.

Region No. 18: Water.

(c) Mesh Definition

r direction: 33 points.

 θ direction: three points.

Buckling $(B^2) = 0.052360$.

(d) Boundary Conditions

Left: $\phi' = 0$.

Right: $\phi = 0$.

Bottom: $\phi^i = 0$.

Top: $\phi^{\dagger} = 0$.

(e) Source

Neutron source density in region 15: 1.0.

(f) Number of Energy Groups: 16.

(3) Convergence Criterion

Sum of flux difference = 10^{-5} .

b. Output Listing

PAGE NO.	
OUPS, 33 X 3 MESH	
), RT, 16 GROUPS, 3	SOMETRY
SC W/ACC),	REACTOR GEOMETRY
SOURCE CALCE WIACC	
SDCANDID	
5.00000+000	
PROB. NO.	

1	401	EOUND	1.5708+000
	BOTTOM	BOUND	0.00.0000.0
	137	BOUNE	1,9548+002
1 1	LEFT	BOUND	0000000000
	G III O III III III III III III III III	TYPE	2D - RO

REGION BOUNDARIES

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ы	6.8310+000	7,2028+000	0 0 0 0 0 0 0 0 0	1.5708+000
•	7.2028+000	7,3818+000	0.00000000	1.5708+000
10	7.3818+000	7,5578+000	0.0000+0000	1.5708+600
s	7,5578+000	7,7277+000	0.00000000	1.5708+000
~	7,7277+000	7,8951+000	0.00000000	1.5708+300
60	7,8951+000	8,0591+000	0.0000+000	1.5708+303
~	8,0591+000	2,3674+001	0.0000+000	1.5708+004
_	2.3674+001	2,3869+001	0.0000*000	1.5708+00-
	2,3869+001	2,4261+001	0.000000000	1.5708+000
C)	2,4261+001	2,5291+001	0.00000000	1.5708+00.
_	2,5291#001	3,0891+001	0.00000000	1.5708+000
4	3,0891+001	4,0291+001	0.0000+000	1.5708+000
	4,0291+001	4,0491+001	0.0000+000	1.5708+000
~	4.04914001	4,3091+001	0.00000000	1.5708+006
17	4,30914001	5,5791+001	0.0000+000	1.5708+000
90	5,5791+001	1,2309+002	0.00000000	1.5708+300
œ	1.2309+002	1,3452+002	0.00+0000	1.5708+000
0	1.3452+002	1.9548+002	0.000+000	1.5708+00

IQ.

CONSTANT INTERVAL PER REGION METHOD - X(R) DIRECTON

X(R) RIGHT	.6290+00	.8310+00	.2028+00	.3818+00	.5578+00	.7277+00	.8951 +00	.0591+00	.3674+00	.3869+00	.4261+00	.5291+00	.0891+00	0291+
NO. MESH	C)	+1	CV		wif	* 1	∓ 1	-1	100		N	e-l	+4	C)
X(R) LEFT	.0000000	,6290+00	.8310+00	.2028+00	.3818+00	.5578+00	.7277*00	.8951400	.0591+00	.3674+00	.3869+00	2,4261+001	,5291+00	.0891+00

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X(R) RIGHT	.0491+00	.3091+00	.5791+00	1,2309+002	.3452+00	.9548+00
NO. MESH	М	8	v-1	4	ed	CV.
X(R) LEFT	.0291+00	.0491+00	.3091+00	5,5791+001	,2309+00	.3452+00

HAX 8 33

X(R) PLUS DELTA X(R) FRCM LEFT TO RIGHT OF REACTOR

MESH NO. ABSCISSA	MESH NO. 1 2 3449+000 6.6290+000 6.8310+000	6,6290+000		7,2028+000	7.3818+000	7 8 9 10 10 10 10 10 10 10 10 10 10 10 10 10	7,7277+000	7.8951+000	10 8,0591+000
MESH NO. ABSCISSA	MESH NO. 11 1269+001 2.3674+001	12 1.8469+001	13	15 2,4065+001	162.4261+001	14 15 19 20 2,3869+n01 2,4065+001 2,4261+001 2,5291+001 3,0891+001 3,5891+001 4,0291+001	18 3,0891+001	3.5591+001	20 4,0291+001
MESH NO. ABSCISSA	21 22 23 4,0358+001 4,0424+001 4,0491+001	22	23	25	26 5,5791+001	24 27 28 29 30 4.1791+n01 4.3091+001 5.5791+001 7.2616+001 8.5441+001 1.0627+002 1.2309+002	28	29	30 31 32 31 31 31 31 31 31 31 31 31 31 31 31 31
MESH NO.	MESH NO. 31 32 33 ABSCISSA 1,3452+002 1,6500+002 1,9548+002	1.6500+002	33						

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MESH DEFINITION

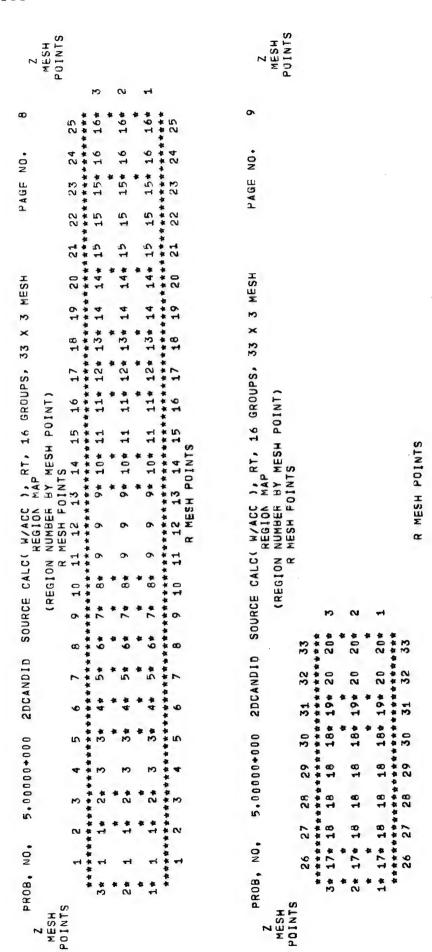
CONSTANT INTERVAL PER REGION METHOD . Y(Z) DIRECTION

Y(Z) TOP	.5708+00	.5708+00	.5708+00	.5708+00	.5708+00	.5708+00	.5708+00	.5708+00	.5708+00	.5708+00	.5708+00	.5708+00	.5708+00	.5708+00	.5708+00	.5708+00	.5708+00	.5708+00	1.5708+000	.5708+00
MESH	ю	2	М	ĸ	М	m	m	ю	Ю	М	M	M	М	M	m	M	М	M	M	ю
0N																				
Y(Z) BOTTOM	.0000+000	.0000+00	,0000+000	.00000+000	.0000000	.00000+000	.00000+00	.00000+000	.00000+00	.000000	.00000+000	.00000+000	.00000+00	.0000+000	.00000+000	.0000000	.00000+000	.00000+00	0 0 0 0 0 0 0 0 0 0 0	.00000+000

Y(Z) PLUS DELTA Y(Z) FRCM BOTTOM TO TOP OF REACTOR

Y(Z) MESH NO. 1 2 2 ORDINATE 5,2360-001 1,0472+000 1,5708+000

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4	.1111	11+00	, 14426	1-00	4046963+00	\$5997858±00	.00000000.	.00000000.	.3272536+0
IU	.1111	11+00	87179	6-00	.2997538+00	.4676681=00	0000000000	00+0000000*	1257400-00
9	1111	11+00	29644	8-00	,2296130+00	4355254+00	0000000000	000000000	0-9589666
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0 0	.1111	11+00	110730	7-00	1505645+00	,4249271=00	.8718817#00	.0171065+00	3983430-0
0.	1111	11+00	,12264	5-00	1276390+00	\$6197320=00	8718817=00	.169n032+00	.1537491-0
10	.111	11+00	14290	6-00	.1083307+00	,1583471 = 00	.8718817w00	.576A954+00	.9404048-0
	1111	11+00	17284	00-0	00+62236400	5566303#00	871881	.64093n9+00	1,7224950-001
	1111	11+00	,23084	2-00	,0821272+00	*3332899#00	8718817=00	,3530592+00	.5904303-0
13	1111	11+00	45457	00-6	.0466844+00	.1643848+00	8718817*00	.374883+00	.0122657-0
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15	.111	11+00	47132	8-00	2991359+00	4047574=00	.8718817*00	.00000000.	.2896646+0
1.6	. 1111	11+00	,82174	00-0	,8665534+00	8281673-00	8718817#00	.0000000000	.7683359+0
17	.111	11+00	, 82862	00-6	4787818+00	.6627896=00	.8718817#00	00+0000000+	.3804955+0
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27	.111	11+00	85854	6-00	,0654040+00	,0029706+00	.9209798*00	,5277813+00	.954984T+0
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60	.1111	11+00	,72107	4-00	.0091386+00	.5960884=00	.9209798e00	00.40000000	.7031253+0
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32	1111	111+000	740	5	081076+	355247	209798.0	1.0171931+000	3,4032156-00
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                   6=1
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  MAX IN=
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CANDIDED HAS COPIED FLUX FROM LUN 49 ONTO LUN 48
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HA ERRLAM	08 2,9898517-00	2.9111673-002	2,8346915-00	2. /6:4046+0n	2.6882833-00	0. 2.617.843-00	0. 2.5375632-00	0. 2.4334572-0n	0: 2.2648892-00	01 2.0095105-0n	0. 2.6142706-00	0. 2.4514059-00	0 2.2999124-0n	0: 2.1586281-00	0- 2,0243232-0n	0- 1.8799213-00	0- 1.699:603-00	0. 1.4238068-00	01 1.0866081-0n	0 1.7059804-0n	00 1.6183320-0n	0. 1.5387583-DR	0 1.4622140'-On	0. 1.3873835-0n	60 1.3047739-00	0. 1.1983428-DR	1.0365961-00	01 7.6325987-00	0. 1.1238348+0n	1.06/1817-00	00 1.0131219+0n	0. 9.6154342-DR	0. 9.1146978-DR	B. 8.5654491+00	00 7.8626821-00	10-1604004 10	51 4.97.4/U05-U1	10 20 2T 40 4 C 10
ALP 1.00000000+5	+00000000	1.0000000+0	+0000000	+00000000	.0171541+	1695438+	.5793872+	.6528357+	.4483759+	.7668773+	+00000000	+00000000.	+00000000	.0171541+	.1695438+	.5793872+	. 6528357+	.4483759+	.7668773+	+00000000.	+00000000	+00000000.	.0171541+	.1695438+	.5793872+	.6528357+	.4483759+	,7668773+	+00000000	+0000000.	+0000000.	.0171541+	.1692438+	.5793872+	.6528357+	* 4400/00/4 400/00/01	+8//00//	
SIGMA 1.0000000+000	.0000000+00	1.0000000+000	0000000+00	0000000+00	8988792-00	8988792-00	.8988792-00	.8988792-00	.8988792-00	.8988792-00	.8988792-00	.8988792-00	.8988792-00	.8988792-00	.8988792-00	.8988792-00	.8988792-00	.8988792-00	.8988792-00	.8988792-00	.8988792-00	.8988792-00	.8988792-00	.8988792-00	.8988792-00	.8988792-00	.8988792-00	.8988792-00	.8988792-00	8988792-00	.8988792-00	8988792-00	.8988792-00	.8988792-00	8988792-00	00-26/8868.	00-20/0000	00.26/0060
CHANGE IN PHI 4.2917830-002	,2484860-00	4,2055924#002	165098/*00	.1210011-00	1492728 = 00	,7217659-00	,3008262-00	.0413691-00	,4631174-00	.3231098*00	,2323329 * 00	,2082398-00	.1844032+00	.1978823-00	4994008-00	,3326081-00	.5020761-00	,2990396 ** 00	.9463547-00	.1237388-00	.1110902*00	.0985801-00	1048384-00	.2558034-00	.6735137-00	.7608594-00	.5100099-00	,4696575-00	4378816-00	3743548-00	,3115636-00	,3395455-00	,0665236-00	.0804036-00	.3321676-00	10/0//0=00	5458070	0040/60646.
K*UPPER 1,0070435+000	,0067996+00	1,0065689+000	.0063506+00	,0061438+00	.0059476+00	,0057583+00	.0055519+00	.0052899+00	.0048856+00	.0043076+00	.0119824+00	.0107151+00	,0095586+00	.0085009+00	,0075150+00	.0064769+00	.0052094+00	,0033516+00	,0020708+00	.0051777+00	.0047478+00	.0043358+00	.0039415+00	.0035579+00	.0031371+00	,0025997+00	.0017649+00	.0009942+00	.0034365+00	0031416+00	,0028669+00	.0026105+00	.0023669+00	.0021057+00	.0017803+00	.0012907+00	0004/02+00	105000500
K-LOWER 9,7633586-001	.7690108-00	9,7745727-001	7800369-00	.7853971-00	.7906479=00	.7958747-00	8017631-00	.8095534=00	.8223668-00	.8421246-00	,8583967-00	,8620101-00	.8655950-00	.8691462-00	,8727179-00	.8767770-00	,8821878-00	.8911357-00	,9120467-00	.8816790-00	.8856446-00	.8894823-00	,8931933-00	.8968409-00	.9008941-00	.9061629-00	.9145889-00	,9336160-00	,9219812-00	.9246977-00	9273568-00	,9299505-00	.9325223-0	.9354027-00	.9391757-00	**************************************	04040040	00-00//404
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CANDIDED HAS COPIED FLUX FROM LUN 42 ONTO LUN 48

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33 JMAXOUT=
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IMAXINE

1,7881408-003 1.6175935-003 1,5019068-003 1,0375962-003 7,2693973-004 3,5052630-003 2,0523857-003 2.0449494-003 1.8706935-003 8,2018253-004 6,8418423-004 ERRLAM 5.5944869-003 5,3128525-003 5.0431601-003 4.5391230+003 4,3038653-003 4.0753992-003 3,8249751-003 3.0086683-003 1,9562945-003 1.7071930-003 .3263957-003 7.7225119=004 4.7854159-003 1.5776599+000 1.0000000+000 1.0171211+008 2.6445757+000 1.000000000.1 1.00000000+000 1.0171211+000 1.5776599+000 1.000*0000001 1.00000000+000 1.00000000.1 1.00000000+000 1.00000000+000 6.3826612+000 2.6445757+000 5.3826612+003 3.4703412+601 000+0000000 1.0171211+006 ALPHA 1.1691692*000 1.1691692+000 3.4703412+001 1.000000+060 1,00000000+000 1,00000000+0000 1,00000000+000 9,8801733-001 9.8801733-001 9,8801733-001 1.0000000+000 9.8801733-001 9.8801733-001 9.8801733-001 9.8801733-001 9,8801733-001 9.8801733-001 9,8801733-001 9.8801733-001 9,8801733-001 9.8801733-001 9.8801733-001 SIGMA 1.0000000+000 9,8801733-001 9,8801733-001 9,8801733-004 9.8801733-001 9,8801733-001 CHANGE IN PHI 2,5154339-003 5,4522171=004 2,4853205-003 2,3970677-003 2,4088864-003 3,6391227-003 7.0213880-002 1,1791858-003 1,1647682=003 1.1505259-003 1,1559151-003 1,7453262-003 2,8691731-003 3,3586550 = 002 5,4277242-004 2,4555556-003 2,7352229*003 5,9846926-003 1,3985631-002 6,7006828=003 5.5650581-004 2,4261393-003 1,3121851=003 5,4958999-004 1,0002565+000 1,0001979+000 1.0001008+000 1.0003601+000 1,0003308+000 K-UPPER 1.0026975+000 1.0014716+000 1,0006622+000 1,0003507+000 1,0002859+000 1.0001599+000 1,0025326+000 1,0023747+000 1,0020791+000 1,0019409+000 1.0016594+000 .,000±093*000 1.0004242+000 1,0022236+000 .,0003173+000 .0002283+000 ..0003912+000 000+9908100 000450875000 大きてつぎ原氏 9,9796637-001 9,9830579-001 9,9841525-001 9,9846832-001 9,9865798-001 9,9906317-001 9.9964661=001 9.9710296-001 9,9733158-001 9.9753996-001 9,9763702-001 9.9773121-001 9,9783445=001 9.9860986-001 9,9836104=001 9.9858033-001 9,9878290-001 9.9960404-001 9.9963313-001 9,9721976-001 9.9743823=001 9,9817184-001 9,9852110=001 9.9961894-00 1.1111111+000 1.111111+000 1.111111+000 1.111111+000 1.111111+000 1.1111111+000 1.111111+000 1.1111111.000 1.111111+000 1.111111+000 1.111111+000 1.1111111000 1.111111+000 1.111111+000 1.111111+000 1.111111+000 1+111111+000 1.111111+000 1.11111111-000 1.111111-000 1.11111100 1.11111114000 1-41111-00 1-111111-000 ITERATION HISTORY ITERATION 9,981

PAGE NO.

SOURCE CALC (M/ACC), RT, 16 GROUPS, 33 X 3 MESH

2DCANDID

5.000000+000

PROB. NO.

Z	ERROR	MAS	AN ERROR WAS ENCOUNTERED IN SUBROUTINE	Z	SUBROUTINE	KEFFCALC	A	AT STATEMENT	NUMBER
-	IME EXCEEDED	EDE	2						

	.0487365-00	.3824139+00	.3590147-00	.7402481-00	.0799144-00	.4131780=00	.7359661-00	.0496992-00	.3540131-00	.6502767-00	.2847737-00	.0405397=00	.4902914-00	.1498606-00	1530547-00	.1334790+00	.0055524+00	.5979689-00	.3509956+00	.4527890-00	.4359596 m 00	.4428157-00	.4316473-00	1326595 00	.1315994-00	.8226343-00	. 8687442=00	.2134490-00	.8399465 00	.1481203-00	9.7529784-010	.2383444 * 01	.4704153-01
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CANDIDED HAS COPIED FLUX FROM LUN

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MAXINE
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33 JMAXOUT=

3. IMAXOUT

33 JMAX1N=

IMAXINE

5. Sample Problem 5

a. Description

(1) Problem Type

Source calculation without Chebyshev acceleration and $\mu = 1/0.9$.

(2) Configuration

Same as Sample Problem 4.

(3) Convergence Criteria

Same as Sample Problem 4.

b. Output Listing

NO. 143	ERRLAM 3.1753051+001	2.3546576 000 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0
SH PAGE N	ALPHA 1.0000000+35	
GROUPS, 33 X 3 ME	SIGMA 1.00000000+000	10000000000000000000000000000000000
ACC 1, RT, 16 G	CHANGE IN PHI 1,0180383+000	4.6202040 1.59978581 1.46766811 1.46766811 1.4713371 1.3858227 1.387328061 1.387328001 1.387328001 1.387328001 1.387328001 1.387328001 1.387328001 1.3873281 1.387281 1.387281 1.387281 1.398661 1.2998661 1.201998661 1.201998661 1.201901
SOURCE CALCINO	K-UPPER 3,1753250+001	2.355273+000 1.2097538+000 1.2296130+000 1.1823985+000 1.1823985+000 1.1823985+000 1.187974+000 1.0079018+000 1.073645+000 1.073645+000 1.073645+000 1.073645+000 1.07581600 1.0652641+000 1.0652641+000 1.0558805+000 1.0558805+000 1.0579953+000 1.0579953+000 1.0579953+000 1.0579953+000 1.0579953+000
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8	00	G#1	47		n	5	78-00				
25	0	22	7#7	L UX	5	5	00-40				
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*	23	18	つまけ	FLUX	10	46	97 ± 00				
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	27	Gel	J=1	FLUX=	9	0 6	14+00	4			
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33 JMAXOUTE

HAXOUT

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UNAXINE

33

I MAXINE

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JMAXOUT
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                                        COPIED FLUX FROM LUN
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PAGE NO.

SOURCE CALCINO ACC), RT, 16 GROUPS, 33 X 3 MESH

2DCAND I D

6,00000+000

PROB. NO.

CANDIDED HAS COPIED FLUX FROM LUN 3 ONTO LUN 48

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33 JMAXDUT#
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2.7760176-002 4.0081391+002 3.9657094-002 2.7181816-002 2,6336618-002 2,6336618-002 2,6897117-002 4.0511626+002 PAGE NO. 1.00000000+000 1.00000000+000 1.0000000+000 .000+0000000. 1.00000000+0000 1.0000000+000 .000+00000000. .00000000. .0000000000. .000+0000000. .000000000. .00000000. .00000000* 1.0000000+000 .0000000+0000 .000000000. .0000000000. .00000000+000 .00000000+0000 .000000000. .0000000000. .00000000+000 .0000000000 .00000000+0000 .0000000000. .00000000000 .000+0000000. 1.0000000+000 .00000000+000 .000+0000000. 1.0000000+000 1.0000000+000 1.00000000+0000 1.0000000+000 000+00000000. 1.0000000001 1,0000000+000 1.000000000.t 1.0000000+000 1.00000000+000 1.0000000+000 1.0000000+000 1.00000000.1 1.00000000+000 1.0000000+000 1.0000000+000 1.0000000+000 1,00000000001 1.00000000001 1.00000000000 1.0000000+000 1.00000000+1 1.0000000+000 1.00000000+000 1.0000000+000 1.0000000+000 1.0000000+000 1.0000000+000 1.00000000001 1.0000000+000 1.0000000+000 1.000.0000.000 1.0000000+000 1.0000000+000 ..000000000. .000+0000000 1.00000000+1 1.000000000.1 1.0000000+000 1.0000000+000 1.0000000+000 1.00000000000 1.00000000001 m SOURCE CALCING ACC 1, RT, 16 GROUPS, 33 X 6.4311637#002 6.3722107#002 6.3137506#002 6.2557792#002 CHANGE IN PHI 7.9343005=002 .1134812-002 6.9846655 002 6.9210708-002 6.7954815*002 6.7334770=002 6.6719927#002 6.6110239-002 6.5505658+002 6.4906139=002 6.1982921-002 6.0847548m002 6.0286965 002 5.9179809-002 5,8633159-002 5,7553526-002 5.4932474#002 5.4421358#002 7.8626135m002 7.7915460m002 .7210910=002 .6512416#002 .5819912=002 .5133335002 .4452614=002 .3777694 m002 .3108512=002 .2445010=002 .1787128-002 .0488006=002 6.8580111-002 6.1412854m002 5.9731065#002 5.8091077-002 5,7020469#002 5.5447908=002 5.6491871-002 5.5967696#002 1,0091332+000 K-UPPER 1.0127452+000 1.0125276±000 .0113346±000 1,0108017+000 1,0106323+000 1.0104669+000 .0101471+000 1,0098415+000 1.0096937+000 .0094075+000 1.0092689+000 .0088702+000 .0087426+000 .0084952+000 .0083751+000 ,0082573+000 .0081418+000 .0080285+000 .0079174+000 .0078083+000 .0077013+000 .0075963+000 .0074931+000 .0073919+000 .0072924*000 .0070988+000 .0070045+000 000+6176900 .0121094+000 1.0119083+000 1,0117122+000 .0115211+000 1,0111527+000 .0109751+000 .0099926+000 .0095491+000 .0071947+000 KaLOWER 9.7223353=001 9,7244620=001 9,7265864-001 9,7287086=001 9,7308286-001 9,7329464=001 9,7350619-001 9,7371751.001 9,7392858-001 9,7413940=001 9,7434994#001 9.7456019 001 9,7477013#001 9,7518896-001 9,7539780=001 9,7560622 001 9,7581418-001 9,7602167-001 9,7622863#001 9,7643504-001 9,7664087-001 9,7684606=001 9,7705060*001 9,7725443-001 9,7745752=001 9,7765984#001 9,7786133-001 9,7806198 001 9,7826172-001 9,7846053-001 9,7865837,001 9,7885520-001 9,7905099-001 9.7924569+001 9,7943926-001 9,7963169=001 9,7982292=001 9,8001293-001 9,8020169 001 9,8038915#001 8057530=001 **2DCANDID** 6.00000+000 K.EFFECTIVE 1.1111111+000 .1111111000 .111111+000 .1111111+000 .1111111+000 .111111+000 .1111111+000 1,1111111+000 .11111114000 ,1111111+000 .1111111+000 .1111111+000 .1111114000 .11111111+000 1111111+000 .111111+000 .11111114000 .1111111+000 1,111111+000 .1111111+000 .1111111+000 .1111111+000 1.1111111+000 .111111+000 11111114000 1111111+000 1.111111+000 1,1111111+000 1111114000 .1111111+000 11111114000 .1111111+000 1,111111+000 1,111111+000 .1111111+000 1,111111+000 1,111111+000 .1111111+000 1.111111+000 1,1111111+000 .1111111+000 ITERATION HISTORY PROB, NO. 10,006 ITERATION

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	18	算	-	FLUX#	17	41575	00				
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		U # 1	FLUX	.9984634=00	
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	88.		FLUX	.5984634-00	

6. Sample Problem 6

a. Description

(1) Problem Type

Real k_{eff} calculation

(2) Configuration

(a) Geometry

 $r\theta$ (full periodic)

(b) Region Definition

The reactor consists of a core region composed of uranium (U^{235} and U^{238}), iron, and aluminum.

(c) Mesh Definition

r direction: 20 points

 θ direction: 20 points

(d) Boundary Conditions

Left: $\phi^{\dagger} = 0$.

Right: $\phi = 0$.

Bottom: periodic.

Top: periodic.

(e) Number of Energy Groups: 2.

(3) Convergence Criteria

 k_{eff} difference = 10^{-6} .

 k_{eff} bounds = 10^{-3} .

Sum of flux difference = 10^{-3} .

Periodic = 10^{-7} .

Maximum number of periodic iterations = 25.

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N

PAGE NO.

TOP BOUND	3.1416+300
BOTTOM	000+0000+0
RIGHT	5,0000+001
LEFT BOUND	000+0000*0
GEOMETRY TYPE	2D * R0

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TOP BDRY.	3,1416+383
BOTTOM BDRY.	0 0 0 0 0 0 0 0 0 0
RIGHT BORY,	5,0000+001
LEFT BDRY,	0 0 0 0 0 0 0 0 0 0
REGION NO.	Ħ

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2 GROUPS,	DEFINITION
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1.01000+001	
NO.	
PROB.	

PAGE NO.

M

CONSTANT INTERVAL PER REGION METHOD - X(R) DIRECTON

X(R) RIGHT	5.0000+001
NO. MESH	20
X(R) LEFT	0.0000000000000000000000000000000000000

IMAX # 20

XKRY PLUS DELTA XKRY FROM LEFT TO RIGHT OF REACTOR

2.5000+000 5.0000+000 7.5000+000 1.0000+001 1.2500+001 1.5000+001 1.7500+001 2.0000+001 2.2500+001 2.5000+001 11 17 18 19 20 2.7500+001 3.2500+001 3.8000+001 3.7500+001 4.0000+001 4.2500+001 4.5003+001 4.500+001 5.0000+001 MESH NO. ABSCISSA MESH NO.

MESH CEFINITION

PAGE NO.

CONSTANT INTERVAL PER REGION METHOD - Y(Z) DIRECTION

7(Z) BOTTOM NO. MESH Y(Z) TOP 0.0000+000 Z0 3.1416+000

JHAX B 20

Y(Z) PLUS DELTA Y(Z) FRCM BOTTOM TO TOP OF REACTOR

	1	□ W # □ C # □ C # □ O # □ D □ D □ D □ D □ D □ D □ D □ D □ D □
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PAGE NO.

PROB, NO. 1.01000+001 2DCANDID RT FULL CIRCLE 2, 2 GROUPS, 20x20 MESH

ERATI	Z	EFFECT IV	5	K UPPER	CHANGE IN PHI	SIGMA	ALPHA	ERRLAM
4 4	***	# 1 TEST	.8949474-00	****				
- 444	el	# 2 TES	9,8656537-0	***				
	1.0	77970+000	1892435-001	1759743+000	1,8323686-001	1,00000000+000	1:00000000+000	8.5704997-001
	0.1	1	11 44					
***	н	H LING	4.9665858-00	***				
		* 2 TEST	-1.2262625-00	****				
	2 1,11	45811+000 6	1438095*001	1362174+000	1.2648776-001	1,00000000+000	1:0000000+0C	5.2183643-001
*	grif.	# 1 TEST	5,0932263-00	****				
***	**	F. 2 TEST	-1.2513913-00	***				
	3 1.17	01205+000 7	4071431-001	1090072+000	1.0301649-001	1.00000000+000	1 * 0 0 0 0 0 0 0 + 0 0 0 C	3.6829291-0n1
1 444	1	H TEST	4.9614216-00					
存在	# #	2 TEST	-1.2162839-00	****				
	CV	8 998+000 8	0331654-001	0892549+000	8.9050481*002	1.00000000+000	1.000000000	2.8593831-001
***	**	1 TESPE	4,7283319-0	1				
***	#	2 TEST	-1,1576486-00	****				
	CV	9657+000 8	4026287-001	0764826+000	7,8539607+002	1,00000000+000	000+00000000t	2,3621975-001
*	T1	1 TEST	4.4505998-00					
-	eri H	2 TEST	-1.0895356-00	****				
	6 1.3	59289+000 8	6473834-001	0668683+000	6.9835569-002	1.0000000000000	1,00060000+600	2.0213001-001
. ****	THE RE	TESE	4.1542189-00					
***		TEST O	** n185465*nn	****		A MANAGEMENT OF A 1-1-1-1-1-1-1-1-1-1-1-1-1-1-1-1-1-1-1		
	7 1.3	00	8231791-001	.0585935+000	6,2358116*002	1,00000000+000	1.000000000	1,7627562-001
***	**	# 1 TESP#	11118529740-00					
		2 TESTS	-9,4811885-00	****			THE TAXABLE WITH THE PROPERTY OF THE PROPERTY	
	"	42831+000 8	.9580962-001	.0515231+000	5.5835605-002	1.00000000+0000	1,00060000+1	1.5571352-001
4	**	TES!	3,5548205-00	****				
***	1 (2 TEST	-8,8012061-00	****				
	9 1.42	59406+000 9	.0660857-001 1	.0454845+000	5,0107598-002	1,000000000000	7.00000000+000	1,3887591-001
-	~	a 1 TEST	3.2646590-0	***				
***	**	s 2 TEST	-8,1557400-00	本本學者本				
	10 1.4	58893+000	,1553015-001	.0403151+000	4.5062839-002	1.00000000+000	1.0008300+800	1,2478498-001
444	-+	1 TEST	2,985482	****				
		2 TEST	•7.5501389-00	****				
	11 1,6	13979+000 9	.2308049-001	.0358750+000	4.0615840-002	1,0000000+000	1.00000000+000	1.1279453-001
444	-1	1 TEST	2,7191209-00					
1 444	7 71	N	-6,9865373	****			A THE OWNER, TWO REPORTS THE REAL PROPERTY AND ADDRESS OF THE PROPERTY PARTY AND ADDRESS.	
	12 1,5	37426+000 9	.2958969-001	.0320475+000	3,6696665=002	1,00000000000	1.00000000+696	1.0245777-001
***	+	B 1 TEST	2,4665394					
444	***	B 2 TEST	-6,4655359-00	****				
	13 1,9	6 000	.3528228-001	.0287368+000	3,3245709-002	1.00000000+000	1.0000000+000	9,3454519-002
**	1 2 2	# 1 TEST	2,2282129-0	****				
404		2 TEST	-5,9864660					
	14 1,5	0020200	.4031661-001	.0258642+000	3.0210848*002	1.000000000000	1.0000000000000000000000000000000000000	8.5547547.0n2
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***	# ti	2 TEST= -5,5478667=00	0				
		G = 1 TEST 1,7942671-006		200 00000000000000000000000000000000000	nonenanne.	~9>±00000.T	7:10=T4/9660*/
* *	4 4 4	1. Ge 2 TEST# 15.1479210*006 44444 1.96675314010 9.4894530*0011 1.0211855 4. Ge 7EST# 1807084**********************************	25+000	2.5606177-002	9.1178142-001	1.0157792+65	7.2342240-002
*	1 69:	1 6 1 2 TEST# 14,7794470100	,	1 4 4 4 4 4 4 4 4 4 4 4 4 4 4 4 4 4 4 4			
なな		//0801#000	000+/*	2.0075502-002	9.11/8142-001	1.1541044400	6,6684649- ₀ n2
***	100 m	1 G# 2 TEST# -4, 4093236=006	6				
* * *		624442+000 7.2627298+0 64 1 TEST# 1.107846	18+000	3,1846270=002	9,11/8142-001	1;5103402+00.	6.0785866-002
* * * *		1 GF 2 TEST - 3,9807501-006					
1		916165+000 9.6148072*001 1.	0152876+000	4,4186528*002	9.1178142-001	2:3467783+60	5.3806868-002
	* **	FORONOS / HINDE I BE CONTROL OF BE CONTROL O					
		5950008#000 9.6844973 001 1.	0127300+000	7.2437320-002	9.1178142*001	4.5096667+0L	4.428 292-002
*	11	G# 1 TEST# 1.2915234*007					
		1 4# 2 EST# #2.00100/3 .5942932#000 9.7916873=00	17+000	1.2125279"001	9.1178142-001	9,6383054+000	2,9952941-002
***	H:	1 G# 1 TEST# -3,6624147-00		,			
4 4	H 6	1 G# 2 TEST# +9.435038	000137	7084747400	447840-04	C	4 4 4 4 4 4 4 4 4 4 4 4 4 4 4 4 4 4 4
*	, m	GR A THENTH IN CONTRACTOR		000-10/1/2/14	T90.34T0.TT.		Zan Lattory . T
***		1 G# 2 TEST# =6.3037942=007					
		5810375+000 9,8899382*001 1.	000+20	9,2307707-003	9,1178142-001	1,0000000000000000	1.4256923-002
***	a (G# 1 TEST# -2,4607971=007					
t t	E CV	97429724000 9.8943986#0	59+000	8,8030827-003	9.1178142-001	1.0000000000	1.3246089-002
***		1 Ge 1 TEST= -2,0709990=00			1	2 2 3	
4	H: C	1 Ga 2 TEST# #1,4476336-007	000.07	2 2 2 2 2 2 2 2 2 2 2 2 2 2 2 2 2 2 2 2			
***	A H	**************************************	0022240+000	0,5502204*005	Y.11/0142*UU1	I+015//82+00	1.23/5035-012
	8	5608647+000 9,9031894=001 1.	17+000	9,3188972-003	9,1178142-001	1,1541;44+00	1.1592723-002
* * * *	n	GB 1 TEST# +1.442970					
		55423344000 9.9078427*001 1.	37+000	1,1680312+002	9.1178142-061	1.5163402+36	1.0779456-002
****	850 f	1 Gm 1 TEST= -1,0177041-00					
de de		1 G# 2 TEST# 3,8679127*007	0	00-0002026	C		7847884-004
***	 	**************************************	000+20	1./363220#006	7.11/0142*UU1	Z + 240 / 102 + 20	9./03/03T*UH3
	50	5411102+000 9.9199434=001 1.	82+000	3,1578584-002	9,1178142-001	4.5096667+06	8,3238335-003
****		1 Gs 1 TEST# 1.148025					
	80	345186+000 9,9299205+001 9.	11-001	6,3574260-002	9.1178142-001	9:6383054+90	5.7720605-003
* * * * *	8 : 8	G# 1 TEST# 3,5097719*007					
	11 12 12 13 14 14 14 14 14 14 14 14 14 14 14 14 14	77979+000 9.9436235*001 9.	9618255-001	6.2475266-003	9.1178142-001	1.000000000	1.8201916+003

1.7148403-003	1,7965441-003	1,8707129,003	1.9289601-003	1.9649565-003	1,9582275-003	1,8561115-003	1.5948026-003	8,9273452-004	9.0546141-004	8,933,610-004	9.8505919-004	1.0094074-003	1,0295260-003	1,1388566-003	1,1297189-003
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9.1178142-001	9,1178142-001	9,1178142-001	9,6005474*001	9.6005474*001	9,6005474-001	9,6005474-001	9.6005474+001	9.6005474-001	9.6005474-001	9,6005474-001	9.6005474-001	9,6005474-001	9.6005474-001	9,6005474-001	9.6005474-001
6,0309751-003	5,8041562-003	5,5723076-003	5,4264489=003	5,9359571*003	7,5392143=003	1,1597475-002	2,3628934*002	6,7512637-002	3,1164338-003	3,0186846-003	2,9125716-003	2,8484522-003	3,1275136-003	3,9870434*003	6,1595312=003
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FLUX NORMALIZATION FACTOR IS 5.84356-001

7. Sample Problem 7

a. Description

(1) Problem Type

Real $k_{\mbox{eff}}$ calculation

(2) Configuration

Same as Sample Problem 6

(3) Convergence Criteria

 k_{eff} difference = 10^{-6} .

 k_{eff} bounds = 10^{-3} .

Sum of flux difference = 10^{-3} .

Periodic = 10^{-4} .

Maximum number of periodic iterations = 5.

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20 1,5925080000 9.6862220011 1.033414000 7.33721372012 9,1264730001 4,5247475+000 8,49992 1.9925080000 9.6862220-001 1.033414000 7.33721372002 9,1264730001 4,5247475+000 8,49992 1.8 1.9 1.9 1.9 1.9 1.9 1.9 1.9 1.9 1.9 1.9	※ななな.		Ga 2 TES	T# #8.9070944	4				
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22 1,5663823000 6,52766274000 40.0000000000000000000000000000000	***	#	G# 2 TE	T# #2.0985024#A	4				
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23 1,5785646000 8,8282866001 1,10022882*000 9,1264730*001 1,0000000*000 2,29788 2 1568750*001 1,0000000*000 1,29788 2 1568750*001 1,0000000*000 1,29788 2 1568750*001 1,0000000*000 1,29788 2 1,264730*001 1,0000000*000 1,29788 2 1,264730*001 1,0000000*000 1,29788 2 1,264730*001 1,10157344*001 1,10157344*001 1,10157344*001 1,10157344*001 1,10157344*001 1,1015734*001 1,1015734*001 1,1015734*001 1,101573*001 1,101573*001 1,101573*001 1,101573*001 1,101573*001 1,101573*001 1,101573*001 1,101573*001 1,101573*001 1,101573*001 1,101573*001 1,101573*001 1,101573*001 1,101573*001 1,101573*001 1,101573*001 1,101573*001 1,101573*001 1,101573*001 1,101573*001 1,101573*001 1,101573*001 1,101573*001 1,101573*001 1,101573*001 1,101573*001 1,101573*001 1,101573*001 1,101573*001 1,101573*001 1,101573*001 1,101573*001 1,101573*001 1,101573*001 1,101573*001 1,101573*001 1,101573*001 1,101573*001 1,101573*001 1,101573*001 1,101573*001 1,101573*001 1,101573*001 1,101573*001 1,101573*001 1,101573*001 1,101573*001 1,101573*001 1,101573*001 1,101573*001 1,101573*001 1,101573*001 1,101573*001 1,101573*001 1,101573*001 1,101573*001 1,101573*001 1,101573*001 1,101573*001 1,101573*001 1,101573*001 1,101573*001 1,101573*001 1,101573*001 1,101573*001 1,101573*001 1,101573*001 1,101573*001 1,101573*001 1,101573*001 1,101573*001 1,101573*001 1,101573*001 1,101573*001 1,101573*001 1,101573*001 1,101573*001 1,101573*001 1,101573*001 1,101573*001 1,101573*001 1,101573*001 1,101573*001 1,101573*001 1,101573*001 1,101573*001 1,101573*001 1,101573*001 1,101573*001 1,101573*001 1,101573*001 1,101573*001 1,101573*001 1,101573*001 1,101573*001 1,101573*001 1,101573*001 1,101573*001 1,101573*001 1,101573*001 1,101573*001 1,101573*001 1,101573*001 1,101573*001 1,101573*001 1,101573*001 1,101573*001 1,101573*001 1,101573*001 1,101573*001 1,101573*001 1,101573*001 1,101573*001 1,101573*001 1,101573*001 1,101573*001 1,101573*001 1,101573*001 1,101573*001 1,101573*001 1,101573*001 1,101573*001 1,101573*001 1,101573*001 1,101573*001 1,101573*001 1,101573*001 1	你在在女		2 TE	T# #2.4645124#0	****				
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27 1,535.665-000 8 908389-001 1,0015543-000 1,1569.69002 9,1264730-001 1,5110725-000 1,007289 28 1,947205-000 9,943553-001 1,001058-000 1,717642-002 9,1264730-001 1,5110725-000 1,0264730-001 29 1,948815-001 1,001054-001 4.0016403-000 1,717642-002 9,1264730-001 4,5247475-000 8,61568 2 151 1,52476-001 8 9202445-001 1,001054-000 5,3284102-002 9,1264730-001 4,5247475-000 8,61568 3 1,52476-59-000 8 9202445-001 1,001054-000 6,3284102-002 9,1264730-001 1,001000000 8,54999 1		56	. 5601267+00	9037519-00	00.	9,2536661,00	.1264730000	.1542734±00	323372-00
28 1,5472505+000 9,9135633*001 1,0010588+000 1,7176642=002 9,1264730=001 2,3497836+000 8,70246 18 1 0* 2 FEST* 1,556964=001 1,0005673*000 3,1370929*002 9,1264730=001 4,5247475+000 8,61586 20 1,5488156+000 9,920445*001 1,000535+000 6,3286162=002 9,1264730=001 4,5247475+000 8,61586 30 1,5483526+000 9,920445*001 1,000335+000 6,3286162=002 9,1264730=001 4,5247475+000 8,61586 31 1,57659+000 8,990544004 ***** 31 1,57659+000 8,5905386*001 1,000335*000 6,3286162=003 9,1264730=001 1,0000000000 8,54999 31 1,57659+000 8,5905386*001 1,0002321+000 5,9178824=003 9,1264730=001 1,0000000000 8,54999 32 1,5182830+000 9,9286266=001 1,0072221+000 5,9178824=003 9,1264730=001 1,00000000000 8,77499 32 1,5182830+000 9,9286266=001 9,9864130=001 5,7738714=003 9,1264730=001 1,0000000000 8,77499 33 1,5182830+000 9,9286266=001 9,9864730=001 1,00000000000 2,7776 34 1,5095312=000 9,9286266=001 9,98644204=001 5,778854=003 9,1264730=001 1,00000000000 2,7776 35 1,51828202000 9,9286266=001 9,98644204=001 5,8758096=003 9,5900781=001 1,1633906+000 2,7776 35 1,51828202000 9,990000 9,9900000 9,9864204=000 1,4484204 35 1,4882202000 9,990000 9,9900000 9,980000000000000		27	\$536685*00	9083309=00	00	1,1589469000	.1264730#00	.5410725+00	7281
18		28	5472505+00	9,9135633"00	00.	1,7176642000	.1264730 00	.3497836+00	.7024855=00
29 1,5468156+000 9,920245=001 1,0006403+000 3,1370929#002 9,1264730#001 4,5247475+000 8, 2 TEST# 2.0108044=0014 4**********************************	4	*	G# 2 TE	T# 1.3569644m	04 44				
State 1 Ge		N	,5408156+00	9202445m00	,0006403+00	11370929#00	.1264730 m00	.5247475+00	61588
30 1,5343526+000 9,9074816=001 1,0003635+000 6,3286162=002 9,1264730=001 9,7180260+000 9, 1	***		G# 2 TE	T# 2.0108044m0	04 state				
		10	,5343526+000	9.9074816-00	.0003635+00	.3286102#00	.1264730 m00	.7180260+00	.619
18 1 Gm 2 TEST	***		G# 1 TES	T# 1.6569548	***				
31 1,5277559400 9,89638264001 1,0431383400 6,46887468003 9,12647308001 1,0000000400 8,	***	*	1 G# 2 TES	Te 3.705951	****		Elizabeth management of the second	9 9 9 9 9 9 9 9 9 9 9 9 9 9 9 9 9 9 9 9	
1 Ge 1 TEST 1.4179050=004 ***** 1 Ge 2 TEST 1.4179050=004 ***** 1		m.	,5277659+000	8,5963826*00	.0431383+00	,4688746m00	,1264730m00	00+00000000	434999
18 1 Ge 2 TEST# 1.8569717=004 ***** 32 1.52148714-00 9.8448256=001 1.0072321+000 5.9170824=003 9.1264730=001 1.0000004000 24* 18 1.52148714-00 9.9260266=001 1.0072321+000 5.9170824=003 9.1264730=001 1.00000004000 24* 33 1.5192830+000 9.9250266=001 9.9864173=001 5.7738714=003 9.1264730=001 1.00000004000 24* 34 1.5093312+00 9.9515667=001 9.9788370=001 5.5778854=003 9.1264730=001 1.00000004000 24* 35 1.5093312+00 9.9515667=001 9.9788370=001 5.5778854=003 9.1264730=001 1.0000000000 24* 35 1.5093312+00 9.951567=001 9.98644204=001 5.875896=003 9.5900781=001 1.016410100 24* 35 1.4982202+000 9.8577142=001 9.9847734=001 7.4643647=003 9.5900781=001 1.5513477=000 24* 36 1.4982202+000 9.960259=001 9.9877734=001 7.4643647=003 9.5900781=001 1.5513477=000 24* 37 1.4980489+000 9.9629146=001 9.9905889=001 1.1478232=002 9.5900781=001 2.5527594+000 24* 38 1.48835652+000 8.965487=004 ********** 38 1.48835652+000 8.965487=004 ***********************************	-			4179050	**				
13 1.521 #671 #671 #671 #672 #672 #672 #672 #672 #672 #673 #673 #673 #673 #673 #673 #673 #673		-	3 Ca 2 2 2 2 2 2 2 2 2 2 2 2 2 2 2 2 2 2	1.8569717	****			54 4 4	
33 1.5182830+000 9.928026690014 ***********************************	4	2	DODAT/SATZC	W-8448296=001	1,0072521+0	*9170824¤00	,1264730#00	*000000000	2+2749583=002
1		9 PC	1 2	Datingthto o	04 44 64 60	SO THE STATE	404647		4040
34 1.5093312+000 9.9515667=001 9.9788370=001 5.5567083=003 9.1264730=001 1.0000000=000 2.12 1 GB 2 TEST 2.7137821=004 ***********************************	神ななな		4 GH 2 7 TR	T# 2.8335850m	7 + + + + + + + + + + + + + + + + + + +	DOB11/80//-	Onano/Losts		0000000
1			5093312+000	9.95156670001	9.9788370*00	.53670A3e00	1264730#00	. 000000000000	9.7270293#603
35 1.5036462+000 9.9544044-001 9.9806475=001 5.3778854=003 9.8900781=001 1.0165101+000 2.624304 1	***		1 G# 2 TE	T# 2.7137821*	10年 本本學亦中				
# 1 G# 2 TEST# 2.5722116=004 ***** 36 1,4982202+000 9.8577142=001 9.9844204=001 5.8758096*003 9.5900781=001 1.1633906+000 2.670610 8 1 G# 2 TEST# 2.4791829**** 37 1,4930489+000 8.9602559**** 37 1,4930489+000 8.9602559**** 38 1,4881516+000 9.9629116**** 38 1,4881516+000 9.9629116**** 39 1,4881516+000 9.9629116**** 39 1,48815652+000 8.9656870**** 39 1,4885552+000 8.9656870**** 4 1 G# 2 TEST# 2.2200091***** 4 1 G# 2 TEST# 2.2200091***** 4 1 G# 2 TEST# 2.2200091***** 5 1,4885552+000 8.9656870******** 5 1,4885652+000 8.9656870************************************		35	,5036462+000	9,9544044m001	9,9806475m00	.3778854m00	.5900781*00	.0166101400	.6243
36 1,4982202+000 9,8577142=001 9,9844204=001 5,8758096=003 9,5900781=001 1,1633906+000 2,670610 p 18 1 Ge 2 TEST# 2,4791829=004 ***********************************	***	-	* 2 TE	T# 2.5722116m	****				
# 18 1 G# 2 TEST# 2.4791829#004 ***** 37 1.4930489+000 8,9602559*001 9,9677734*001 7.4648647#003 9,5900781**001 1.5513477*000 % # I# 1 G# 2 TEST# 2.365466**004 ***** 38 1.4881516**000 9,9629116**001 9,9905889**001 1.1478232#002 9,5900781**001 2.5227594**00 % # I# 1 G# 2 TEST# 2.2200094**004 ***** 39 1.4885652**000 8.9654870**001 9.9934015**004 2.33344#&%**002 0.8900784**004 E.8448046**000 #			4982202+000	9,9577142*00	9,9844204.00	.8755096m00	.5900781 .00	1633906+00	,67061
37 1,4930469+000 8,9602559+001 9,9877734+001 7,4648647+003 9,5900781+001 1,5513477+000 % # 1% 1 Ge 2 TEST# 2,3654666+004 +++++ 38 1,4881516+000 9,9629116+001 9,9905889+001 1,14752327002 9,5900781+001 2,5227594+000 % # 1# 1 Ge 2 TEST# 2,2200091+004 +++++	***	82	1 Gm 2 TE	Ta 2.4791829m	****				
# 4	4	-	49304894000	8,9602559°0	9,9877734-00	.4643647m00	* 5900781=00	+5913477+00	.75175
18 14 GB TESTA 2. ZZ2000944004 Parkodoli Tara/SZZZZZZZZZZZZZZZZZZZZZZZZZZZZZZZZZZZ		* M	A PRINT ALDON	0.0420444	20 10 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0	000000000000000000000000000000000000000	000000000000000000000000000000000000000		4×95220-8
3 4 48856524000 8.9656870*004 9.9634015*004 9.34844838000 9.80007948004 5.8448648000 9.80007948065	***) M	TO TO THE PARTY OF THE	T* 9.9200004	7 * YYUZOGY#U	14/9236490	0041070044	*2227294400	104400//0/*
		. Po	48356524000	8-9656A70R0A+	9.9034015en	747648300	ROPHTRA BOOM	S44 ROALABO	6.774 A 6 Q 3 m 0 B 3

* * *	# 4 #	G# 2 TEST# 2,011405 4794157*000 9,9122504*0 G# 1 TEST# 2,493264	6,5917200m002	9,5900781=001	1,7442605+001	1,1938159*002
* * * * * * * * *	# 4 H	.8040680 69434=00 .1958531	5.7492100*003	9,5900781*001	1,0000000*000	2,1265644*001
* * *	# 4 H	4723224+000 9,7343919-001 1,	3.6995102=003	9.5900781=001	1.0000000*000	4,9911045=002
	43	4689550 + 000 9 99544111001 1,0008	2,9249081-003	9,5900781=001	1,000000000001	5.1961128-003
	4 4	57268+000 9,9702909=001 9,9	2,7972014#003	9,5900781*001	1.00000000001	2,1519208*003
4 d	# 4 i	Gm 2 TEST# 1,5037825m004 4626296#000 9,9767140m001 9,	2,6830751#003	9,5900781#001	1,000000000000	£.1609664*003
2	46	4596666+000 9,9780692=001 9,	2,5689133#003	9,5900781*001	1.00000004000	1.2411545=003
性 性 性 体	# 4 7	EST# 1.413562 0 9.9786984*0	3.5112302003	9,5730206#001	1,4300000+000	£,3670241*003
**	# 4 60	41541+000 9,9795448=001 9,9	3,3282589#003	9,5730206*001	.430000#00	1,3276032+003
**	H 4	EST# 1.2889303+004 ***** 0 9.9804581+001 9.9939309+0	3.1139198#003	5730206*00	.4300000+00	
* * * *	H EV	1 Gm 2 TEST# 1.1746841=004 ***********************************	.0584175m00	5730206#00	0040000000	4343287*0
教教教教		1 G# 2 TEST# 4.1305426-004 ***** 4420316-000 0.9824886-004 0.0046182-0	9464004-00	R770904-00		0000 BA-00
***	*	1 G* 2 TEST* 1.1082419*004 ****			*********	*********
1	52	4449100+000 0,9831012*001 9,	1,8462891#003	9,5730206+001	1.0000000.000	£.4630793-003
	ر ا ا	9,9839328-0	1,7470336#003	9,5730206=001	1,0000000000000000000000000000000000000	1,4799112=003
**		1 G# 2 TEST# 1.0026808*004 ***** 444104224000 9.98444488004 0.000573880	0040428079	4424057	10 THE RESERVE OF THE	407004
		4392332+000 9,9853453-001 1,0000254	8067458 = 00		1.1608654+000	167-00
	56	0 9,9860644=001 1,0000517=0	726	4624057	1.54004364000	52332*0
		.4346347*000 9.9887849*001 9.9988340*00	2628925400	4624057*00	5.1993116+000	.0049126-00
		,4334356+000 9,9842589 001 9,9991542 00	5015903 00	. 4624057e00	1,4310121+001	,4895292*00
位 作 作	4 Q	1 G# 27338+000 9,8737030001 1,000	1,88164709003	9.46240578001	1.0000000+000	2.1085620-002
	61	4319057+000 9,9650765-001 1,00	026171*0	.4624057#00	.000000	.5674005
		.4311227+000 9,9880181-001 1,01	726105#00	4624057*00	.000000000	2658
	63	,43035034000 9,9927825m001 9,99	869054#00	4624057000	.00000000.	.1045610*
	40 A	.4296121+000 9.9937309+001 9.99	6.2783131¤004	4624057*00	.0000000	3,7511783=004
	200	42856774000 7,47595700m001 7,97,428584854001 0,000	55445#UU	A171081=00	44000000	8051700
	67	2753924000 9,9947452F001 9,99	524929 • 00	9,81719819001	300	4
	89	,4268973+000 9,9929813=001 9,9	410629*00	8171981 .00	4300000+00	,5777033
	69	,4263128+000 9,9948778=001 9,99	4,9815141m004	9,8171981 .001	00+00	,9512606=

	70	,4257653+00	956970-00	9992594-00	,6686425m00	8171981 * 00	00000000+	,5534290=00
	/1	4252357+00	,9947969=00	9995427 • 00	,5190272=00	8171981=00	.0000000	.7457310*00
	15	4247396+00	,9959814#00	9998016*00	.3132814#00	8171981+00	.0170102+00	,8202407-00
	7	4242593#00	9948457-00	00-5866666	.8098750.00	.8171981*0C	·1679099+00	.1527885±00
	4	4236111400	00-9660966	0000111+00	,0470024 = 00	8171981 . 00	.5718723+00	.0054266-00
	2,2	,4235856+00	9945687=00	999995-00	.6445575-00	,8171981 # 0C	* 6171420 ± 00	,4277734±00
	10	4227883+00	9957523-00	9666	2,1149195*003	9,8171981 - 001	6,1709454+000	8,8954000+004
4	- 1		7,7844004-001	On a cratter of	00+4021400+	8171951#00	.8566669*00	00*08#/696
* *	Ŧ []	7 0	7848044	4444				
	-	1.4216070+000	9.8263406-001	1.0399164*00	5429497m00	984 80 0	00+00000	7282319*00
	29	1,4215118+00	9937867-00	0023387+00	7597861500	8171981#06	000000000000000000000000000000000000000	9600482=00
	80	1,4213619+00	9969095=00	0003852+00	8482635+00	817198100	00+0000000	9424920*00
	81	1,4212024+00	9974048-00	0001385+00	5345622#00	8171981.00	00+0000000	9797317.00
	82	1,4210480+00	9977126-00	000038400	0562867 + 00	3027242#00	4300000+00	2228776=00
	8	1,4208975+00	9979647=00	0000527+00	9387591#00	3027242#00	4300000+00	5618958*00
	4	1,4207502+00	9981110-00	0000125+00	8541942#00	3027242*00	4300000+00	0137124.00
	00	1,4206061+00	9982944-00	00828400	2485432*00	3027242=00	000000000	6110556-00
	9 1	1,4204660+00	9983791=00	9997410 . 00	2048511=00	3027242#00	000000000	3618819=00
	87	1,4203300+00	9984842+00	9996417 .00	1620696*00	3027242*00	00000000	1574627=00
	80	1,4201983+00	9985769=00	9996182+00	1208444#00	3027242#00	0000000000	0412520-00
	00	1,4200710+00	9986591-00	9996872-00	0973643#00	6409409*00	0166997+00	0180163-00
	06	1,4199482+00	998758200	9997472=00	2095021 00	6409409=00	1643996+00	889688*00
	16	1,4198301+00	9988544=0(9997815 00	552389500	6409409#00	5558974*00	2747863-00
	0	1,4197169+00	9989699 01	9997716*00	4276916#00	6409409#00	5432997+00	0167039#00
400	56	1,4196096+00	9991266#00	9997019-00	1086547m00	6409409*00	6467117+00	7524259 00
	4 1	1,4195105+00	9992182=00	9995366+00	5945322#00	6409409#00	9109036*00	1833464-00
	2 0	10+902+614	9992336#0(9994/89=00	0364680=00	6409409=0	000000000	4536115=00
	000	1.419640140	9993286-00	9995330+00	8154809=00	6409409#0	000000000	0438150=00
) a	1, 41, 70, 01, 01, 01, 01, 01, 01, 01, 01, 01, 0	0 5 1 5 5 5 5 5 5 5 5 5 5 5 5 5 5 5 5 5	000000000000000000000000000000000000000	5991922 B 0 0	6409409#00	0000000	6367372m00
	0 0	1,41914854000	9,99936/28001	9,999997	0,02//434#000 7. 47#040-00F	9,64094094001	1,02416314000	6.2519180 = 0.00 7.04 6 4 0 1 0 0 0 0
		1.4190484+0	9993942	9997906=00	447207=0	440040040	040404040	0447057
	101	1,4189816+0	9994168#0	9997920+00	4032787400	6409409#0	2623419+00	7518970=00
		1,4189194+0	9994611=0	9997474-00	8244153 00	6409409#0	6807014+00	8633120-00
	0	1,4188675+0	9995428#0	0002200666	3603973#00	6409409#0	.00000000	5246638 00
	-	1,418817400	9996032#0	9997186*00	2118775000	6409409*0	0000000000	1545722-00
	W 14	1,418768840(,9996126m0	9997680 • 00	0772978#00	640940900	.00000000	,5546408**00
	L 2.	1.418722040	9996194=0	9998153*00	0849286#00	6409409=0	0380914#00	9389365.00
		1 1 41 60 / 1 40	0.0294684	004/168666	3988798#00	6409409m0(4236408400	2566004-00
	₩.	1 . 61 60 GGV + 0	9996335=0	999871300	0910633=00	6409409*0	.9986668*00	3788103-00
	_ •		0-0169666	000000000000000000000000000000000000000	7340477 m00	6409409=0	3774189+00	0550564=00
		T * T C C C C C C C C C C C C C C C C C	08440/4/4	000018000	0.4696504	040740980	0000000	0440510#00
		1.410202024	08/01/664	0.5000	8108502#00	6407409#0	.00000000	87//230=00
		1 1 1 1 1 1 1 1 1 1 1 1 1 1 1 1 1 1 1 1	0502/444	0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0	700170#00	040840840	0000000	2528237=00
	7 7	1 44844040	0 1 4 4 4 4 0 C O	0.0000000000000000000000000000000000000	047846/80	040740940	040040600	15215045#B
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(v) C	· · · · · · · · · · · · · · · · · · ·	1 UN1 FLUXE 1,78 1 UN1 FLUXE 1,77 1 UN1 FLUXE 1,74 1 UN1 FLUXE 1,78	1 UHL FLUXH 1,69833589 1 UHL FLUXH 1,59661681 1 UHL FLUXH 1,52404074	1 CEP 7 COX P. 44148811111 CEP 7 CX P. 4414881111 CX P. 4495454111 CX P. 449545111 CX P. 44451111 CX P. 4445111 CX P. 444511 CX P. 44451 CX P. 444511 CX P. 444511 CX P. 444511 CX P. 444511 CX P. 44451 CX P. 444511 CX P. 444511 CX P. 444511 CX P. 444511 CX P. 44451 CX P. 444511 CX P. 444511 CX P. 444511 CX P. 444511 CX P. 44451 CX P. 444511 CX P. 444511 CX P. 444511 CX P. 444511 CX P. 44451 CX P. 444511 CX P. 444511 CX P. 444511 CX P. 444511 CX P. 44451 CX P. 444511 CX P. 444511 CX P. 444511 CX P. 444511 CX P. 44451 CX P. 444511 CX P. 444511 CX P. 444511 CX P. 444511 CX P. 44451 CX P. 444511 CX P. 444511 CX P. 444511 CX P. 444511 CX P. 44451 CX P. 444511 CX P. 444511 CX P. 444511 CX P. 444511 CX P. 44451 CX P. 444511 CX P. 444511 CX P. 444511 CX P. 444511 CX P. 44451 CX P. 444511 CX P. 444511 CX P. 444511 CX P. 444511 CX P. 44451 CX P. 444511 CX P. 444511 CX P. 444511 CX P. 444511 CX P. 44451 CX P. 444511 CX P. 444511 CX P. 444511 CX P. 444511 CX P. 44451 CX P. 444511 CX P. 444511 CX P. 444511 CX P. 444511 CX P. 44451 CX P. 444511 CX P. 444511 CX P. 444511 CX P. 444511 CX P. 44451 CX P. 44	1 Can FLOX# 1, 1423006 1, 120909966 1, 12084848 1, 12084848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120848 1, 120	1 CH1 FLUXE 7,8949261111 CH1 FLUXE 6,6958771111 CH1 FLUXE 5,40648451	1 JR1 FLUX= 4,1595168** 1 JR1 FLUX= 2,9275785* 1 JR1 FLUX= 1,7229637*	Jai FLUX# 5,87514624 S COPTED FLUX FROM LUN	0 JMAXIN= 20 IMAXOU∓= 20 JMAXOUT=	1 LEL FLUXE 3,5455439494001 LEL FLUXE 3,545534394094001 LEL FLUXE 3,54553439409400	1 Jan Fluxa 3,5455339400 1 Jan Fluxa 3,5455339400 1 Jan Fluxa 3,545533940	1 LET FLUX# 3,5355339#00 1 LET FLUX# 3,5355339#00 1 LET FLUX# 3,5355339#00	1 JR1 FLUXE 3,5355339#00 1 JR1 FLUXE 3,5355339#00 1 JR1 FLUXE 3,5355339#00	1 LET FLUXE 3,53553353503010 LET FLUXE 3,5355353535353510 LET FLUXE 3,53553535450400	1 LET FLUXE 3,545533991 LELUXE 3,5455339339	1 JM1 FLUXH 3,5355339+00	TAS COPIED FLUX FROM LUN 49 ONTO LUN 48	
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PAGE NO.

PROB, NO. 1,02000+001 20CANDID RT FULL CIRCLE 1, 2 GROUPS, 20x20 HESH

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40	575140	004	2,63346#003	00
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20	57515a	004	2,63346=003	20

8. Sample Problem 8

a. Description

(1) Problem Type

Real keff calculation

(2) Configuration

(a) Geometry

хy

(b) Region Definition

The reactor consists of three regions as follows:

Region No. 1: Internal thermal column (water).

Region No. 1: Core region composed of enriched uranium, stainless steel, and water.

Region No. 3: Reflector composed of beryllium and water.

(c) Mesh Definition

x direction: nine points

y direction: nine points

Buckling $(B^2) = 0.0027416$

(d) Boundary Conditions

Left: $\phi' = 0$.

Right: $\phi = 0$.

Bottom: $\phi' = 0$.

Top: $\phi = 0$.

(e) Number of Energy Groups: 18.

(3) Convergence Criteria

 k_{eff} difference = 10^{-5} .

 k_{eff} bounds = 10^{-5} .

Sum of flux difference = 10^{-5} .

Up-scattering = 10^{-3} .

Maximum number of up-scattering iterations = 0.

Output Listing	
	ng
Output	
	Output

K CALC(UPSCAT-NO INNERS), XY, 18 GROUPS, 9x9MESH 7,10000+000 2DCANDID PROB. NO.

M PAGE NO.

REACTOR GEOMETRY

RIGHT BOUND

BOTTOM BOUND

BOUND

0.0000+0000 BOUND LEFT

2D-XY WITH NO SYMMETRY

GEOMETRY TYPE

3,6000+001

0.0000000000

3,6000+001

REGION BOUNDARIES

REGION NO.

LEFT BORY. 0.000000000

RIGHT BDRY. 3.6000+001 2,1000+001

2.1000+001

TOP BORY. 3.6000+001

BOTTOM BDRY.

K CALC(UPSCAT-NO INNERS), XY, 18 GROUPS, 9x9MESH 2DCAND I D

7,10000+000

PROB. NO.

MESH DEFINITION

PAGE NO.

INCREMENT METHOD - X(R) DIRECTION

INCREMENTS X(R)

M

0000+0000

INCREMENTS X(R) 6.0000+000

INCREMENTS X(R) 2.1000+001

3,6000+001

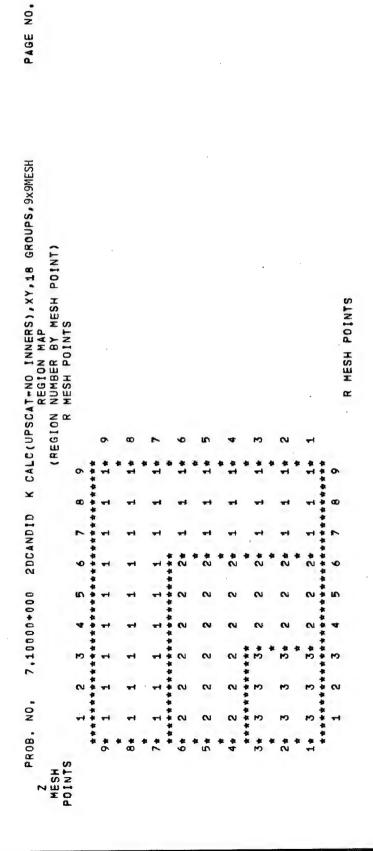
X(R)

INCREMENT

X(R)

X(R) PLUS DELTA X(R) FROM LEFT TO RIGHT OF REACTOR

2,0000+000 4,0000+000 6,0000+000 1,1000+001 1,6000+001 2,1000+601 2,6000+001 3,1000+001 3,6000+001 MESH NO. ABSCISSA



Z MESH POINTS

ø

K CALC(UPSCAT-NO INNERS), XY, 18 GROUPS, 9X9MESH MESH DEFINITION **2DCANDID** 7,10000+000 PROB. NO.

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PAGE NO.

INCREMENT METHOD - Y(Z) DIRECTION

Y(Z)	INCREMENTS	4(2)	Y(Z) INGREMENTS	Y(Z)	Y(Z) INCREMENTS	Y(Z)	Y(Z) INCREMENTS	(Z) A
0.0000+0000	ю	6,0000+000	ю	2.1000+001	ю	3,6000+001		
CMAX # 9								

Y(Z) PLUS DELTA Y(Z) FROM BOTTOM TO TOP OF REACTOR

5.3616415-002

1,5239983+000

9,2779316=001

4,5160331*003

1.0273433+000

,0268248+000

9,6999782=001

.6906845-001

1,1855687+000 1,1585613+000 1,1344588+000

,0226849+000

9,2779316-001

.2779316-00:

2,4807272±003 3,0991304±003

1572362+000

5.6812614-002

5.2470464-00 4.0878474-00 3.4017724-00 1.9438036-00 1.779692-00	1.4539751-00 1.3202855-00 1.1571603-00 9.7089145-00	5.21n3328-00 2.23n6894-00 2.6471622-00 2.9454448-00 3.1430401-00 3.2033749-00	2.6732010-00 1.8646090-00 1.6705683-00 6.6052310-00 5.0775032-00	5.0810198-00 4.7983453-00 4.4841820-00 3.6918401-00 2.592222-00 3.7098155-00 1.7742848-00	400444044
. 8058668+00 . 1361863+00 . 0000000+00 . 0000000+00	1572362+00 5239983+00 4036257+00 8058668+00	.00000000+00 .0000000+00 .0160608+00 .1572362+00 .523983+00	. 6058668+00 . 1361863+00 . 0000000+00 . 0000000+00 . 0000000+00	.1572362+00 .5239983+00 .4036257+00 .8058668+00 .1361863+00 .0000000+00	1.0160608+000 1.1572362+000 2.4036257+000 4.8058668+000 1.1361863+000 1.0000000+000 1.0000000+000
.2779316-00 .2779316-00 .2779316-00 .2779316-00	2779316+00 2779316+00 2779316+00 2779316+00	.2779316*00 .2779316*00 .2779316*00 .2779316*00 .2779316*00	2779316-06 2779316-06 2779316-06 2779316-06	2779316-90 2779316-90 2779316-90 2779316-90 2779316-90 2779316-90	9.27793161001 9.27793161001 9.27793161001 9.27793161001 9.27793161001 9.27793161001
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.0240860+00 .0157558+00 .0193509+00 .0101893+00 .0092406+00	.0072837+00 .0066783+00 .0059781+00 .0054918+00	.0031631+00 .0022096+00 .0021195+00 .0020644+00 .0020036+00	. 0015639+00 . 0017639+00 . 000784+00 . 9994424+00 . 9994424+00	.0000372+00 .0000658+00 .0000423+00 .9998768-00 .9999771-00	1.0000324+000 1.0000456+000 1.0000533++000 1.0000533+000 1.000034+000 1.000034+000 1.0000245+000
7161557*00 7487736*00 8533320*00 9075131*00 9146095*00	9274993800 9347542*00 9440654800 9578287*00	9795280100 9997893100 9947235100 9911898100 9885739100	999967800 99919470800 99919471800 99936739800 9943649800	.9952910*00 .9958599*00 .9962055*00 .9967309*00 .9972846*00 .9982028*00	9,9983428 9,9984071 9,9984071 9,9985758 9,9997458 9,9996682 9,9997659 9,9997659
.11022009+00 .1022009+00 .1041487+00 .1063391+00 .1085791+00	.1130016+00 .1151726+00 .1173271+00 .1194827+00	1234704+00 12551111+00 1255837+00 1279047+00 1390780+00	1130 1311 1311 1311 1311 1311 1311 1311	1308623+00 1308624+00 1309040+00 1303472+00 1302167+00 130036+00	1.12989666 1.12989166 1.129813860 1.12978466 1.129789460 1.129789460 1.1297894600 1.1297894600
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9 JHAXOUTE

9 IMAXOUTS

9 JMAXINE

IMAXINE

PROB	• NO.	7.10000+	0+000 2DCANDID	K CALCIUPSO	NI ON-LA	CALCIUPSCAT-NO INNERS), XY, 18	GROUPS, 9x9MESH	PAGE	NO. 35
# N T T	RATION K=E 1.129	HISTORY FFECTIVE 8051+000	K-LOWER 9,9998009-001	K*UP 1,0000165+	PER CHA	NGE IN PHI 654417-007	SIGMA 1,00000000001	ALPHA 1*0000000+00	ERRLAM 3.6408281*005
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CANDID2D	HAS COP	IED FLUX	FROM LUN 42	ONTO LUN 4	30				
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IMAXIN= CANDID2D	9 JMA	JMAXIN=	9 IMAXOUT# SECTIONS ONLY (9 JMAXOUT= ON LUN 48	0.				

9. Sample Problem 9

a. Description

(1) Problem Type

Real $k_{\mbox{eff}}$ calculation

(2) Configuration

Same as Sample Problem 8

(3) Convergence Criteria

 k_{eff} difference = 10^{-5} .

 k_{eff} bounds = 10^{-5} .

Sum of flux difference = 10^{-5} .

Up-scattering = 10^{-3} .

Maximum number of up-scattering iterations = 4.

b. Output Listing

PROB, NO. 7,40000+000 20CANDID K CALC(UPSCAT-4 INNERS), XY,18 GROUPS,9X9MESH

3

PAGE NO.

ERATION	PEFFECTIV	-LONE	BODOE	HANGE IN PE	E	PE	RRLA
	8,3038143+000	9,5068875-003	1,2558345+002	4.3654464+000	1.00000000000	1.0000000+000	4,2557394+002
11,522							
N	,2631728+00	15176762-00	2459383+00	+2831032+00	0.0000000	00+0000000.	0+4071409.
2	4393983+00	12170893*00	.1480156+00	*1803486*00	.0000000.	.000000000	.7263067+00
4	3612331+00	2219240-00	8938318+00	.7305079#00	0.0000000	004000000	3716394+00
S	3882127+00	.8890472*00	6586894+00	.4066188#00	00000000	004000000.	0697847+00
9	3329679400	.3182750=00	5267534+00	.1511913000	0.00000000	00*0000000	9492586+00
~	1414289+00	7979081=00	3791054+00	4523168m00	0.000000	00+0000000	9931457-00
œ	840	7,2671541-001	1,2529651+000	2,1837922=001	1.00000000+000	1.000000+000	524971=00
O *	4695024+00	,6962049=00	1587620+00	.1303869e00	0.0000000	00+00000000	8914151-00
10	,0658863400	.0241793-00	0915385+00	11092755=00	.0000000.	.0000000	8942062-00
11	,6566499+00	,0460113=00	.0429081+00	*0493433m00	0.0000000	.00000000	3830701-00
12	,2605305+00	10731056=00	.0068928*00	.9421047=00	0.0000000	000000000	9958227.00
13	8892950+00	1348032-00	7966078=00	.7993827 00	04000000	00000000	6618046 00
4	,5495610+00	,2260274=00	,5875197 * 00	,6360260m00	0+00000000	.00000000	3614923*00
5	2443466+00	*3419074-00	,7687034.00	4645363×00	0.0000000	00+0000000	4267960+00
16	9742723+00	,3680301-00	9834623 00	3140674#00	951791100	.0154874+00	6154322-00
17	,7384542+00	.3949307-00	.0186389+00	,2983286#0D	9517911#0	1508751+00	.7914588 nO
4	,5356010+00	,4476962=00	.0330864+00	,4265885 » 0 0	951791100	4964344*00	,8831675.00
19	,3661121+00	,5755470=00	.0366544+00	.7166436m00	9517911=0	.2906058*00	7909970=00
20	,2371086+00	,9033539+00	,0268595+00	*0676628m00	951791100	.2387819*00	3652409=00
4	,1812599+00	*0258330*00	0064553+00	,2653688m00	951791100	*3283273*00	.0387196 .00
22	1909983400	8636850=00	2287559+00	.7700904e00	9517911=0	00.0000000	4238744.00
23	1934157+00	3431022-00	,1451811+00	.5280497 m00	,9517911=0	00+00000000	1087092-00
2	1934611+00	,4681090 m	,0930245+00	,6458479#00	9517911*0	0000000000	,4621365-00
22	1911389+00	,5364743=00	0598116*00	.0981404m00	9517911=0	*0000000	0041641600
50	1869251+00	,5940398=00	,0382765+00	*7227147#00	951791100	00+00000000	,8872538=00
27	1812550+00	,6541581=00	0239937+00	4922435*00	,951791150	00*0000000	8177849=00
28	1745891+00	,7173442"00	0143129+00	.368828/#00	951791100	.00000000	.2578492 00
ON I	,1673815+00	,7649582-00	,0076165+00	.307876700	9517911#0	00+00000000	.1120692=00
30	1600268400	8008421 - 00	,0028986+00	*2721646#00	*9517911*0	*00000000	.2814378.00
31	1528564*00	,8276021"00	.9976644=00	,2380115#00	951791100	.00000000	.7006228-00
20	1461332+00	8483546=00	.0017720+00	11936027#00	951791100	.00000000	,6936550=00
9	1400524400	,8626552°00	.0034291+00	.1525128 0 0	951791100	*0154874*00	,6863543-00
34	,1347484+00	8814512-00	0046759400	,2213406 m00	951791100	1508751+00	.6530793.00
35	1303427+00	8944319000	,0052510+00	4387388 p0	.9517911*0	.4964344±00	.5407840=00
91	1270362+00	.9198381°00	*004924400	.8539977.00	,9517911#0	*2906058*00	.3013641=00
37	1252417+00	,9383351=00	.0039839+00	,3693536 00	951791100	.2387819+00	0750440.00
00 t	1256701+00	,9567048¤00	,0032277+00	8496889#00	951791100	273+00	5572486m00
65	1274371400	,9800849*00	.0031026+00	,3268429m00	951791100	00+0000000+	*0941444*00
4	10001474CD	.0047227ens	000745000	. 40.784.00-0n	0547044		TAKOLA 17.08

4	•	1299427400	9900472=00	047303740	7477485	054 7044 = 00		0.448 n
4	-	1307741400	9884956*0P	004025040	1404070*00	054701450	201010101010101010101010101010101010101	ES 6264-00
7.7	• •	1 24 44 044 0	00-0007.00	0440400	100000	051704100	100000000000000000000000000000000000000	200000000000000000000000000000000000000
2	-	000000000000000000000000000000000000000	1207470711	00105000	1102186800	921/911011	1208721 * 00	70/031*00
4	ri	10104040101	, 88/5954m00	00136880	3267640#00	951/911 00	4964344+00	5092692 m 0 0
4	*1	1320453+00	9888784-00	0010993+0	7590498=00	9517911=00	2906058+00	114944 = 00
4		1319526+00	9911761-00	00007504+0	3639090#00	9517911 00	2387819400	327762-00
47	+1	,1315069+00	9922560-00	0001502+0	3472385#00	9517911=00	3283273+00	2464429 a 0 0
4	+	1309944*00	9911703-00	00000000	8038822 × 00	951791100	00000000000	5365315=00
49	*	1305617+00	9982264#00	0000515+0	7155991=00	951791100	0000000	888786*00
20		1301952+00	9932596#00	0001800+0	0059867 00	9517911.00	0000000000	5408051=00
51		1298895+00	9940738-00	0002712+0	408970500	951791100	00+000000	5384261-00
52		1296431+00	9946888*00	0003317+0	9052465#00	147848500	0158320+00	5280538.00
53	**	1294542+00	9951861-00	0003636+0	9931683#00	1478485=00	1546906+00	4497704=00
54	+1	1293239+00	9956824=00	0003594+0	7314543000	147848500	5128835+00	9111580 = 00
50	wi	1292594+00	9963069=00	0003193*0	1949894#00	1478485#00	3572358+00	3859101.00
56	w t	1292780+00	9973143=00	0002493+0	0309717#00	147848500	5624124+00	1782397#00
21	-1	1294040+00	9994022-00	0002168+0	1259751#00	1478485#00	9205939*00	7657722=00
58	**	1295379+00	9990194.00	00003169+0	3909836#00	1478485 = 00	000000000	1493753-00
29	**	1296504+00	9998503=00	0002629+0	0521706#00	1478485*00	00+000000	7789274000
09	#1	1297453+00	9995178-00	0002168+0	8507498=00	1478485*00	00000000	5503729-00
61	w i	,1298243+00	.9992793=00	0001829+0	7148504=00	147848500	0158320+00	5501792-00
62	**	,1298876+00	9991283=00	0001584+0	7620939#00	147848500	1,546906+00	4558589-00
63	44	1299349*00	9990877-00	0001377+0	0139285*00	1478485	5128835*00	2894613000
49	₩.	,1299	9.9991654=001	,0001163+0	054853#0	1478485 = 0	3572358+00	9977705=00
62	T	1299692+00	9993334-00	,0000873+0	1071750#00	147848500	5624124+00	5395039-00
99	w f	1299421+00	9994249=00	00000310+0	8557506*00	147848500	9205939+00	8528264.00
29	- 1	1299067+00	9991329=00	0000365+0	5172700 m 00	1478485 00	00000000	2322664×00
90 9		+1298777+00	.9992749=00	0000013+0	4476171m00	1478485=0	000000000	9777819-00
69	*1	1298536+00	. 9994063m00	0000149+0	8455701 ** 00	1478485#0	.0000000	4234529-00
7.0	+1	1298340+00	,9995064m00	0000020040	4446257=00	147848500	.0158320+00	9922098-00
7.7		1296165+00	9995801-00	0000240+0	\$5276195¤00	1478485 0	1546906400	5985892#00
72	**	12980/3400	9996422=00	0000246+0	1266791=00	147848500	5128835+00	0383987 - 00
13	₩.	1296009+00	. 9997050=00	,0000222+0	3090554000	147848500	.3572358+00	1744195=00
4		1298007+00	,9997862m0	,0000177+0	8226787 . 00	147848500	,5624124*D(9051665-00
2	-1	1296085+00	0999339=00	0000153+0	3724914=00	147848500	.9205939+00	1928339=00
76	πÍ	1298180+00	\$999453*0C	00000228+0	7108544w00	1478485 = 0	.00000000	1256408=00
11		1298258+00	9999837=00	.0000191+0	4472486m00	1478485=0	00000000	0688254=00
78		,1298323+00	9999627=0(,0000157+0	295250000	1478485*0	.00000000	9388855=00
29	v t	1298377+00	9999468=01	.0000131+0	1933961#00	147848500	*0158320+0	8398132*00
80		1298419+00	9999371+00	.0000112+6	,2203299#00	1478485=0	1546906+00	7488142-00
81	wł.	1298451+00	9999350=0(.0000096*	3873963m0C	1478485 0	5128835+0	6113307-00
CVI 000	*-1	1298469+00	9999409m0(+09000000	,7152460m0C	147848500	.3572358+0	3904093 = 00
83	**	1298471+00	.9999530n0	1465000000	1255766#00	1478485 0	5624124+0	0581891-00
80	wi	1298451+00	. 9999593	0000194	.5528437#00	784	9,9205939+000	6.0067105-006

1

APPENDIX B

LGOEDIT2

1. Identification

a. <u>Title.</u> Load-and-go tape edit
b. <u>Language.</u> FORTRAN-63, one COMPASS subroutine
c. Machine. 3600

d. Programming Sanford Elkin, D. B. Taylor, A. L. Rago, J. Zapatka, G. J. Duffy

2. Purpose

LGOEDIT2 provides a facility for editing and updating load-and-go tapes for the 3600 by means of deleting, replacing, and inserting binary decks and control cards at specified positions on the tape. LGOEDIT2 is expected to be useful primarily when the LGO tape contains MAIN, OVERLAY, and SEGMENT control cards, or when the order of loading is important.

3. Usage

LGOEDIT2 creates a new LGO tape by editing an old LGO tape with information from an edit tape (the edit tape can be the standard input unit). LGOEDIT2 first reads a card from standard input using the format 3I2. Colums 1 and 2 must contain the logical unit number of the old LGO tape, colums 3 and 4 the logical unit number of the new LGO tape, and columns 5 and 6 the logical unit number of the edit tape. (All must be numbers from 1 to 49, except that columns 5 and 6 can contain 60.) All further information is obtained from an edit tape.

LGOEDIT2 recognizes five control cards. These are:

Column	1-4	<u>5-8</u>	9-16
	7DEL		NAME
	7REP		NAME
	7INS		NAME
	7BYP	.	NAME
	7DLC		NAME

[†] Control Data Corporation.

where NAME is the BCD name of the subroutine of the LGO tape, and it must be left justified beginning in column 9. Columns 5-8 are not used.

LGOEDIT2 reads a control card from the edit tape and copies the old LGO tape onto the new one until it finds subroutine NAME. It then does the operation required by the control card and returns to the edit tape to read another control card. When it finds an end of file on the edit tape, editing is completed, and the remainder of the old LGO tape is copied onto the new one.

The control cards cause LGOEDIT2 to perform the following actions:

⁷DEL will delete subroutine NAME.

REP will replace subroutine NAME with binary decks from the edit tape, until another control card is read. Note that NAME may be replaced by many subroutines.

7/9INS will insert binary decks from the edit tape onto the new LGO tape <u>after</u> subroutine NAME has been copied. Note that many decks may be inserted after NAME.

⁷BYP will copy the old LGO tape onto the new LGO tape, up to and including NAME. The reason for this is that on an overlay tape, the same subroutine name may occur in several overlays or segments. If we desire to edit some of these subroutines, ⁷BYP allows us to bypass an earlier occurrence of NAME.

70LC will delete all cards between the IDC card of subprogram NAME and the final card of the subprogram preceding NAME.

This DLC card was installed mainly to remove or replace loader control cards (e.g., bank, main, overlay, segment) or octal corrector cards.

LGOEDIT2 never looks at the names of the subroutines on the edit tape, but only at the NAMES on the control cards.

4. Restrictions

The NAMEs referenced on the edit tape must appear in the same order as on the old LGO tape, because LGOEDIT2 will not backspace the tapes to search for missing names. Both the edit tape and the old LGO tape must be terminated by an end of file. We cannot follow a REP card with an INS card with the same NAME on both.

5. Timing

A large LGOEDIT2 job was timed with the following resulting statistics:

a. LGOEDIT2 portion of job:

Number of DLC control cards = 2

Number of DEL control cards = 5

Number of REP control cards = 8

Number of binary subroutines transferred = 143.

LGOEDIT2 time = 1 min 35 sec (22% of job time)

b. FORTRAN called nine times to compile 11 subroutines (1650 source cards).

FTN time = 2 min 50 sec (40% of job time)

c. Translation of LOAD/GO tape to Overlay tape.

Translation time = 2 min 44 sec (38% of job time)

Total time for all three runs in job = 7 min 9 sec

6. LGOEDIT2 Output Edit

LGOEDIT2 will produce lists of the subprogram names on old and new LGO tapes, along with a record of the edit procedures carried out and attempted.

7. Typical LGOEDIT Job

⁷Job,999,204051,25

Acct. card.

 $\frac{7}{9}$ /Mount Tapes L5882(44) and L8878(49)

7EQUIP, 20=60

7EQUIP,21=61

7EQUIP,29=(SNARG·OVERLAY·29,1),SV

7EQUIP,39=(SNARG.OVERLAY.LOAD.AND.GO,1),SV

⁷EQUIP,44=(SNARG·XLIBIT·DEBUG,6)RO,SV

ZEQUIP,49=(SNARG·OVERLAY·LOAD·AND·GO),RO,SV

FILE,14

7REP....SNARGID

FILE END

7FTN, L, A, X=14.

SNARGID Source Deck

SCOPE

⁷FILE,14

?REP...INITILIZ

FILE END

FTN, L, A, X=14

INITILIZ Source Deck

SCOPE

LGOEDIT2 Binary Deck

7RUN, 60, 10000, 7,

4 9 3 9 1 4

78END OF FILE card for end of this run

LOAD,39

7RUN,60,10000,7,

Data Deck

Blank Card

⁷End of File card for this job.

To place the following on LUN 14:

7REP....SNARGID

SNARGID binary deck.

REP...INITILIZ

INITILIZ binary deck.

7(EOF)

→ Transfers all binary card records from 49 to 39 replacing only those belonging to SNARGID and INITILIZ in the process. These two are transferred from 14 to 39 in lieu of their old versions. After the LGOEDIT2 run, LUN39 was loaded to run the program.

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